博士論文

# Silver phosphate microcrystalline film-based photoanodes for photoelectrochemical water oxidation

(微結晶リン酸銀膜光電極を用いた水の酸化反応)



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## Abstract

Since the first report of electrochemical photolysis of water using a TiO<sub>2</sub> electrode under ultraviolet light irradiation, photoelectrochemical (PEC) water splitting process has captivated tremendous research interest as a promising technology for sustainable production of the clean energy carrier  $-H_2$ , from solar light and water. The advances in PEC technology substantially depend on the development of highly stable and efficient photoelectrodes, especially photoanodes, since the oxygen evolution reaction (OER, *i.e.* water oxidation) happening on the anode is usually more thermodynamically complex and dynamically sluggish than hydrogen evolution reaction (HER, *i.e.* water reduction) happening on the cathode. The photoanode allowing the occurrence of water oxidation is thus more dominant for development of overall water splitting technology. Traditional TiO<sub>2</sub>-based electrodes usually suffer a lot from the low visible-light absorption induced by their large bandgaps. Meanwhile, sulfides and (oxy)nitrides show poor stability over long-term operation in actual conditions. These have highlighted WO<sub>3</sub>, BiVO<sub>4</sub>,  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>, and Ag<sub>3</sub>PO<sub>4</sub> as novel alternatives to the above-mentioned photoanode materials towards efficient and stable water oxidation. In 2010, as an important benchmark report, Yi et al. demonstrated that Ag<sub>3</sub>PO<sub>4</sub> as the photoanode material, with the presence of AgNO<sub>3</sub> in electrolyte as an electron scavenger, could offer much higher photocatalytic activity for water oxidation than other well-established anodic catalysts including BiVO4 and WO3. This superior performance could be further improved through the strategies like facet and morphology engineering. However, till now still not much work has been reported on using  $Ag_3PO_4$  as the photoanode material for PEC water splitting due to the difficulty in fabrication of highly adhesive Ag<sub>3</sub>PO<sub>4</sub> layers on conductive substrates as membrane

photoanodes. More importantly, the two major intrinsic shortcomings of Ag<sub>3</sub>PO<sub>4</sub>, *i.e.* the poor structural stability and low bandgap absorption resulted from the native defects of silver vacancy ( $V_{Ag}$ ), silver interstitial (Ag<sub>i</sub>), and simultaneously decomposition-induced precipitation of Ag<sup>0</sup> nanoparticles (NPs) on the surface, have together limited its further usage as photoanode materials for PEC water oxidation.

In this thesis, solution-processed large-area  $(4.5 \times 4.5 \text{ cm}^2) \text{ Ag}_3\text{PO}_4$  microcrystalline films are first prepared via a room-temperature, air-exposed immersion reaction. The multiple effects of the native defects in Ag<sub>3</sub>PO<sub>4</sub>, *i.e.* V<sub>Ag</sub>, Ag<sub>i</sub>, and surface precipitation of Ag<sup>0</sup> NPs, are first demonstrated experimentally. By changing the addition amount of the oxidation agent H<sub>2</sub>O<sub>2</sub> during the synthetic reaction, Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films with different amounts of Ag<sup>0</sup> NPs on the surface are obtained. It is found that despite the favorable superior conductivity and plasmonic effect having improved the interfacial charge transfer and light absorption of the Ag<sub>3</sub>PO<sub>4</sub> film photoanode, the defective film with the largest amount of Ag<sup>0</sup> NPs on the surface finally exhibits the lowest current density, due to the disadvantageous effects involving the decreased crystalline quality, easier decomposition of active material of Ag<sub>3</sub>PO<sub>4</sub>, decreased electrochemical surface area with a high content of Ag<sup>0</sup> NPs covering the surface of Ag<sub>3</sub>PO<sub>4</sub>, and also serving as the recombination centers. As a result, even the Ag<sub>3</sub>PO<sub>4</sub> microcrystalline film which has achieved the highest photocurrent density fails to demonstrate considerable stability. These results have further highlighted the significance of mitigation of native defects of Ag<sub>3</sub>PO<sub>4</sub> fundamentally towards enhanced water oxidation stability and efficiency of the microcrystalline film photoanodes.

Typically, the post-treatment of vacuum annealing at different temperatures (i.e. 100,

200, 300, 400 and 500 °C) has been investigated as a strategy for solving the problem of native defects and low light absorption of Ag<sub>3</sub>PO<sub>4</sub>. It is found that below 400 °C is the Recovery stage for Ag<sub>3</sub>PO<sub>4</sub> microcrystals when defect amount decreases and therefore crystallinity increases. From 400 °C, the |Recrystallization| stage and also precipitation of Ag<sub>3</sub>PO<sub>4</sub> bulk crystals start, and afterwards till 500 °C is the |Recrystallization + Grain growth| stage for Ag<sub>3</sub>PO<sub>4</sub>, when more faceted and larger bulk crystals of Ag<sub>3</sub>PO<sub>4</sub> precipitate extensively. In this series of Ag<sub>3</sub>PO<sub>4</sub> film photoanodes, the same variation in  $[J_{\text{light}} - J_{\text{dark}}]$  vs. E curves is found as the light absorption profiles, *i.e.* 400 °C > 500 °C >  $300 \text{ }^{\circ}\text{C} > 200 \text{ }^{\circ}\text{C} > \text{non-annealed} > 100 \text{ }^{\circ}\text{C}\text{-annealed } \text{Ag}_{3}\text{PO}_{4} \text{ film, indicating that light}$ absorption, especially the bandgap absorption could be the dominant factor to the PEC performance. Particularly, much improved PEC water oxidation stability and efficiency of the 400 °C-annealed Ag<sub>3</sub>PO<sub>4</sub> film photoanode are demonstrated compared with the non-annealed Ag<sub>3</sub>PO<sub>4</sub> film, suggesting that the methodology of vacuum annealing could be utilized for improving the PEC performance of Ag<sub>3</sub>PO<sub>4</sub> film photoanodes because of its ability to mitigate the unfavorable effects of the native defects like  $V_{Ag}$ , remove the Ag<sup>0</sup> NPs on the surface, and substantially increase the bandgap light absorption of these Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films.

Next, from the viewpoint of advanced structural design, in consideration that the usage of microcrystalline film photoanode might cause the problem of decreased active surface area despite its relatively higher stability than nanocrystalline powder-deposited photoanode, the possibility of introducing porous structures into these microcrystalline Ag<sub>3</sub>PO<sub>4</sub> film photoanodes is also investigated towards further improved photoactivity and stability. Particularly, it is demonstrated that the porous degree of the as-obtained Ag<sub>3</sub>PO<sub>4</sub> films could be tuned by changing the reaction time from 1 h to 2 h, 5 h, 9 h, 20

h and 24 h, with the usage of specific reaction parameters. These porous  $Ag_3PO_4$  films, in spite of having left the native defects not recovered and  $Ag^0$  NPs on the surface not removed, finally could demonstrate improved surface areas and light absorption, owing to which improved PEC performance has also been achieved. Specifically, the optimal highly porous 20-h  $Ag_3PO_4$  microcrystalline film photoanode shows superior water oxidation efficiency with a photocurrent density of 4.32 mA/cm<sup>2</sup> which is almost three times higher than that of non-porous 1-h  $Ag_3PO_4$  microcrystalline film (1.48 mA/cm<sup>2</sup>) at 1.0 V *vs.* Ag/AgCl. This value of photocurrent density is also higher than that of the 400 °C-annealed non-porous  $Ag_3PO_4$  film photoanode.

Overall, it could be concluded that: a) The instability of the Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films is originated from the native defects involving  $V_{Ag}$  and Ag<sup>0</sup> NPs precipitated on the surface, thus the methodology of vacuum annealing, typically at 400 °C, turns out to be more favorable for improving the stability because it can fundamentally solve, or at least mitigate the various problems induced by native defects; b) Light absorption is a crucial factor to the PEC performance, whereas the surface area which would largely determine the final process of PEC water oxidation, *i.e.* the surface redox reaction is also important, thus the porous Ag<sub>3</sub>PO<sub>4</sub> microcrystalline film photoanodes can achieve even higher efficiency in spite of leaving the defect and absorption issues not solved; c) Trying to make a combination of these two techniques, *i.e.* taking both the fundamental properties of the material (*e.g.*, defects and absorption) and the advanced structural design (*e.g.*, larger surface area) into concern, may provide further possibility towards more stable and efficient Ag<sub>3</sub>PO<sub>4</sub> film photoanodes for PEC water oxidation.

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## 1. Introduction

#### **1.1.** Photoelectrochemical water splitting

Since the first report of electrochemical photolysis of water at a TiO<sub>2</sub> electrode under ultraviolet light irradiation,[1] photoelectrochemical (PEC) water splitting has attracted tremendous research interest as a promising technology for sustainable production of clean energy carrier – hydrogen, from solar light and water.<sup>[2]</sup> Generally speaking, one PEC cell should be composed of three major parts, which are the working electrode, the counter electrode and the electrolyte solution, respectively. One common approach is to utilize the specific semiconductor with desirable photoactivity as the working electrode, for efficient production of photo-excited electron-hole pairs (i.e. the photoelectrons and photoholes). In such a traditional system, the oxidation of water (*i.e.* the oxygen evolution reaction, OER) will happen on the surface of the photoanode. Meanwhile, the reduction of water (i.e. the hydrogen evolution reaction, HER) will accordingly happen on the surface of the photocathode with the assistance of external bias and under simultaneous light irradiation.[3] Typically, when the semiconducting photoelectrode has been immersed into the electrolyte solution, the process of charge transfer will be happening across the as-formed semiconductor-electrolyte junction (SEJ) in order to keep the electrostatic equilibrium of the PEC system. Consequently, after such an electrostatic equilibrium has been established, the Fermi level  $(E_{\rm F})$  of the electrons (i.e. photoelectrons) of the semiconductor, as well as the redox potential of the electrolyte solution, will be aligned at the as-formed SEJ interface. This will afterwards result in an upward bending of the electronic band in the *p*-type semiconductor for a *p*-type SEJ, and accordingly a downward bending of the electronic band in the *n*-type semiconductor for an *n*-type SEJ. In fact, this charge transfer process will lead to the removal of majority charge carriers within a characteristic region near the surface of the semiconductor, termed as depletion layer, where a built-in electric field is generated and therefore the "bending of the electronic band" will takes place. Typically in a PEC cell, depletion layer-induced band bending could reversely help the transfer and separation of photogenerated charge carriers. Hence in this regard, it is of great importance to understand such band bending in PEC process to achieve optimized PEC performance. Typically, the band bending phenomenon is usually generated from the surface states of semiconductors as well as the metal–semiconductor contact (*i.e.* ohmic contact). Particularly in PEC process, the applied bias ( $V_{app}$ ) and electrolyte molecules absorption will also lead to the occurrence of band bending.[4]



**Figure 1.1.** Effect of  $V_{app}$  ( $V_A$  in the figure) on an *n*-type semiconductor photoelectrode.  $\Phi_M$  is the work function of the metal.  $\Phi_{SC}$  is the potential drop across the depletion layer. Reproduced with the permission from Ref. 4.

The surface of the semiconducting photoelectrode will absorb the molecules in the electrolyte solution upon contacting each other. The conduction band (CB) of acceptor molecules will consequently shift downwards until equilibrium is established, resulting

in the electron transfer from the semiconductor to absorbed molecules. This will cause the formation of a Helmholtz layer, which is the region between the absorbed ions on the semiconductor surface and closest ions in electrolyte. The potential drop across the Helmholtz layer ( $V_{\rm H}$ ) changes by –59 mV per unit pH at 25 °C.[5] Meanwhile, the band bending could also be induced by  $V_{\rm app}$  between the semiconductor working electrode and reference electrode. The difference in the potential would be distributed over both the depletion layer and Helmholtz layer. Further, since  $V_{\rm H}$  is constant, any change in the  $V_{\rm app}$  will have an influence on the depletion layer. For instance, a positive  $V_{\rm app}$  applied on *n*-type semiconductor will lead to an increase in the depletion layer (*i.e.* upward band bending, Figure 1.1a), whereas the band bending degree will be reduced if a negative  $V_{\rm app}$  has been applied (Figure 1.1b). In contrast, a negative  $V_{\rm app}$  is needed for increasing the depletion layer of a *p*-type semiconductor and a positive  $V_{\rm app}$  is needed for reducing the depletion layer.[4]

If an *n*-type semiconductor is immersed into an electrolyte solution which contains  $H_2O/O_2$  redox couples, as has been discussed above, the charge transfer occurs across the SEJ interface until electrostatic equilibration. Then the semiconductor electrode will have an excess of positive charge distributed over the depletion layer, and the electrolyte will have an excess of negative charge distributed over Helmholtz layer. Further, the light irradiation will bring about large quantities of non-equilibrium photogenerated charge carriers, resulting in a built-in electric field near the semiconductor surface, and thus the generation of the photovoltage (*i.e.* open-circuit voltage,  $V_{OC}$ ). The maximum current generated in this electric field is termed as the short circuit current ( $J_{SC}$ ). Thus the maximum power point  $P_{PA} = V_{OC} \times J_{SC}$ . An overall energy diagram of the SEJ of an *n*-type semiconductor is shown in Figure 1.2.[6] Briefly, given that the doped *n*-type

semiconductor has a uniform distribution of donor ions  $N_D$  that are balanced by the same quantity of free electrons  $n_0 = N_D$  in the neutral region, where the band edge is flat. The doping amount can be indicated as  $\zeta_{nb} = E_{cb} - E_{Fn,b}$  and then we can write  $\zeta_{nb} = k_B T$  $\ln(N_c/n_0)$  where  $N_C$  is the effective density of states in the CB and  $k_B T$  is the thermal energy. The potential in the Helmholtz layer is  $V_H = \phi_s - \phi_{el}$ . The change in  $V_H$  produces the modification of the position of the band edge  $E_{cs}$ . The potential  $V_{sc}$  is related to the height of the barrier in the semiconductor, that is  $V_{sc} = \phi_b - \phi_s$ , where  $\phi_b$  is the electrostatic potential in semiconductor neutral region, and  $\phi_s$  is the potential at the semiconductor surface. The measured voltage can be expressed as  $V_r = (E_0 + \chi + \zeta_{nb})/q + V_{sc} + V_H$ . The CB edge at position x in the semiconductor is modified by the local electrostatic potential  $E_c(x) = -\chi - q[\phi(x) - \phi_{el}]$ . The energy of the CB edge at the surface is  $E_{cs} = -\chi - qV_H$ . The potential of the CB edge  $V_{cs}$  is determined by the expression  $E_{cs} = E_0 - qV_{cs}$ , therefore  $V_{cs} = (E_0 + \chi)/q + V_H$ . Hence, the flatband potential can be formulated as  $V_{fb} = (E_0 + \chi + \zeta_{nb})/q + V_H$ , or  $V_{fb} = V_{cs} - \Delta V_H + \zeta_{nb}/q$ , indicating that flatband potential directly correlates with the position of the CB edge.

In many situations, the applied bias  $V_{app}$  is of central interest for the analysis of PEC performance of a semiconductor electrode. Given that the applied bias  $V_{app}$  with respect to an equilibrium initial situation, only modifies the  $V_H$  and the semiconductor barrier height  $V_{sc}$ , the  $V_{app}$  therefore can be stated as  $V_{app} = -(E_{Fn,b} - E_{redox})/q$ ,  $V_{app} = V_r - V_{r0}$  or  $V_{app} = \phi_b - \phi_{b0}$ . The built-in potential  $V_{bi}$  is an equilibrium quantity  $V_{bi} = \phi_{b0} - \phi_{s0}$ , thus the  $V_{app}$  can be consequently stated as  $V_{app} = V_{sc} - V_{bi} + (V_H - V_{H0})$ . Particularly as in the Figure 1.2, a negative  $V_{app}$  must be applied to the semiconductor in equilibrium with the electrolyte in order to flatten the bands (*i.e.* reducing the depletion layer).



**Figure 1.2.** The energy diagram of SEJ of an *n*-type semiconductor electrode.  $V_r$  is the measured voltage with respect to standard hydrogen electrode (SHE).  $V_{app}$  is the applied bias.  $\phi$  is the local electrostatic potential.  $E_{Fn}$  is the Fermi level of electrons which is a function of the position.  $E_{Fn,b}$  is the Fermi level of electrons at the back contact.  $E_{Fp}$  is the Fermi level of holes.  $E_0$  is the energy of SHE (*i.e.* -4.44 eV).  $E_{redox}$  is the Fermi level of electrolyte species at redox level.  $E_c$  is the energy of CB edge.  $E_v$  is the energy of valence band edge.  $\chi$  is electron affinity.  $\zeta_{nb}$  is the semiconductor barrier height at depletion layer.  $V_{cs}$  is the potential of the CB at semiconductor surface.  $V_{fb}$  is the flatband potential, and  $V_{bi}$  is the built-in potential. Reproduced with the permission from Ref. 6.

For practical usages, a variety of semiconductors have been investigated for use as cost-effective, highly efficient and stable photoanodes and photocathodes for PEC water splitting, as summarized in Figure 1.3.[7] Specifically, TiO<sub>2</sub> has been recognized as the favorable light absorber for water splitting towards both HER and OER. Nevertheless, TiO<sub>2</sub> due to its large bandgap of 3.0–3.2 eV, generally could absorb only the ultraviolet light, which is just a small fraction of sunlight energy (3–5%). Sulfides and selenides, on the other hand, suffer a lot from their serious photo-corrosion problem, despite their considerable visible-light absorption capability. Compared with TiO<sub>2</sub>, tungsten trioxide (WO<sub>3</sub>) possesses a smaller bandgap of 2.4–2.8 eV, making it possible to utilize both the ultraviolet and visible light within the solar spectrum. Nonetheless, WO<sub>3</sub> still displays a low solar energy conversion efficiency because of the relatively low reduction potential for electrons, which determines its quite positive CB position. Besides, hematite ( $Fe_2O_3$ ) is another well-recognized and widely studied photoanode material attributed to its low cost, high stability, along with superior visible-light absorption. However, despite the tremendous investigations which have been conducted so far on this material, there are still a lot of challenges for further employment of  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> towards PEC water oxidation. For example, the flat-band potential of  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> is too positive for reduction of hydrogen, and therefore will need a large overpotential for driving the OER to happen. Moreover, the relatively low absorption coefficient and unsatisfactory conductivity for the majority carrier of  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> are both unfavorable for the achievement of a superior solar energy conversion efficiency. Hence in this regard, it is still quite challenging for researchers to develop novel photoanode materials of higher solar energy conversion efficiency and better stability towards PEC water oxidation.

#### 1.2. Ag<sub>3</sub>PO<sub>4</sub>: Photoactivity and drawbacks for PEC water oxidation

The advances in PEC technology substantially rely on the development of highly efficient and stable photoelectrodes, especially photoanodes, considering that the OER happening on photoanodes is usually more complex and dynamically sluggish than the HER happening on the cathodes, and thus is more dominant in the water splitting process: For one thing, the production of one O<sub>2</sub> molecule requires removal of four electrons and two protons from two H<sub>2</sub>O molecules, which are twice than those for HER; For another, water oxidation is also an uphill reaction with a positive Gibbs free energy from a thermodynamic standpoint.<sup>[8]</sup> Traditional TiO<sub>2</sub>-based photoanodes suffer a lot from low visible-light absorption induced by the large bandgaps. Meanwhile, sulfides and (oxy)nitrides such as CdS, [9] Ta<sub>3</sub>N<sub>5</sub>, [10] and TaON [11] demonstrate poor stability over long-term operation in actual conditions. These have highlighted WO<sub>3</sub>, BiVO<sub>4</sub> and  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> as novel alternatives to the above-mentioned photoanode materials towards efficient and stable water oxidation.[3,12–18] The first report of using silver phosphate (Ag<sub>3</sub>PO<sub>4</sub>) for photocatalytic water oxidation can be traced back to as early as 1988 by Tennakone et al..[19] In 2010, as an important benchmark report, Yi et al. further demonstrated that Ag<sub>3</sub>PO<sub>4</sub> as a photoanode material, with the presence of AgNO<sub>3</sub> in electrolyte as an electron scavenger, could display much higher photocatalytic activity for water oxidation than well-recognized visible-light-driven OER catalysts including BiVO<sub>4</sub> and WO<sub>3</sub>,[20] and this superior performance could be further improved through the strategies like morphology and facet engineering.<sup>[21]</sup>

Till now, the photocatalytic activity of Ag<sub>3</sub>PO<sub>4</sub> for pollutant degradation has been widely investigated and optimized.[22–27] In addition, Ag<sub>3</sub>PO<sub>4</sub> is regarded as an ideal

candidate for water photooxidation due to its low toxicity, moderate bandgap with an absorption edge of ~530 nm, highly positive valence band edge position at ~2.85 V vs. reversible hydrogen electrode (RHE), which can afford large thermodynamic driving force for photoholes to oxidize water.[28–29]



**Figure 1.3.** Schematic illustration of the Ag<sub>3</sub>PO<sub>4</sub> structure, displaying both [AgO<sub>4</sub>] and [PO<sub>4</sub>] tetrahedral units. Ag<sub>3</sub>PO<sub>4</sub> is of BCC structure with the typical lattice parameter of ~6.004 Å, consisting of isolated [PO<sub>4</sub>] tetrahedra with the P–O distance of ~1.539 Å, as well as every six Ag<sup>+</sup> distributing among twelve sites of the two-fold symmetry. In the modeled structure, each isolated [PO<sub>4</sub>] unit is bonded to three neighboring [AgO<sub>4</sub>] units by O atoms, and the existence of two O–Ag–O bond angles ( $\alpha$ ,  $\beta$ ) indicates the [AgO<sub>4</sub>] units are highly distorted in the lattice, which is resulted from the inductive effect of the high-electronegativity [PO<sub>4</sub>] tetrahedra. Reproduced with the permission from Ref. 30.

In general, the efficiency of PEC water oxidation on the photoanode surface is determined by four fundamental processes: 1) Light absorption (*i.e.* photoexcitation of

charge carriers); 2) Separation of photoexcited electron–hole pairs; 3) Charge migration towards the SEJ interface; 4) Surface redox reaction with water molecules. Particularly for Ag<sub>3</sub>PO<sub>4</sub>, the fairly higher intrinsic mobility of photoelectrons than photoholes, and the inductive effect of [PO<sub>4</sub>] tetrahedron units originated from high electronegativity together facilitate the separation of electron–hole pairs, which provides solid support for its excellent photocatalytic activity.[29] Figure 1.3 presents a schematic illustration of the crystal structure of Ag<sub>3</sub>PO<sub>4</sub>, where such a unique inductive effect of the [PO<sub>4</sub>] units has been discussed.

However, the large numbers of intrinsic point defects [i.e. silver vacancy ( $V_{Ag}$ ) and silver interstitial  $(Ag_i)$ ] resulted from the high ionic mobility of  $Ag^+$  in silver compounds like Ag<sub>3</sub>PO<sub>4</sub> and silver halides will easily trap photogenerated electrons, [30–32] leading to non-negligible precipitation of metallic Ag<sup>0</sup> on the surface. It has been revealed both theoretically and experimentally that such  $V_{Ag}$  and  $Ag_i$  can exist from room-temperature growth of Ag<sub>3</sub>PO<sub>4</sub>, and are highly mobile, which is determined by their low formation energy and migration barrier.[33-35] Although this characteristic could be used in some special occasions requiring tunable Ag<sup>+</sup> release for bactericidal activity, [36] it is not good for the photocatalytic or PEC processes because the decomposition of Ag<sub>3</sub>PO<sub>4</sub> means the loss of the photoactive semiconductor material, making it unstable during long-time reactions. In addition, the real conditions in PEC water oxidation process (i.e. aqueous solution and light irradiation) make it even easier for Ag<sub>3</sub>PO<sub>4</sub> to be reduced to Ag<sup>0</sup> due to the high electrode potential of Ag-Ag<sub>3</sub>PO<sub>4</sub> (0.45 V vs. RHE).[37] Meanwhile, the precipitated  $Ag^0$  sites on the surface, due to their better electron affinity than Ag<sub>3</sub>PO<sub>4</sub>,[38–39] might easily attract and store electrons, and thereafter become recombination centers for photoholes, lowering the photocurrent density as a result. To

ease this problem, Wu *et al.* incorporated a  $Ag^0$ -nanoplate underlayer in their  $Ag_3PO_4$  photoanode to drain the electrons away to the cathode, as shown in Figure 1.4 below.[40] Nevertheless, this couldn't be a fundamental solution since the  $Ag^0$  nanoparticles (NPs) will be continuously precipitated at the  $Ag_3PO_4$ /electrolyte interface during long-term OER operation, offsetting such an advantageous draining effect.



**Figure 1.4.** Schematic illustrations showing (a) the energy band diagram, as well as (b) the charge-flow direction (*i.e.* draining of photoelectrons) of the Ag–Ag<sub>3</sub>PO<sub>4</sub> composite photoanode for PEC water oxidation. Reproduced with the permission from Ref. 40.

Different types of photoelectrodes, at the same time, could have a great influence on the performance, especially the stability. Microcrystalline film photoelectrodes usually demonstrate better structural stability than nanocrystalline powder-deposited electrodes due to the superior anti-aggregation ability and stronger binding force between active material layers and substrates. More importantly, the increased crystallite size leads to decreased amount of surface defects and dangling bonds, which will easily become self-degradation sites of the active semiconductor materials.[41] Nonetheless, due to the challenges of the poor stability and the difficulties in fabricating  $Ag_3PO_4$  membranes, it is considered to be not easy to obtain highly adhesive  $Ag_3PO_4$  layers on conductive substrates.[42–44] Recently, Cao *et al.* reported direct growth of highly porous  $Ag_3PO_4$ microcrystalline film photoanodes on as-purchased silver foils with large surface areas, which consequently demonstrated enhanced OER stability and efficiency.[45] However, the technical problem of native defects-induced precipitation of  $Ag^0$  NPs on the surface of  $Ag_3PO_4$  was still not solved.

Besides, as another fundamental process determining the PEC conversion efficiency, light absorption of as-reported Ag<sub>3</sub>PO<sub>4</sub> photocatalysts is far from satisfactory. Despite its high photoactivity, the bandgap absorbance in the region of  $\leq -530$  nm of Ag<sub>3</sub>PO<sub>4</sub> powder photocatalyst (fabricated into thick film for measurement) can hardly exceed 0.7 (*i.e.*  $\leq 80\%$  absorption of incident light),[34,46–47] which is also much lower than other traditional semiconductors such as BiVO<sub>4</sub>,[26] WO<sub>3</sub>,[48] and Co<sub>3</sub>O<sub>4</sub>.[49] Specifically for solution-processed Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films, although the advantageous highly porous surface structure could help to increase light absorption greatly, the maximum absorption is still below 80% across the ultraviolet–visible region.[45] Overall, despite the arguably highest photoactivity of Ag<sub>3</sub>PO<sub>4</sub> among common semiconductors, there are still two big issues, *i.e.* native defects-induced Ag<sup>0</sup> precipitation, and poor bandgap light absorption, which need to be solved for further improvement towards highly stable and efficient photoanodes for PEC water oxidation.

## **1.3.** Experimental demonstration of the multiple effects of $Ag^0$ NPs

Now that the easy decomposition of  $Ag_3PO_4$  and consequently precipitation of  $Ag^0$ NPs on the surface is a well-recognized but unavoidable problem originated from the intrinsically high ionic mobility of  $Ag^+$  and native defects of  $V_{Ag}$  and  $Ag_i$  in  $Ag_3PO_4$ , as discussed above, it can be meaningful if a more comprehensive understanding could be achieved experimentally, regarding the multiple effects of surface precipitation of  $Ag^0$ NPs on the various physicochemical properties and PEC water oxidation performance of  $Ag_3PO_4$  microcrystalline film photoanodes.



**Figure 1.5.** SEM images of the  $Ag_3PO_4$  microcrystalline films with different amounts of surface precipitation of  $Ag^0$  NPs: (a–b) very little amount, (c–d) moderate amount, (e–f) large amount, respectively.



**Figure 1.6.** Representative XRD patterns of the Ag<sub>3</sub>PO<sub>4</sub> films with different amounts of Ag<sup>0</sup> NPs on the surface. Panel (b) is an enlarged view of the region 37.5 °  $\leq 2\theta \leq 45.0$  °, showing the characteristic diffraction peaks of metallic silver phase.



**Figure 1.7.** Representative Raman spectra of the  $Ag_3PO_4$  films with different amounts of surface precipitation of  $Ag^0$  NPs. The different singly, doubly, and triply degenerate vibrational modes are labelled by  $A_1$ , E, and T<sub>2</sub>, respectively.



**Figure 1.8.** Representative high-resolution XPS spectra of the (a) Ag 3d region, and (b) Ag MVV Auger electron region, of the  $Ag_3PO_4$  films with different amounts of surface precipitation of metallic  $Ag^0$  NPs.

Typically, different amounts of metallic  $Ag^0$  precipitation on the surface of  $Ag_3PO_4$ microcrystalline films are intentionally induced by applying different addition amounts of the oxidation agent  $H_2O_2$  during the synthesis experiment. As is shown in Figure 1.5, SEM images demonstrate the different amounts of  $Ag^0$  NPs precipitated on the surface of as-obtained  $Ag_3PO_4$  microcrystalline films, termed as the "very little amount" (Figure 1.5a–b), the "moderate amount" (Figure 1.5c–d), and the "large amount" (Figure 1.5e–f), respectively. The increase in the amount of surface precipitated  $Ag^0$  NPs from the "very



**Figure 1.9.** EIS Nyquist plots of the Ag<sub>3</sub>PO<sub>4</sub> film photoanodes with different amounts of surface precipitation of metallic Ag<sup>0</sup> NPs. The insert displays the equivalent circuit consisting of a serial resistor ( $R_S$ ), a resistor–capacitor circuit modelling the parallel combination of the interfacial charge transfer resistance ( $R_{CT}$ ), and constant phase element (CPE) attributed to the SEJ interface.



**Figure 1.10.** Representative diffuse-reflectance spectra (DRS) of the  $Ag_3PO_4$  films with different amounts of surface precipitation of metallic  $Ag^0$  NPs.



**Figure 1.11.** PEC water oxidation performance of the Ag<sub>3</sub>PO<sub>4</sub> film photoanodes with different amounts of Ag<sup>0</sup> NPs precipitated on the surface: (a) the dependence of current density on applied potential. The solid and dashed lines stand for the photocurrent and dark current densities of each photoanode, respectively. Inset shows an enlarged view of the potential range from 0.22 to 0.46 V *vs*. Ag/AgCl. (b) The current density–time (J–t) curve of the photoanode fabricated from the Ag<sub>3</sub>PO<sub>4</sub> film with the "moderate amount" of Ag<sup>0</sup> NPs on the surface during the stability test at an applied potential of 1.0 V *vs*. Ag/AgCl.



**Figure 1.12.** The dependence of photocurrent density minus dark current density ( $J_{\text{light}} - J_{\text{dark}}$ ) on applied potential of the Ag<sub>3</sub>PO<sub>4</sub> film photoanodes with different amounts of Ag<sup>0</sup> NPs precipitated on the surface. Inset shows an enlarged view of the potential range from 0.22 to 0.46 V *vs.* Ag/AgCl.



**Figure 1.13.** SEM images exhibiting the surface morphology after the stability test of the photoanode fabricated from the  $Ag_3PO_4$  film with "moderate amount" of  $Ag^0$  NPs precipitated on the surface: (a) overall morphology; (b) the interfacial area of the eroded and uneroded part; (c–d) typical morphology of an uneroded (c) and eroded (d) area.

little amount" to the "large amount" Ag<sub>3</sub>PO<sub>4</sub> film has also been verified by the increased diffraction intensity of the characteristic peaks of metallic Ag<sup>0</sup> phase in XRD patterns (Figure 1.6), as well as the slightly positive shift of peak positions in the XPS narrow scans of the Ag 3d region (Figure 1.8a). Specifically, because the 3d peaks for Ag<sup>+</sup> are at 367.8 and 373.8 eV respectively, while the corresponding peak centers for Ag<sup>0</sup> are at 369.2 and 375.2 eV, such a gradually positive shift could help to confirm the increase in the amount of Ag<sup>0</sup> NPs on the surface of the Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films.[47,50–52] The similar phenomenon is also observed in the Ag MVV Auger electron region (Figure 1.8b).[53–54] Meanwhile, it is clear that as the amount of the surface precipitated Ag<sup>0</sup> NPs gets increased, the surface roughness of the Ag<sub>3</sub>PO<sub>4</sub> microcrystalline film will also increase (Figure 1.5), accompanied by the decrease in the crystallinity (*i.e.* increase in the amount of native defects involving  $V_{Ag}$  and Ag<sub>i</sub>) which are revealed by the decreased intensities of the characteristic diffraction peaks of Ag<sub>3</sub>PO<sub>4</sub> in the XRD patterns (Figure 1.6), and also the Raman scattering bands in as-recorded Raman spectra (Figure 1.7).

It should be noticed that despite the fact that the  $Ag^0$  NPs on the surface of  $Ag_3PO_4$ may become recombination centers and also decrease the electrochemical surface areas (ECSAs) of the microcrystalline film photoanode, which are both unfavorable for PEC water oxidation performance, these  $Ag^0$  NPs could be advantageous as well from other aspects. For one thing, the increased amount of metallic  $Ag^0$  NPs precipitated on the surface could help to promote the interfacial conductivity, and therefore the interfacial charge transfer rate of the  $Ag_3PO_4$  film photoanodes, as witnessed by the EIS spectra in Figure 1.9. For another thing, these  $Ag^0$  NPs could also help to achieve enhanced light absorption of the  $Ag_3PO_4$  photoanode due to their favorable plasmonic effect.[43,55] It could be seen clearly in Figure 1.10 that as the amount of  $Ag^0$  NPs on the surface of the

Ag<sub>3</sub>PO<sub>4</sub> microcrystalline film photoanode increases, light absorption get improved in the region of  $\leq -600$  nm. Indeed, this is favorable for improving the PEC performance. However, the maximum light absorbance across the whole ultraviolet-visible light band of  $\sim 0.7$  which is achieved by the Ag<sub>3</sub>PO<sub>4</sub> microcrystalline film having "large amount" of Ag<sup>0</sup> NPs precipitated on the surface, is still far from satisfactory, as has already been discussed in the above section. Moreover, this sample (*i.e.* "large amount" one) and the accordingly fabricated Ag<sub>3</sub>PO<sub>4</sub> film photoanode finally shows the lowest photocurrent density (Figure 1.11a and Figure 1.12) and positively shifted onset potential (the inset in Figure 1.11a and Figure 1.12), which should be mainly resulted from the increase in the defect amount.[3,56-57] Besides, it is also noticeable that the microcrystalline Ag<sub>3</sub>PO<sub>4</sub> film photoanode with the "moderate amount" of Ag<sup>0</sup> NPs on the surface, despite having achieved the highest photocurrent density among the three different samples, as shown in Figure 1.11a and Figure 1.12, has still failed to demonstrate superior stability due to the non-negligible defect-induced decomposition of the active material of Ag<sub>3</sub>PO<sub>4</sub>, and thus precipitation of metallic  $Ag^0$  NP aggregations (*i.e.* the while aggregations in Figure 1.13c), as can be known from Figure 1.11b and Figure 1.13. Hence in this regard, there is still much work needed to be done in order to fundamentally reduce the native defects of Ag<sub>3</sub>PO<sub>4</sub> towards improved PEC water oxidation stability and efficiency.

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## 2. The effects of vacuum annealing boosting light absorption and water oxidation stability of novel microcrystalline Ag<sub>3</sub>PO<sub>4</sub> film photoanodes

#### 2.1. Supplementary introduction

It should be noted that the generation of native defects, *i.e.* the  $V_{Ag}$  and simultaneous precipitated metallic Ag<sup>0</sup> from Ag<sub>3</sub>PO<sub>4</sub> photocatalysts, plays a completely different role in the case of photocatalytic dye degradation than PEC water oxidation using Ag<sub>3</sub>PO<sub>4</sub> photoanodes. Typically in photocatalytic dye degradation reactions using Ag<sub>3</sub>PO<sub>4</sub> as the photocatalyst, the synergistic effect of  $V_{Ag}$  and metallic Ag<sup>0</sup> precipitation may contribute to efficient separation of photogenerated electron–hole pairs, and consequently better photocatalytic performance. Unlike the multiple negative effects on the water oxidation performance of Ag<sub>3</sub>PO<sub>4</sub> photoanodes, such opposite positive effects may have provided a promising strategy towards enhanced performance of Ag<sub>3</sub>PO<sub>4</sub> photocatalysts. Hence in this regard, *in-situ* chemical etching using ethanol during the synthesis of Ag<sub>3</sub>PO<sub>4</sub>,[1] or post-etching using ammonia solution,[2] as well as post-annealing in air at different temperature and time conditions of the as-obtained Ag<sub>3</sub>PO<sub>4</sub> powder photocatalyst,[3–7] have been exploited till now for intentionally introducing the  $V_{Ag}$  and metallic Ag<sup>0</sup> NPs on the surface, in order to achieve higher photocatalytic efficiencies and stability.

Light absorption is primarily determined by the bandgap, which in turn could also be modified by morphology, crystalline quality (*i.e.* crystallinity) and crystal structure, *etc.* In view of this, in consideration that annealing is a common method used to improve the crystallinity and thus light absorption of semiconductor materials, [8–9] whereas it may also cause undesired precipitation of  $Ag^0$  NPs on the surface of  $Ag_3PO_4$  photoanodes, as has been discussed above in the Introduction section, it seems to be a double-edged sword for improving the PEC water oxidation performance of  $Ag_3PO_4$  photoanodes. Herein this chapter, instead of the traditional air annealing methodology, we investigate the novel vacuum annealing on the modification of the crystallinity and properties of microcrystalline  $Ag_3PO_4$  film photoanodes. It is amazingly found that such a vacuum annealing technique could help reduce the native defects, *i.e.* the  $V_{Ag}$  and its induced precipitation of metallic  $Ag^0$  NPs on the surface, of the obtained  $Ag_3PO_4$  films, affording substantially promoted bandgap light absorbance to as high as 0.9, and hence significantly improved PEC water oxidation stability and efficiency. Furthermore, the different mechanism of how the vacuum annealing process reduces the native defects of  $Ag_3PO_4$  microcrystalline films, with a comparison to traditional air annealing process, has also been discussed.

#### 2.2. Experimental details

Synthesis and annealing of  $Ag_3PO_4$  microcrystalline films. The as-purchased silver foils (99.98%, 0.1 mm thick, Nilaco) were first cut into pieces (4.5 × 4.5 cm<sup>2</sup>) and cleaned using dilute nitric acid solution (Sigma-Aldrich) and acetone ( $\geq$  99.5%, Wako), respectively. Briefly, an aqueous solution was first prepared by mixing 7.2 g of sodium phosphate monobasic (NaH<sub>2</sub>PO<sub>4</sub>,  $\geq$  99.0%, Sigma-Aldrich), 482.8 g of hydrogen peroxide (H<sub>2</sub>O<sub>2</sub>, 30.0–35.5%, Sigma-Aldrich), and 10 g of polyvinylpyrrolidone (PVP,  $M_W \approx 55,000$ , Sigma-Aldrich) under stirring at room temperature till the solution
became transparent. A piece of Ag foil was then immersed into the mixture solution and kept for 90 min. Afterwards, the foil was taken out of the solution and rinsed with ethanol ( $\geq$  99.5%, Wako) and distilled water (resistivity  $\geq$  18 M $\Omega$ ·cm) thoroughly, and dried in air naturally. Next, the obtained Ag<sub>3</sub>PO<sub>4</sub> film was cut into smaller pieces (1.5 × 1.5 cm<sup>2</sup>), loaded onto an alumina boat at the center of a quartz tube, and annealed in vacuum at 100 °C, 200 °C, 300 °C, 400 °C, and 500 °C, respectively for 2 h. Before annealing, the quartz tube was pumped using a mechanical pump for  $\geq$  20 min. During annealing, the ramp rate was kept at 1 °C/min for all samples, and after annealing, all the resulted Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films were cooled down naturally to room temperature for further usage.

*Characterization.* Scanning electron microscope (SEM) images were recorded using two different field-emission SEMs (FE-SEM, S4700 and S4800, Hitachi, Japan). X-ray diffraction (XRD) data were acquired on a Rigaku Miniflex II-MW (Japan) X-ray diffractometer working with Cu-K<sub> $\alpha$ </sub> radiation (1.5406 Å). X-ray photoelectron spectroscopy (XPS) analysis was carried out using a PHI 5000 VersaProbe spectrometer (ULVAC-PHI, Japan) with an Al-K<sub> $\alpha$ </sub> X-ray source (1486.6 eV). Raman spectra were collected on a laser confocal Raman micro-spectrometer (in*Via* reflex, Renishaw, UK) equipped with an integrated optical microscope (DM-2500, Leica, Germany). The excitation light source of a 632.8 nm He–Ne laser, laser power of 1.7 mW (10% × 17 mW), grating of 1800 1/mm, exposure time of 5.0 s, and accumulation of three times were applied for all Raman measurements. UV–visible diffuse-reflectance spectra (DRS) were measured on a JASCO V-670 (Japan) spectrophotometer equipped with an integrating sphere setup. Dark-field scattering spectra as well as microscopic images were taken by an Olympus BX61 industrial optical microscope (Japan) equipped with

the dark-field attachment along with Olympus MPLFLN-BD and LMPLFLN-BD objective lenses. This dark-field microscope is also connected with an IsoPlane-160 (Princeton Instruments, United States) spectrometer.

*PEC and electrochemical measurements.* For evaluation of the PEC water oxidation performance of the Ag<sub>3</sub>PO<sub>4</sub> films, a three-electrode cell was assembled and used for all measurements. The Ag<sub>3</sub>PO<sub>4</sub> photoelectrode, a Pt wire, an Ag/AgCl electrode, and the aqueous solution of Na<sub>3</sub>PO<sub>4</sub> (0.02 M, pH  $\approx$  11.3) were utilized as the working electrode, counter electrode, reference electrode and electrolyte solution, respectively. Potential applied on working electrode was controlled by a VersaSTAT 4 potentiostat (Princeton Applied Research, United States). Before every measurement, the electrolyte solution was purged with high-purity nitrogen gas for  $\geq$  20 min under continuous magnetic stirring. During the measurements, the Ag<sub>3</sub>PO<sub>4</sub> photoanodes were illuminated with a 500 W xenon lamp delivering an irradiance of 285 mW/cm<sup>2</sup>. The scan rate was 10 mV/s from negative to positive potential for recording each current density–potential (*J*–*E*) curve. The electrochemical impedance spectroscopy (EIS) data were acquired under dark conditions at an applied potential of 1.0 V vs. Ag/AgCl varying the frequency between 100,000 (*i.e.* start frequency) and 500 Hz (*i.e.* end frequency) of a sinusoidal perturbation with 10 mV amplitude.

### 2.3. Results and discussion

Large-area ( $4.5 \times 4.5 \text{ cm}^2$ ) microcrystalline Ag<sub>3</sub>PO<sub>4</sub> films were first grown on as-purchased silver foils *via* a facile room-temperature and air-exposed solution process. This solution contained NaH<sub>2</sub>PO<sub>4</sub> as the PO<sub>4</sub><sup>3+</sup> sources, and H<sub>2</sub>O<sub>2</sub> as the initial agent for oxidation from  $Ag^0$  to  $Ag^+$ , where the overall reaction could be described as follows:

$$2 \operatorname{NaH_2PO_4} + 2 \operatorname{H_2O_2} + 6 \operatorname{Ag} \rightarrow 2 \operatorname{Ag_3PO_4} + 2 \operatorname{NaOH} + 2 \operatorname{H_2O} + \operatorname{H_2\uparrow}$$

The as-obtained Ag<sub>3</sub>PO<sub>4</sub> film, as shown in the SEM image in Figure 2.3a, turned out to be microcrystalline with irregular and faceted microcrystals on the surface, indicating that large crystals of Ag<sub>3</sub>PO<sub>4</sub> could be generated from the reaction of PO<sub>4</sub><sup>3-</sup> and Ag<sup>+</sup> at room temperature. The color of as-obtained microcrystalline Ag<sub>3</sub>PO<sub>4</sub> film was dark khaki, as shown in Figure 2.1. After vacuum annealing for 2 h at different temperatures, the color of the Ag<sub>3</sub>PO<sub>4</sub> film changed gradually. From Figure 2.1, it is clear that when higher annealing temperature was applied, the color of resulted Ag<sub>3</sub>PO<sub>4</sub> film would be brighter. The Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films annealed at  $\geq$  300 °C, especially the 400 °C- and 500 °C-annealed films, display an obviously changed color appearance to even bright canary yellow. This drastic color change suggests that the optical properties of the 400 °C- and 500 °C-annealed Ag<sub>3</sub>PO<sub>4</sub> films may have also been changed, and could help to exclude the possibility of the existence of a high content of Ag<sup>0</sup> or Ag<sub>2</sub>O NPs on the surface of the Ag<sub>3</sub>PO<sub>4</sub> film after undergoing a vacuum annealing process, as these may result in darker color (*e.g.*, green yellow or dark khaki) of the films just like the non-annealed one.[2,5]

The decreased amount of  $Ag^0$  NPs on the surface of  $Ag_3PO_4$  is directly confirmed by the dark-field microscopic analysis, as displayed in Figure 2.2. Dark-field imaging technique is widely used for probing  $Ag^0$  nanostructures due to its sensitive response to the surface plasmon resonance (SPR) of silver,[10–11] whereas it is difficult to probe  $Ag^0$  NPs accurately by SEM because many silver-based compounds can be destroyed, or reduced to give non-original  $Ag^0$  NPs by high-energy electron beams during measurements.[12–13] It is clear that in non-annealed Ag<sub>3</sub>PO<sub>4</sub> microcrystalline film (Figure 2.2a), large amount of Ag<sup>0</sup> NPs (*i.e.* brighter white dots) were found to exist originally. This is because a part of Ag<sup>+</sup> in as-formed Ag<sub>3</sub>PO<sub>4</sub> film may be reduced by the high-concentration electrons induced by the violet redox reaction, in consideration of the higher ionic mobility of Ag<sup>+</sup> than PO<sub>4</sub><sup>3–</sup>, as well as the high mobility of electrons in these Ag<sub>3</sub>PO<sub>4</sub> materials. In addition, the easy formation of intrinsic defects at room temperature during the solution growth of Ag<sub>3</sub>PO<sub>4</sub> might have also contributed to this phenomenon because Ag<sup>0</sup> NPs would generated when interstitial Ag<sub>i</sub> was reduced by electrons. It is amazing that after annealing the Ag<sub>3</sub>PO<sub>4</sub> films in vacuum, as shown in Figure 2.2b–f, the amount of Ag<sup>0</sup> NPs on the surface decreases significantly. There are hardly any distinct Ag<sup>0</sup> aggregations that could be observed in the films annealed at  $\geq$  300 °C (Figure 2.2d–f), which is in good consistency with the change in color of the annealed Ag<sub>3</sub>PO<sub>4</sub> films.



**Figure 2.1.** Optical images displaying the obvious color change from the non-annealed  $Ag_3PO_4$  film (1.5 × 1.5 cm<sup>2</sup>) to the films annealed at different temperatures.



**Figure 2.2.** Dark-field microscopic images of the (a) non-annealed  $Ag_3PO_4$  film, and the  $Ag_3PO_4$  films annealed in vacuum at the temperature of (b) 100, (c) 200, (d) 300, (e) 400, and (f) 500 °C, respectively.

This unusual result of vacuum annealing that the amount of metallic  $Ag^0$  NPs on the surface of  $Ag_3PO_4$  got reduced as the annealing temperature was increased might be attributed to its different circumstance compared to air annealing. Regarding the traditional air annealing processes of  $Ag_3PO_4$ ,[3–7] although most reports didn't give detailed explanation for the generation of more  $Ag^0$  NPs after annealing, some researchers proposed that the metallic  $Ag^0$  NPs might come from the reaction between  $Ag^+$  and thermally excited electrons during annealing.[4,7] However, this cannot help to



**Figure 2.3.** SEM images recorded at different magnifications of the (a) non-annealed  $Ag_3PO_4$  film, and  $Ag_3PO_4$  films annealed in vacuum at the temperature of (b) 100, (c) 200, (d) 300, (e) 400, and (f) 500 °C, respectively. For each sample, panels (1) and (2) display the low- and high-magnification SEM images.

understand the reduce in the amount of  $Ag^0$  NPs on the surface when annealing  $Ag_3PO_4$ in vacuum. Therefore, according to our experimental results, it is speculated that the precipitation of more  $Ag^0$  NPs during air annealing is actually originated from the thermal decomposition of intermediate products. Typically, if  $Ag_3PO_4$  is annealed in air

at a high temperature, due to the existence of reactive species in atmosphere involving the H<sub>2</sub>O, CO<sub>2</sub> and O<sub>2</sub> molecules, unstable intermediates with low decomposition temperatures, such as AgOH, Ag<sub>2</sub>CO<sub>3</sub>, AgO and Ag<sub>2</sub>O can easily form – after all, Ag<sub>3</sub>PO<sub>4</sub> itself tends to be decomposed in a humid environment due to its poor structural stability.[6] For instance, the decomposition of Ag<sub>2</sub>CO<sub>3</sub> into Ag<sup>0</sup> NPs, AgO and Ag<sub>2</sub>O usually occurs at ~170 °C; the transformation of AgO to Ag<sub>2</sub>O occurs in 100-200 °C; the complete thermal decomposition of Ag<sub>2</sub>O into Ag occurs at 400 °C.[14–17] In this case, large quantities of Ag<sup>0</sup> NPs could be precipitated on the surface of Ag<sub>3</sub>PO<sub>4</sub> films after annealing in air at  $\geq 400$  °C, due to the complete decomposition of intermediates. In contrast to air annealing, if annealing Ag<sub>3</sub>PO<sub>4</sub> in vacuum, no such intermediates could form to generate Ag<sup>0</sup> NPs due to the absence of reactive gas molecules. Moreover, the originally existed Ag<sup>0</sup> NPs from room-temperature growth of the films may be melted, and thus disappear during high-temperature vacuum annealing process. Previous reports have demonstrated that the melting temperature of metallic nanostructures, including Ag<sup>0</sup> NPs depends greatly on the size.[18] The melting point of  $Ag^0$  NPs, which was calculated to be ~480 °C using a classic thermodynamic description, could be lowered to ~400 °C at a decreased size, and can be further decreased to ~100 °C via the interaction with an ionic substrate.[19-21] Therefore, in consideration of the well-recognized superionic Ag<sup>+</sup> conductive properties of Ag<sub>3</sub>PO<sub>4</sub>,[22–23] it is rational that the melting point of Ag<sup>0</sup> NPs supported on Ag<sub>3</sub>PO<sub>4</sub> surfaces could be decreased to quite low values, resulting in the reduce of the amount of Ag<sup>0</sup> NPs during annealing in vacuum at elevated temperatures. The melted Ag<sup>0</sup> NPs may become chemisorbed Ag<sup>0</sup> atoms, enter the defective lattices of Ag<sub>3</sub>PO<sub>4</sub> via diffusion, and finally recover the intrinsic  $V_{Ag}$  defects, facilitating for higher crystalline

quality of the annealed Ag<sub>3</sub>PO<sub>4</sub> films.



**Figure 2.4.** Low-magnification SEM images of the (a) non-annealed  $Ag_3PO_4$  film, and the  $Ag_3PO_4$  films annealed in vacuum at the temperature of (b) 100, (c) 200, (d) 300, (e) 400, and (f) 500 °C, respectively.

The morphological evolution of the  $Ag_3PO_4$  microcrystalline films annealed at different temperatures is revealed by SEM images shown in Figure 2.3 and Figure 2.4 (*i.e.* larger-scale images). It is clear that the surface roughness of the  $Ag_3PO_4$  microcrystalline films increases gradually after annealing. Specifically, the surface roughness increases slightly from non-annealed to 200 °C-annealed  $Ag_3PO_4$  microcrystalline films (Figure 2.3a–c), and then increases significantly from 300 °C- to 500 °C-annealed films (Figure 2.3d–f). This should be resulted from the crystal growth

of Ag<sub>3</sub>PO<sub>4</sub> during annealing. It is notable that in the films annealed at 400 °C, some faceted bulk crystals could be observed (Figure 2.3e and Figure 2.5a), and when the annealing temperature further goes to 500 °C, as shown in Figure 2.3f and Figure 2.5b, more bulk crystals of Ag<sub>3</sub>PO<sub>4</sub> with larger size and more faceted surfaces can be identified. The non-negligible large amount of Ag<sub>3</sub>PO<sub>4</sub> bulk crystals could be observed more obviously when the inside of the 500 °C-annealed film has been exposed and shown in Figure 2.6.



**Figure 2.5.** Enlarged SEM images and detailed analysis of the precipitated  $2^{nd}$ -phase bulk crystals of Ag<sub>3</sub>PO<sub>4</sub> in the resulted films annealed at (a) 400 °C and (b) 500 °C. Panels (a) and (b) correspond to the panels (e<sub>2</sub>) and (f<sub>2</sub>) of Figure 2.3, respectively.

Based on these results, it could be deduced that the annealing process of the Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films could be divided into 3 stages depending on different annealing temperatures: 1) From room temperature (*i.e.* non-annealed) to 300 °C is the |Recovery| stage, in which the amount of intrinsic defects  $V_{Ag}$  and Ag<sub>i</sub> decreases, and the crystallinity increases; 2) From 300 °C to 400 °C, likely more close to 400 °C, the |Recrystallization| stage starts, when the precipitation of Ag<sub>3</sub>PO<sub>4</sub> bulk starts; 3) From 400 °C to 500 °C and higher temperature is the |Recrystallization + Grain growth| stage, when more and larger bulk crystals of Ag<sub>3</sub>PO<sub>4</sub> are precipitated.



**Figure 2.6.** SEM images recorded at different magnifications showing the inside of the 500 °C-annealed  $Ag_3PO_4$  film. This sample was prepared by tearing up the  $Ag_3PO_4$  film using a conductive tape to make the inside exposed. Red arrows in panel (b) indicate the resulted potholes when the second-phase bulk crystals of  $Ag_3PO_4$  were moved away.

To verify the component and phase purity of the non-annealed and annealed Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films, two X-ray techniques, XRD and XPS have been exploited for analyzing the films, as shown in Figure 2.7. Specifically in the XPS spectra exhibited in Figure 2.7b, for all the non-annealed and annealed microcrystalline film samples, the

characteristic peaks of Ag 3d<sub>3/2</sub>, Ag 3d<sub>5/2</sub>, P 2p and O 1s are identified to be at 373.2, 367.2, 131.9 and 529.9 eV of the binding energy, which show perfect matches with the literature values for Ag<sup>+</sup>, as well as P<sup>5+</sup> and O<sup>2-</sup> in PO<sub>4</sub><sup>3-</sup> units in Ag<sub>3</sub>PO<sub>4</sub>.[24–25] Meanwhile, in XRD patterns of all the non-annealed and annealed films (Figure 2.7a), all distinct diffraction peaks can be well indexed to the cubic P(-4)3n phase of silver orthophosphate (Ag<sub>3</sub>PO<sub>4</sub>, PDF#70-0702, labelled in blue in Figure 2.7a). Both these results confirm that the component of the films (*i.e.* pure-phase Ag<sub>3</sub>PO<sub>4</sub>) was retained essentially after annealing in vacuum at a relatively high temperature ( $\geq 400$  °C). The relative crystallinity of the Ag<sub>3</sub>PO<sub>4</sub> films annealed in vacuum at different temperatures, as presented in Table 2.1, is evaluated from the relative intensity of the strongest (210) diffraction peaks of all samples, using that of non-annealed Ag<sub>3</sub>PO<sub>4</sub> film as the reference.[6] It is clear that the Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films annealed at elevated



**Figure 2.7.** Representative (a) XRD patterns and (b) XPS spectra of the as-obtained microcrystalline  $Ag_3PO_4$  films annealed at different temperatures. The abbreviation NA stands for the non-annealed  $Ag_3PO_4$  film.

**Table 2.1.** Relative intensities of three characteristic diffraction peaks in the measured XRD patterns (Figure 2.7a) of all the obtained  $Ag_3PO_4$  microcrystalline films, using the intensities of non-annealed (NA)  $Ag_3PO_4$  as the reference. The relative intensity of the strongest peak, *i.e.* the (210) crystal plane, is used to indicate the relative crystallinity of all the obtained  $Ag_3PO_4$  films.

Temperature	NA	100 °C	200 °C	300 °C	400 °C	500 °C
(2 1 0)	1.00	1.14	1.42	1.78	2.05	2.35
(1 1 0)	1.00	0.93	0.76	1.51	1.28	0.98
(2 2 2)	1.00	1.92	2.10	1.85	3.28	2.20

temperatures show higher crystallinity, confirming that the intrinsic defects of  $V_{Ag}$  and  $Ag_i$  could be recovered to afford better crystal quality of  $Ag_3PO_4$  after annealing.

Besides, it should be noted that in the film annealed at 500 °C, a second cubic P(-4)3n phase of Ag<sub>3</sub>PO<sub>4</sub> (PDF#70-0702, labelled in gray in Figure 2.7a) can be clearly identified, which is supposed to be corresponding to the bulk crystals of Ag<sub>3</sub>PO<sub>4</sub> precipitated in high-temperature annealing, as have already been revealed in SEM images in Figure 2.3e–f and Figure 2.5. In the 400 °C-annealed microcrystalline film, the amount of such second-phase bulk crystals may be still too low for detection by XRD, whereas in the 500 °C-annealed film, the considerable amount of these bulk crystals (Figure 2.6) is already enough to give distinct X-ray diffraction signals of the new peaks. More importantly, when compared to the original phase of the Ag<sub>3</sub>PO<sub>4</sub> microcrystalline film, each characteristic peak of the second Ag<sub>3</sub>PO<sub>4</sub> phase shows a

slight shift to a higher  $2\theta$  value, indicating that non-negligible lattice contraction must have occurred in these second-phase Ag<sub>3</sub>PO<sub>4</sub> bulk crystals. It is assumed that during the high-temperature annealing in the oxygen-free conditions (*i.e.* vacuum) of the Ag<sub>3</sub>PO<sub>4</sub> films, a small amount of oxygen vacancies  $(V_0)$  may be generated in the lattice, despite the recovery of silver atoms-based intrinsic defects at the same time. It is well recognized that the formation of V<sub>O</sub> in oxide lattice would cause lattice contraction, [26–28] and the  $V_0$  in Ag<sub>3</sub>PO<sub>4</sub> has a much higher formation energy compared to intrinsic  $V_{Ag}$  and  $Ag_i$  defects because of the strong P–O covalent bond that must be broken to remove O atoms. [29–30] Therefore, it is rational that such  $V_0$  can be formed when annealing Ag<sub>3</sub>PO<sub>4</sub> in vacuum at  $\geq$  400 °C: For one thing, the relatively high temperature could afford the energy to break the strong P-O covalence; For another, the oxygen-free conditions could facilitate the release of lattice oxygen as O<sub>2</sub>. Finally, the existence of V<sub>0</sub> in the lattice of the recrystallized and precipitated second-phase Ag<sub>3</sub>PO<sub>4</sub> bulk crystals has resulted in such contraction. It is sometimes considered that such contraction derived from Vo, might in turn lead to the rearrangement of  $Ag_3PO_4$  cells, along with the disappearance of  $V_{Ag}$ .[2] In addition, each characteristic diffraction peak of the second-phase Ag<sub>3</sub>PO<sub>4</sub> shows lower intensities than the original phase, implying that the recrystallized Ag<sub>3</sub>PO<sub>4</sub> bulk crystals are more defective with relatively low crystallinity than the original Ag<sub>3</sub>PO<sub>4</sub> phase, which could also be regarded as an indirect evidence of the V<sub>0</sub> in the bulk crystals.

Meanwhile, it is noticeable that the photoactivity of  $Ag_3PO_4$  depends greatly on the exposure of specific crystal planes (*i.e.* facet engineering). Previous studies have pointed out that the {111} and {110} planes, especially the {111} plane of  $Ag_3PO_4$ , are the most favorable for photooxidation of water.[31–34] First, the {111} plane possesses

the highest surface energy because of the over-abundance of the dangling P–O bonds, which are designated oxidation sites. Second, the exposure of {111} plane also causes slicing of P–O bonds, which consequently lowers the layer-by-layer distance along the [111] direction. Further, because the photoholes in Ag<sub>3</sub>PO<sub>4</sub> possess the smallest mass, and thus greatest mobility along the [111] direction, the exposure of {111} facet could help to achieve more efficient charge separation and faster charge transfer of photoholes for water oxidation reaction on the surface.[31,35] In view of this, relative intensities of the (222) and (110) planes have also been calculated from the XRD patterns (Figure 2.7a), and shown in Table 2.1. Typically, it is found that only the 300 °C- and 400 °C-annealed Ag<sub>3</sub>PO<sub>4</sub> film (300 °C > 400 °C), whereas the Ag<sub>3</sub>PO<sub>4</sub> film annealed at 400 °C exhibits much increased diffraction intensity, and thus possibly much enhanced exposure rate of the (222) plane. Overall, it is supposed the facet exposure of the 400 °C-annealed film might be the most advantageous for photooxidation of water, as highlighted in yellow in Figure 2.7a.

Figure 2.8 displays the Raman spectroscopy analysis of as-obtained Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films before and after annealing at different temperatures. Raman spectroscopy is regarded as a powerful methodology for estimating the short-range order of solids.[36–37] It is revealed in Figure 2.8a that for the non-annealed Ag<sub>3</sub>PO<sub>4</sub> film (*i.e.* the black curve), five characteristic Raman scattering bands could be identified at 238, 423, 564, 909, and 1015 cm<sup>-1</sup> of the Raman shift, which can be assigned to the Raman-active vibrational modes of the rotation (238 cm<sup>-1</sup>), symmetric bending (423 cm<sup>-1</sup>), and asymmetric bending (564 cm<sup>-1</sup>) of [PO<sub>4</sub>] tetrahedron units, as well as the symmetric stretching (909 cm<sup>-1</sup>) and asymmetric stretching (1015 cm<sup>-1</sup>) of O–P–O

bonds, respectively.[38–39] After the Ag<sub>3</sub>PO<sub>4</sub> film was annealed at elevated temperature, several bands show obvious red-shifts. Specifically, the three bands located at ~423, 564 and 1015 cm<sup>-1</sup> shift to 408, 549 and 1001 cm<sup>-1</sup> respectively for the 400 °C-annealed Ag<sub>3</sub>PO<sub>4</sub> film. This red-shift could be explained by the phonon confinement effect and surface relaxation induced by increased crystal size after annealing at high temperature.[40–41] Meanwhile, the slightly increased intensities of the characteristic Raman bands indicate the improved crystallinity and lattice order of the Ag<sub>3</sub>PO<sub>4</sub> films annealed at elevated temperatures, which is in consistency with the above XRD results.

Furthermore, significant difference between the Raman spectra of the non-annealed and 500 °C-annealed Ag<sub>3</sub>PO<sub>4</sub> films may provide more information, as disclosed by a direct comparison of these two samples (Figure 2.8b). It is surprising that the Ag<sub>3</sub>PO<sub>4</sub> film after 500 °C-annealing behaves just like bulk Ag<sub>3</sub>PO<sub>4</sub>, with nearly no obvious Raman signals ascribed to the modes of [PO<sub>4</sub>] units being detected (*i.e.* the yellow background region). In Ag<sub>3</sub>PO<sub>4</sub> lattice, O–P–O bonds (*i.e.* the blue background region) are basic units, while the [PO<sub>4</sub>] tetrahedrons could be regarded as advanced combinations of O-P-O bonds. Naturally, due to the inductive effect of [PO4] units, the [AgO<sub>4</sub>] tetrahedrons are highly distorted in the lattice with two different O-Ag-O bond angles  $(\alpha, \beta)$  being found. The native defect of  $V_{Ag}$  generated from the separation of interior silver atoms from the supercell further makes such distortion of the [AgO<sub>4</sub>] units even more severe, contributing to the flexibility of the [PO<sub>4</sub>] units as a result.[29] Such flexibility finally brings about the distinct signals of the vibrational modes of the [PO<sub>4</sub>] units in the non-annealed, defective Ag<sub>3</sub>PO<sub>4</sub> film. On the contrary, the disappear of the modes of [PO<sub>4</sub>] units in the 500 °C-annealed film suggests the recovery of defects and consequently increase in lattice order that make  $[PO_4]$  units more fixed and difficult

to generate Raman vibrational signals, and can also help to confirm that the rearrangement of structural units and consequent recrystallization of  $Ag_3PO_4$  bulk crystals do not involve the linking of [PO<sub>4</sub>] tetrahedra.[42]



**Figure 2.8.** Room-temperature Raman spectra of (a) all the obtained  $Ag_3PO_4$  films, and (b) a more detailed comparison of the non-annealed and 500 °C-annealed  $Ag_3PO_4$  films. The abbreviation NA in panel (a) stands for the non-annealed  $Ag_3PO_4$  film. The different singly, doubly, and triply degenerate vibrational modes are labelled by  $A_1$ , E, and  $T_2$ , respectively.



**Figure 2.9.** Room-temperature (a) DRS spectra and (b) dark-field scattering spectra of the obtained Ag<sub>3</sub>PO<sub>4</sub> films before and after annealing at different temperatures.

Poor light absorption of the pristine  $Ag_3PO_4$  films, as has been already discussed in the Introduction section, is an essential issue needed to be addressed. DRS spectra of the  $Ag_3PO_4$  films before and after annealing at different temperatures, as exhibited in Figure 2.9a, have been recorded to unveil their substantially different light absorption profiles. It could be known that the absorption edges attributed to the charge transfer response of  $Ag_3PO_4$  from its valence band to the conduction band, of the non-annealed  $Ag_3PO_4$  film, as well as the films annealed at  $\leq 200$  °C appear at 560–570 nm, while the

absorption edges of the films annealed at  $\geq 300$  °C show up at 520–530 nm, which are quite near the standard value (~530 nm). Such a distinct red-shift in the absorption edge of the non-annealed and low-temperature annealed Ag<sub>3</sub>PO<sub>4</sub> films should be resulted from the large amount of native defects, especially  $V_{Ag}$  originally existed in the lattice. Previous first-principles calculations have disclosed the influence of  $V_{Ag}$  on the electronic properties of Ag<sub>3</sub>PO<sub>4</sub>.[29-30,43-45] Ideally, intrinsic V<sub>Ag</sub> in Ag<sub>3</sub>PO<sub>4</sub> would act as defect states near the valence band edge to extend the absorption edge of  $Ag_3PO_4$ in visible light band, typically from 530 nm to 560 nm.[29] Further, it is clear that for Ag<sub>3</sub>PO<sub>4</sub> films annealed  $\leq 200$  °C, including the non-annealed film, non-negligible broad absorption in the sub-bandgap region of  $\geq \sim 560$  nm could be identified. This may come from three aspects: 1) The plasmon-enhanced visible light absorption arising from extensively existed  $Ag^0 NPs[46-47]$  on the surface of these defective  $Ag_3PO_4$  films, as can be seen in Figure 2.2a-c; 2) Defect-state absorption from the large numbers of native  $V_{Ag}$  and  $Ag_i$  defects, as such a phenomenon has been widely discussed in other semiconductors like TiO<sub>2</sub> and SrTiO<sub>3</sub>;[48-49] 3) Lower light scattering in the sub-bandgap region, as confirmed by the dark-field scattering spectra in Figure 2.9b. Therefore, after annealing the Ag<sub>3</sub>PO<sub>4</sub> films at  $\geq$  300 °C, the recovery of native defects like  $V_{Ag}$ , decrease in the amount of  $Ag^0$  NPs on the surface, as well as increased light scattering in the sub-bandgap region (Figure 2.9b) in the resulted Ag<sub>3</sub>PO<sub>4</sub> films, have together decreased the sub-bandgap absorption to a relatively lower level.

More importantly, it is noticeable that the bandgap absorption of these Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films get greatly improved after annealing. The bandgap absorption of the Ag<sub>3</sub>PO<sub>4</sub> film essentially increases if increased annealing temperature was applied, following the variation of the 400 °C-annealed film > 500 °C > 300 °C > 200 °C >

non-annealed > 100 °C-annealed Ag<sub>3</sub>PO<sub>4</sub> film. Particularly for the Ag<sub>3</sub>PO<sub>4</sub> films annealed at  $\geq$  400 °C, the light absorbance exceeds 0.7 across the whole bandgap region with sharp onset slopes near the absorption edge, and can even reach as high as ~0.9 in the ultraviolet band. Overall, this substantial improvement could be ascribed to the improved crystalline quality as has been confirmed by XRD and Raman results.[5] Meanwhile, the possibly existed small amount of  $V_0$  in the recrystallized Ag<sub>3</sub>PO<sub>4</sub> bulk crystals in 400 °C- and 500 °C-annealed Ag<sub>3</sub>PO<sub>4</sub> films might have also made contributions to achieving the superior absorbance, since the capability of  $V_0$  to enhance the light absorption and redox reactivity of photocatalysts has been widely reported in literature.[6,50–52]



**Figure 2.10.** High-magnification SEM image of the 100 °C-annealed Ag<sub>3</sub>PO<sub>4</sub> film. Red arrows indicate the positions where pores generated by gas (*e.g.*, water molecules) release or desorption.

Besides, the reason why the 100 °C-annealed  $Ag_3PO_4$  film shows lower absorption than non-annealed film is maybe because such a low annealing temperature has introduced more defects along with pores into the film *via* the release of adsorbed gas



**Figure 2.11.** PEC water oxidation performance of the Ag<sub>3</sub>PO<sub>4</sub> microcrystalline film photoanodes: (a) the dependence of photocurrent density minus dark current density  $(J_{\text{light}} - J_{\text{dark}})$  on applied potential. Inset shows an enlarged view of the potential range from 0.330 to 0.405 V *vs.* Ag/AgCl. NA stands for the non-annealed Ag<sub>3</sub>PO<sub>4</sub> photoanode. (b) Current density–time (J-t) curves of the photoanodes fabricated from non-annealed and 400 °C-annealed Ag<sub>3</sub>PO<sub>4</sub> films during the stability test at an applied potential of 1 V *vs.* Ag/AgCl.

(*e.g.*, water desorption), while is still not high enough for the recovery of native defects, as witnessed by the high-magnification SEM image displayed in Figure 2.10. In addition, the slightly lower light absorption of the 500 °C-annealed  $Ag_3PO_4$  film than the 400 °C-annealed one is maybe because the extensively distributed bulk crystals on the surface of the 500 °C-annealed  $Ag_3PO_4$  microcrystalline film, as revealed in Figure 2.3f and Figure 2.4f, may have decreased the surface area and also increased undesired light reflection of the film. In generally, it is considered that the 400 °C-annealed  $Ag_3PO_4$  film shows the most appealing light absorption profiles among all the obtained microcrystalline films.



**Figure 2.12.** Comparisons of the photocurrent density (red curves) *vs.* dark current density (black curves) of different  $Ag_3PO_4$  microcrystalline film photoanodes.



**Figure 2.13.** Electrochemical impedance spectroscopy (EIS) Nyquist plots of the non-annealed and 400 °C-annealed Ag<sub>3</sub>PO<sub>4</sub> film photoanodes. The insert shows the equivalent circuit consisting of a serial resistor ( $R_S$ ), a resistor–capacitor circuit showing the parallel combination of the interfacial charge transfer resistance ( $R_{CT}$ ), and constant phase element (CPE) attributed to the interface between the electrode surface and the electrolyte solution.

The substantially improved physicochemical properties with regard to native defects like  $V_{Ag}$ , surface precipitation of  $Ag^0$  NPs, and bandgap light absorption of the  $Ag_3PO_4$ microcrystalline films obtained after vacuum annealing, together with their intrinsically high photoactivity, suggest the possibility towards improved stability and efficiency of PEC water oxidation. Figure 2.11a summarizes the dependency of photocurrent density minus dark current density ( $J_{light} - J_{dark}$ ) on applied potential. It is obvious that the same variation in [ $J_{light} - J_{dark}$ ] *vs. E* curves are found as the light absorption profiles, *i.e.* 400 °C > 500 °C > 300 °C > 200 °C > non-annealed > 100 °C-annealed microcrystalline film photoanode, indicating that light absorption, especially the bandgap absorption could be the dominant factor to the PEC performance of these  $Ag_3PO_4$  film photoanodes. Figure 2.12 compares the photocurrent density *vs.* dark current density of all the photoanodes, confirming that the current density of these  $Ag_3PO_4$  microcrystalline film photoanodes is mainly originated from the photocatalytic effect under light irradiation, instead of the electrocatalytic effect.

Particularly, much improved PEC water oxidation stability (Figure 2.11b) as well as efficiency of the 400 °C-annealed Ag<sub>3</sub>PO<sub>4</sub> film photoanode are demonstrated compared to the non-annealed Ag<sub>3</sub>PO<sub>4</sub> film, suggesting that the methodology of vacuum annealing



**Figure 2.14.** SEM images of different magnifications showing the surface morphology after the stability test of the photoanodes fabricated from (a-c) non-annealed and (d-f) 400 °C-annealed Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films.



Figure 2.15. XRD patterns of the photoanodes fabricated from the non-annealed and 400 °C-annealed Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films after the stability test.



**Figure 2.16.** Representative XRD patterns of the non-annealed (*i.e.* blue-background part) and 400 °C-annealed (*i.e.* orange-background part)  $Ag_3PO_4$  microcrystalline film photoanode before (*i.e.* dashed curves) and after (*i.e.* solid curves) the stability test. The dark yellow dashed lines indicate the peak positions of metallic  $Ag^0$  phase.

could be useful for improving the PEC performance of the Ag<sub>3</sub>PO<sub>4</sub> microcrystalline film photoanodes because of its ability to reduce the native defects like  $V_{Ag}$ , remove the Ag<sup>0</sup> NPs precipitated on the surface, and also increase the bandgap light absorption of these Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films. The significantly decreased amount of defects in the 400 °C-annealed Ag<sub>3</sub>PO<sub>4</sub> microcrystalline film photoanode, meanwhile, can also be known from its obviously shifted onset potential to a more negative value, as shown in the inset of Figure 2.11a.[53–55]

Electrochemical impedance spectroscopy (EIS) analytical technique is a strong tool to study the charge transfer process at the SEJ interface of PEC water oxidation reaction. In the derived Nyquist plots, the arch at the high frequency stands for the process of charge transfer, and its diameter indicates the charge carrier mobility; *i.e.* the smaller diameter the semicircle has, the faster charge carriers can transfer.[56] Herein, Figure 2.13 displays the slightly decreased interfacial charge transfer rate from the non-annealed Ag<sub>3</sub>PO<sub>4</sub> microcrystalline film photoanode to the 400 °C-annealed one. This is rational considering that defects are well recognized to be favorable for the conductivity,[57] nevertheless such a slight decrease in conductivity of the annealed microcrystalline Ag<sub>3</sub>PO<sub>4</sub> film photoanodes should have minor influence on the PEC water oxidation efficiency (*i.e.* expressed as the photocurrent density) when compared with the distinctly improved light absorption.

Specifically, the much improved PEC water oxidation stability, as could be known from Figure 2.11b, can be understood from the following two aspects: For one thing, the surface erosion of the non-annealed Ag<sub>3</sub>PO<sub>4</sub> microcrystalline film photoanode is much stronger than the 400 °C-annealed one (Figure 2.14); For another, the decomposition of

Ag<sub>3</sub>PO<sub>4</sub> into metallic Ag<sup>0</sup> in stability test of the 400 °C-annealed microcrystalline film is less serious as well, as witnessed by the characteristic diffraction peaks corresponding to the metallic Ag<sup>0</sup> phase in the XRD patterns displayed in Figure 2.15 and Figure 2.16. Both points indicate less loss of the active material of Ag<sub>3</sub>PO<sub>4</sub> in the 400 °C-annealed microcrystalline film photoanode than the non-annealed one. Overall, it can be deduced that the reduced defects in the lattice of Ag<sub>3</sub>PO<sub>4</sub>, as well as the strong interaction among adjacent Ag<sub>3</sub>PO<sub>4</sub> crystallites after annealing,[3] have together inhibited the precipitation of metallic Ag<sup>0</sup> NPs on the surface, and dissolution of Ag<sup>+</sup> from Ag<sub>3</sub>PO<sub>4</sub> into aqueous solution, and finally contributed to the enhanced stability for PEC water oxidation.

## 2.4. Conclusions

In this chapter, solution-processed large-area Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films are first prepared *via* a room-temperature, air-exposed immersion reaction. The post-treatment of vacuum annealing at different temperatures (*i.e.* 100, 200, 300, 400 and 500 °C) has been investigated as a strategy for solving the problem of native defects and low light absorption of Ag<sub>3</sub>PO<sub>4</sub>. It is found that below 400 °C is the |Recovery| stage for Ag<sub>3</sub>PO<sub>4</sub> microcrystals when defect amount decreases and therefore crystallinity increases. From 400 °C, the |Recrystallization| stage and also precipitation of Ag<sub>3</sub>PO<sub>4</sub> bulk crystals start, and afterwards till 500 °C is the |Recrystallization + Grain growth| stage for Ag<sub>3</sub>PO<sub>4</sub>, when more faceted and larger bulk crystals of Ag<sub>3</sub>PO<sub>4</sub> precipitate extensively. In this series of Ag<sub>3</sub>PO<sub>4</sub> film photoanodes, the same variation in [*J*<sub>light</sub> – *J*<sub>dark</sub>] *vs. E* curves are found as the light absorption profiles, *i.e.* 400 °C > 500 °C > 300 °C > 200 °C > non-annealed > 100 °C-annealed Ag<sub>3</sub>PO<sub>4</sub> film, indicating that light absorption, especially the bandgap absorption could be the dominant factor to the PEC performance. Particularly, much improved PEC water oxidation stability and efficiency of the 400 °C-annealed Ag<sub>3</sub>PO<sub>4</sub> film photoanode are demonstrated compared with the non-annealed Ag<sub>3</sub>PO<sub>4</sub> film, suggesting that the methodology of vacuum annealing could be utilized for improving the PEC performance of Ag<sub>3</sub>PO<sub>4</sub> film photoanodes because of its ability to reduce the native defects like  $V_{Ag}$ , remove the Ag<sup>0</sup> NPs on the surface, and increase the bandgap light absorption of these Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films.

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# **3.** Porous microcrystalline Ag<sub>3</sub>PO<sub>4</sub> film photoanodes with improved light absorption and surface area for PEC water oxidation

### **3.1.** Supplementary introduction

Microcrystalline film-based photoelectrodes generally exhibit better stability than nanocrystalline powder-deposited photoelectrodes.[1] Nevertheless, the increase in the particle/crystallite size of the active semiconductor material of the photocatalyst would unfortunately result in the decrease of active surface areas, and consequently reduce the photocurrent densities of the photoelectrodes. Hence in view of this, introducing porous structures into the microcrystalline semiconductor film photoelectrodes seems to be an efficient strategy to realize both improved photoactivity and stability. Recently, porous structures have been demonstrated to be able to improve the PEC performance of the photocathode materials such as NiO,[2] and Cu<sub>2</sub>ZnSnS<sub>4</sub>,[3] and also a lot of photoanode materials such as  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>,[4] WO<sub>3</sub>,[5] and fluorine-doped rutile TiO<sub>2</sub>.[6]

In this chapter, we have prepared large-area ( $4.5 \times 4.5 \text{ cm}^2$ ) and solution-processed Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films with tunable porosity controlled by the reaction time as novel photoanodes for PEC water oxidation. Enhanced PEC water oxidation activity as well as stability were displayed by the Ag<sub>3</sub>PO<sub>4</sub> porous films with the optimized porosity obtained for extended reaction times ( $\geq 20$  h), suggesting great potential of these porous Ag<sub>3</sub>PO<sub>4</sub> photoanodes to support efficient solar-to-fuel energy conversion and other PEC applications.

#### **3.2. Experimental details**

The as-purchased silver foils (99.98%, 0.1 mm thick, Nilaco) were first cut into pieces (4.5×4.5 cm<sup>2</sup>) and cleaned by dilute nitric acid solution (Sigma-Aldrich) and acetone ( $\geq$  99.5%, Wako), respectively. Briefly, an aqueous solution was first prepared by mixing 7.2 g of sodium phosphate monobasic (NaH<sub>2</sub>PO<sub>4</sub>,  $\geq$  99.0%, Sigma-Aldrich), 103.8 g of hydrogen peroxide (H<sub>2</sub>O<sub>2</sub>, 30.0–35.5%, Sigma-Aldrich) and 10 g of polyvinylpyrrolidone (PVP, M<sub>W</sub>  $\approx$  55,000, Sigma-Aldrich) in distilled water (380 mL) under stirring at room temperature till the solution became transparent. A piece of Ag foil was then immersed into the mixture solution and kept for different times varied from 1 h to 24 h. Afterwards, the foil was taken out of the solution and rinsed with anhydrous ethanol ( $\geq$  99.5%, Wako) and distilled water thoroughly and dried in air naturally for fabrication into photoanodes.

For materials characterization, the scanning electron microscope (SEM) images were recorded using a field-emission SEM (FE-SEM, S4700, Hitachi, Japan). X-ray diffraction (XRD) data were acquired on a Rigaku Miniflex II-MW (Japan) X-ray diffractometer working with Cu-K $\alpha$  radiation (1.5406 Å). X-ray photoelectron spectroscopy (XPS) analysis was carried out with a PHI 5000 VersaProbe spectrometer (ULVAC-PHI, Japan) using an Al-K $\alpha$  X-ray source (1486.6 eV). The electron analyzer was operated at the pass energy of 23.5 eV for the narrow scans. Raman spectra were collected on a laser confocal Raman micro-spectrometer (in*Via* reflex, Renishaw, UK) equipped with an integrated optical microscope (DM-2500, Leica, Germany). The excitation light source of a 632.8 nm He–Ne laser, laser power of 1.7 mW (10% × 17 mW), grating of 1800 1/mm, exposure time of 5.0 s, and accumulation of three times
were applied for all Raman measurements. UV-visible diffuse-reflectance spectra (DRS) were obtained on a JASCO V-670 (Japan) spectrophotometer equipped with an integrating sphere.



**Figure 3.1.** (a) Schematic illustration for the gas bubbles release-induced etching into porous structures during the preparation of microcrystalline Ag<sub>3</sub>PO<sub>4</sub> films; (b) Digital photography image of the original silver foil substrate ( $4.5 \times 4.5 \text{ cm}^2$ ) and the finally obtained Ag<sub>3</sub>PO<sub>4</sub> film grown on silver after 24 h of the reaction; (c–e) Digital photography images of the chemical reaction phenomena after immersing the silver foil into the mixture solution for 5 min (c), 5 h (d) and 24 h (e), respectively.

For evaluation of the PEC performance, all measurements were carried out in a

three-electrode cell. The as-fabricated Ag<sub>3</sub>PO<sub>4</sub> photoanodes, a Pt wire, an Ag/AgCl electrode, and the aqueous solution of Na<sub>3</sub>PO<sub>4</sub> (0.02 M, pH  $\approx$  11.3) were utilized as the working electrode, counter electrode, reference electrode and electrolyte, respectively. Potential of working electrode was controlled with a VersaSTAT-4 potentiostat (Princeton Applied Research). Before each measurement, the electrolyte solution was stirred and purged with high-purity nitrogen gas for above 20 min. During measurements, the Ag<sub>3</sub>PO<sub>4</sub> working electrodes were illuminated with a 500 W xenon lamp delivering an irradiance of 285 mW/cm<sup>2</sup>. The scan rate was 10 mV/s for collecting each current density–potential (*J*–*E*) curve from negative to positive bias. The electrochemical impedance spectroscopy (EIS) data were acquired using the same electrochemical measurement system under dark conditions at an applied potential of 0.1 V vs. Ag/AgCl varying the frequency between 100,000 and 1 Hz of a sinusoidal perturbation with 10 mV amplitude.

#### **3.3. Results and discussion**

The Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films were grown on silver foil substrates *via* a facile room-temperature and air-exposed solution process without any post-treatment process (*e.g.*, annealing). Figure 3.1 shows a schematic illustration of the origin of the porosity of the Ag<sub>3</sub>PO<sub>4</sub> films — induced by the release of H<sub>2</sub> bubbles. Briefly in the reaction, pieces of  $4.5 \times 4.5$  cm<sup>2</sup> Ag substrates were immersed into a mixture solution containing NaH<sub>2</sub>PO<sub>4</sub> as the PO<sub>4</sub><sup>3-</sup> source and H<sub>2</sub>O<sub>2</sub> as the oxidation agent. Ag<sub>3</sub>PO<sub>4</sub> was then generated on the substrate from the chemical reaction described as follows:

$$2 \operatorname{NaH_2PO_4} + 2 \operatorname{H_2O_2} + 6 \operatorname{Ag} \rightarrow 2 \operatorname{Ag_3PO_4} + 2 \operatorname{NaOH} + 2 \operatorname{H_2O} + \operatorname{H_2\uparrow}$$

Figure 3.1c-e show the gas release phenomenon at the time points of 5 min, 5 h as well as 24 h during the reaction. Overall, it can clearly be observed that as the reaction proceeded, a large amount of  $H_2$  bubbles were generated from the reaction and released from the surface of the substrates. Specifically, at the initial stage of the reaction (*i.e.* 5 min), the size of the gas bubbles turned out to be quite large, suggesting that the reaction was rather violent at this time due to the strong oxidation activity of hydrogen peroxide. As the time reached 5 h, the size of the bubbles decreased obviously, making it look more like mist rather than bubbles, indicating that the reaction became milder at this stage owing to the decreased concentration of reactants. Finally, the amount of the mist-like bubbles released from the surface of the substrate decreased when the reaction time reached 24 h as can be understood from the gradual exhaustion of reactants.



**Figure 3.2.** Digital micrographs of the Ag<sub>3</sub>PO<sub>4</sub> microcrystalline samples obtained with different reaction times on as-purchased silver foils (the most left one). All Ag<sub>3</sub>PO<sub>4</sub> samples were  $4.5 \times 4.5$  cm<sup>2</sup> (as shown in Figure 3.1b) and were cut into  $1.0 \times 1.5$  cm<sup>2</sup> to fabricate the photoanodes.

The digital micrographs in Figure 3.1b reveal that after the reaction, the color of as-obtained Ag<sub>3</sub>PO<sub>4</sub> films were close to dark yellow, revealing their possibly good

photoactivity when serving as photoelectrodes. The digital micrographs of all samples are shown in Figure 3.2, demonstrating the color change more completely.



**Figure 3.3.** Representative (a) XRD pattern and (b–d) XPS spectra of Ag 3d (b), P 2p (c) and O 1s (d) region of the obtained porous microcrystalline Ag<sub>3</sub>PO<sub>4</sub> films. The binding energy in panels (b–d) is calibrated by using the contaminant C 1s peak at 284.8 eV as the reference. All data are shown using the 20-h Ag<sub>3</sub>PO<sub>4</sub> microcrystalline film sample.

To determine the component and phase purity of the obtained porous films on the silver substrates, the XRD pattern and XPS spectra of the samples were recorded, as exhibited in Figure 3.3. It is clear in Figure 3.3a that except the small diffraction peak at  $2\theta \approx 44.3^{\circ}$  corresponding to the (200) crystal plane of silver (Ag, PDF#04-0783) which should come from the silver substrate, all distinct diffraction peaks could be well

indexed to the cubic P(-4)3n phase of silver phosphate (Ag<sub>3</sub>PO<sub>4</sub>, PDF#70-0702). High crystallinity of the obtained photoanode material could be deduced from the distinct shapes and high intensities of the characteristic diffraction peaks in Figure 3.3a. Besides the XRD analysis, XPS narrow scans have also been performed to verify the chemical environment of the different elements. The XPS spectra of Ag 3d, P 2p and O 1s regions are shown in Figure 3.3b–d, and all binding energies have been calibrated by using the contaminant C 1s peak at 284.8 eV as the reference. In Figure 3.3b, the Ag  $3d_{3/2}$  and Ag  $3d_{5/2}$  peaks could be observed at 373.19 and 367.20 eV of binding energy, respectively, suggesting the presence of Ag<sup>+</sup> in as-synthesized porous microcrystalline films. In Figure 3.3c–d, the peak of P 2p at around 131.94 eV and O 1s at 529.90 eV should be originated from the crystal lattice of P<sup>5+</sup> and O<sup>2-</sup> in PO4<sup>3-</sup> units according to literature reports.[7–8]



**Figure 3.4.** (a) TEM and (b) HRTEM images of a fragment obtained from as-prepared 1-h Ag<sub>3</sub>PO<sub>4</sub> film by scratching, grinding and sonification. The fine particles in the surroundings [*i.e.* panel (a)] came mainly from the electron beam-induced precipitation of silver nanoparticles, which is a common phenomenon in Ag<sub>3</sub>PO<sub>4</sub> nanomaterials.[9]



**Figure 3.5.** SEM images of the obtained porous  $Ag_3PO_4$  photoanodes after (a) 1 h, (b) 2 h, (c) 5 h, (d) 9 h, (e) 20 h and (f) 24 h of the reaction. Inserts in each panel are SEM images recorded at a higher magnification for each sample.

Moreover, the HRTEM image of as-obtained sample, as shown in Figure 3.4b, in which the (211) crystal plane of  $Ag_3PO_4$  could be clearly identified, together with the above-discussed XRD and XPS results, proves that the pure phase  $Ag_3PO_4$  was formed successfully on the silver substrate.



**Figure 3.6.** Cross-sectional SEM images of the (a–b) nearly non-porous 1-h Ag<sub>3</sub>PO<sub>4</sub> microcrystalline film and (c–d) highly porous 20-h Ag<sub>3</sub>PO<sub>4</sub> film. The insert in panel (b) shows an enlarged image of the area marked by the red box. The high-resolution SEM images in panels (b) and (d) were both recorded with a 30 °-tilt angle. The thicknesses of the Ag<sub>3</sub>PO<sub>4</sub> layers of these two samples are both estimated to be around 10  $\mu$ m.

The morphologies of the obtained Ag<sub>3</sub>PO<sub>4</sub> films on silver substrates are revealed by the SEM images of Figure 3.5 and Figure 3.6. Typically, the appearance of the Ag<sub>3</sub>PO<sub>4</sub> film obtained after 1 h of the reaction (Figure 3.5a) looks similar to other previously reported solution-processed semiconductor microcrystalline electrodes like  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>[10] and Cu<sub>2</sub>O[11] with hardly no porous structures on the surface being identified. However, as the reaction proceeded till 2 h long, the surface of Ag<sub>3</sub>PO<sub>4</sub> became much rougher. This roughness was ascribed to the etching effect of the H<sub>2</sub> gas bubbles released (Figure 3.5b). Further, when the reaction time was extended to 5 h, 9 h, 20 h and finally 24 h,

distinct interconnected porous structures could clearly be distinguished from the SEM images of the samples (Figure 3.5c-f). The porous structures might be resulted from the continuous formation and release of H<sub>2</sub> bubbles during the chemical reaction generating Ag<sub>3</sub>PO<sub>4</sub> from the Ag substrate and  $PO_4^{3-}$  in the solution. Cross-sectional SEM images displayed in Figure 3.6 of the nearly non-porous 1-h Ag<sub>3</sub>PO<sub>4</sub> film and highly porous 20-h sample further help to confirm the porosity degree of the films. Besides, from the cross-sectional SEM images, it is also found that the thickness of the Ag<sub>3</sub>PO<sub>4</sub> film grown on silver foils was almost independent with the reaction time when it was longer than 1 h, as the thicknesses of both the 1-h and 20-h Ag<sub>3</sub>PO<sub>4</sub> films were just similar to each other. According to the recorded cross-sectional SEM images shown in Figure 3.6, the nearly non-porous 1-h Ag<sub>3</sub>PO<sub>4</sub> microcrystalline film and the highly porous 20-h Ag<sub>3</sub>PO<sub>4</sub> film exhibit the similar thickness of ~10  $\mu$ m. Therefore, it is assumed that when the reaction time was longer than 1 h, the thickness of as-formed Ag3PO4 film would hardly grow, and hence was almost independent with the reaction time. Based on this, it is further assumed that the growing process of Ag<sub>3</sub>PO<sub>4</sub> films can be described roughly as the following 2 different stages:

1. The 1<sup>st</sup> stage ( $0 \le \text{reaction time} \le 1$  h): In this stage, because the film was initially quite thin and kept growing thicker gradually, the position where the H<sub>2</sub> bubbles drastically generated was still quite near the top surface of the film. Therefore, these bubbles were likely to be released quickly from the top surface of Ag<sub>3</sub>PO<sub>4</sub>. Meanwhile, according to our observation, the generated Ag<sub>3</sub>PO<sub>4</sub> was not enough to cover everywhere on the silver foil in this stage and there were still a lot of exposed silver surfaces, as illustrated in Figure 3.7b. In this case, the bubbles were more likely to be released from the Ag<sub>3</sub>PO<sub>4</sub>–Ag foil interfaces on the silver side, and therefore would

have less influence on the formed Ag<sub>3</sub>PO<sub>4</sub> film. In addition, the millimeter-sized H<sub>2</sub> bubbles in this stage may be too large to introduce micrometer- or even nanometer-sized pores, as revealed by the SEM images of the porous Ag<sub>3</sub>PO<sub>4</sub> photoanodes (Figure 3.5). Therefore, there are almost no pores being observed in the 1-h sample. In fact, it could be seen from Figure 3.7a that if the reaction time was further decreased to 20 min (*i.e.* < 1 h), the produced Ag<sub>3</sub>PO<sub>4</sub> microcrystalline film would be more faceted. Nevertheless, these samples are not included and discussed in this chapter since the as-formed Ag<sub>3</sub>PO<sub>4</sub> film at this stage could not cover the whole surface of the Ag foil substrate.



**Figure 3.7.** (a) SEM image of the obtained Ag<sub>3</sub>PO<sub>4</sub> microcrystalline film after 20 min of the reaction; (b) Schematic illustration of the coverage of the generated Ag<sub>3</sub>PO<sub>4</sub> film on the silver foil during the first reaction stage ( $0 \le$  the reaction time  $\le 1$  h).

2. The 2<sup>nd</sup> stage (reaction time  $\geq 1$  h): As the reaction proceeded into this stage, the as-formed Ag<sub>3</sub>PO<sub>4</sub> film was already quite thick and had covered the whole silver foil substrate. Meanwhile, the size of the bubbles also decreased as the reaction became milder. In this stage, due to the gradual exhaustion of the reactants, the film thickness might grow only a little. In this situation, the small H<sub>2</sub> bubbles could diffuse slowly from the Ag-Ag<sub>3</sub>PO<sub>4</sub> interfaces at the bottom, to the upper surface of the Ag<sub>3</sub>PO<sub>4</sub> film, and bring about large numbers of of pores during this process.

The small nanoparticles (*i.e.* brighter white dots) which could be observed in Figure 3.5a–b might be the precipitation of silver from  $Ag_3PO_4$  taking place during the violent chemical reaction at this stage. As the reaction became milder for longer reaction times, the amount of such white dots being observed on the surface seems to decrease, as seen in Figure 3.5c–f. The precipitation of silver nanoparticles did not come from the silver foil substrate which was covered by as-formed thick  $Ag_3PO_4$  film, but was from the decomposition of the  $Ag_3PO_4$  on the surface. It is well known that silver based semiconductors (*e.g.*,  $Ag_3PO_4$  and AgI) are important fast ionic conductors.[12–14] Intrinsically, the ionic mobility of  $Ag^+$  in such materials is fairly higher than the other anions. Particularly in  $Ag_3PO_4$ , the ionic mobility of  $Ag^+$  is higher than  $PO_4^{3-}$ ,[15] and the mobility of electrons is higher than holes as well.[16]

In the synthetic experiment, because the amount of the reactants gradually decreased as time went by, it is reasonable that the chemical reaction was more drastic during the initial and early periods than later periods, resulting in the phenomenon that the gas generation and release were also more drastic initially. There were plenty of electrons transferring between silver species and hydrogen species during the redox reaction, and



**Figure 3.8.** Images recorded with an optical scattering (dark field) microscope of the as-synthesized  $Ag_3PO_4$  microcrystalline films with different reaction times varied from (a) 1 h, (b) 2 h, (c) 5 h, (d) 9 h, (e) 20 h, to (f) 24 h.

the dominant reaction was the formation of  $Ag_3PO_4$  solid film and  $H_2$  from NaH<sub>2</sub>PO<sub>4</sub>, H<sub>2</sub>O<sub>2</sub> and Ag, as described in the chemical equation. However, in consideration of the higher ionic mobility of Ag<sup>+</sup> than PO<sub>4</sub><sup>3–</sup>, and easy transport of electrons in as-formed Ag<sub>3</sub>PO<sub>4</sub> solid, it is also rational that in such a high-concentration flow of the electrons during the redox reaction, a part of the Ag<sup>+</sup> in Ag<sub>3</sub>PO<sub>4</sub> might be reduced by electrons. As a result,  $Ag^0$  NPs were precipitated. Afterwards, as the reaction became milder for longer reaction times, the amount of  $Ag^0$  precipitation tended to decrease as well, due to the decrease of the electron flow in  $Ag_3PO_4$  in milder redox reactions. In fact, the reduction of  $Ag^+$  in  $Ag_3PO_4$  by electrons, and consequent precipitation of  $Ag^0$  NPs could be regarded as a common phenomenon in  $Ag_3PO_4$ -based materials. The reduction of  $Ag^+$  in  $Ag_3PO_4$  into  $Ag^0$  NPs by the strong electron beam irradiation during the TEM characterization can also be regarded as an indirect evidence, as exhibited in Figure 3.4. This phenomenon of the surface precipitation of  $Ag^0$  NPs on as-prepared  $Ag_3PO_4$  films could also be identified using the dark-field scattering microscopy technique, as shown in Figure 3.8 (*i.e.* white dots in the dark-field images).

In summary, the highly porous structures of the  $Ag_3PO_4$  microcrystalline films obtained with relatively longer reaction times (*i.e.* 9 h, 20 h and 24 h) should have the potential to increase the active surface areas and, hence, enhance the light absorption and subsequent PEC water oxidation activity.

On the other hand, it should also be noted that the generation of such distinct porous structure also depends greatly on the addition of oxidation agent H<sub>2</sub>O<sub>2</sub> and surfactant. If H<sub>2</sub>O<sub>2</sub> was added drop by drop, no highly porous structures could be formed even after 8 h of the reaction, as revealed in Figure 3.9. This indicates that the initial supersaturation condition which can be realized by adding H<sub>2</sub>O<sub>2</sub> once at the same time is necessary for the generation of porous films. Meanwhile, the surfactant of PVP has also played an important role for the generation of high-quality Ag<sub>3</sub>PO<sub>4</sub> surfaces. It is clear in Figure 3.10 that PVP of  $M_W \approx 40,000$ , instead of PVP of  $M_W \approx 55,000$ , is not enough to protect the Ag<sub>3</sub>PO<sub>4</sub> microcrystals to get defined and clean surfaces.



**Figure 3.9.** SEM images of the  $Ag_3PO_4$  microcrystalline films prepared by adding the oxidation agent  $H_2O_2$  drop by drop, and after the reaction time of (a–b) 1 h, (c–d) 2 h, (e–f) 5 h and (g–h) 8 h, respectively.



**Figure 3.10.** SEM images of the Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films prepared by utilizing the PVP of  $M_W \approx 40,000$  as the surfactant, and after the reaction time of (a–b) 2 h and (c–d) 9 h, respectively.

Raman spectra of the samples, as displayed in Figure 3.11, confirm the existence of  $[PO_4]$  units and O–P–O bonds in the sample materials. Specifically, seven characteristic Raman scattering peaks could be experimentally identified at the Raman shifts of around 147, 238, 416, 561, 718, 910 and 1003 cm<sup>-1</sup>. These peak positions could be assigned to the Raman-active vibration modes of rotation and translation of  $[PO_4]$  units (147 and 238 cm<sup>-1</sup>), symmetric bending (416 cm<sup>-1</sup>) and asymmetric bending (561 cm<sup>-1</sup>) of  $[PO_4]$  units, symmetric stretching (718 and 910 cm<sup>-1</sup>), and asymmetric stretching (1003 cm<sup>-1</sup>) of O–P–O bonds, respectively.[9] The observation of multiple resonant Raman peaks indicates that the synthesized Ag<sub>3</sub>PO<sub>4</sub> films possessed favorable optical quality.[17] In addition, the slightly increased intensities of the characteristic Raman

peaks of the samples obtained with extended reaction times suggest a gradual improvement of the crystallinity of  $Ag_3PO_4$  as the reaction proceeded.



**Figure 3.11.** Room-temperature Raman spectra of the porous  $Ag_3PO_4$  microcrystalline photoanodes obtained with different reaction times. Insert displays a representative optical micrograph image of the measured area on the  $Ag_3PO_4$  surface during the micro-Raman characterization.

DRS spectra of the obtained Ag<sub>3</sub>PO<sub>4</sub> films, as shown in Figure 3.12, indicate that as the extended reaction time brought about promoted porosity, the Ag<sub>3</sub>PO<sub>4</sub> films grown on silver substrates demonstrated decreased reflectance and hence increased absorption over the whole ultraviolet–visible light band. This might be attributed to the enhanced light confinement within the interconnected pores of the obtained Ag<sub>3</sub>PO<sub>4</sub> films,[18] and consequently lead to improved light capture capability for PEC water splitting. The absorption edges of all  $Ag_3PO_4$  samples turn out to be close to each other at around 530 nm, which is consistent with the standard value.



**Figure 3.12.** UV–visible DRS spectra of the synthesized Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films grown on silver substrates with different reaction times.

In addition, a small peak can be seen in the wavelength range of 400–450 nm, as indicated in blue background in Figure 3.12. This small contribution may come from the scattering of precipitated  $Ag^0$  NPs. It seems that as the reaction time increases from 1 h to 24 h, the intensity of this peak decreases slightly. This phenomenon has also been confirmed by the optical scattering microscopy of the  $Ag_3PO_4$  films (Figure 3.8) since the number of the blue–green scattering dots originating from  $Ag^0$  NPs decreases in the samples obtained with longer reaction times. These results imply that when longer reaction times were used, less  $Ag^0$  NPs would be precipitated on the surface of  $Ag_3PO_4$ during the reaction, which is consistent with the SEM observation results (Figure 3.5). The enhanced light absorption of the obtained porous  $Ag_3PO_4$  films together with their intrinsically high photoactivity suggest the potential of the porous films as highly efficient photoanodes for PEC water oxidation.



**Figure 3.13.** PEC water oxidation performance of the obtained  $Ag_3PO_4$  photoanodes: (a) LSV curves of all the photoanodes fabricated using  $Ag_3PO_4$  films obtained at different reaction times. The solid lines and dashed lines stand for the photocurrent densities and dark current densities, respectively; (b) Current density–time (*J*–*t*) curves of the porous  $Ag_3PO_4$  photoanodes fabricated using 1 h, 20 h and 24 h reaction times recorded during 1-h continuous PEC water splitting tests at an applied potential of 1.0 V *vs.* Ag/AgCl.



**Figure 3.14.** Monochromatic incident photon-to-electron conversion efficiency (IPCE) profiles of the obtained nearly non-porous 1-h Ag<sub>3</sub>PO<sub>4</sub> photoanode and highly porous 20-h Ag<sub>3</sub>PO<sub>4</sub> photoanode measured at 1 V *vs.* Ag/AgCl.



**Figure 3.15.** The dependence of  $J_{\text{light}} - J_{\text{dark}}$  (*i.e.* photocurrent density – dark current density) on the applied bias for all the porous Ag<sub>3</sub>PO<sub>4</sub> microcrystalline photoanodes obtained after different reaction times. These curves are derived from Figure 3.13a.

Figure 3.13a displays the applied bias-dependent photocurrent densities of all the fabricated Ag<sub>3</sub>PO<sub>4</sub> photoanodes together with the dark current densities. It is clear that as the reaction time increased from 1 h to 20 h, the photocurrent densities of the Ag<sub>3</sub>PO<sub>4</sub> photoanodes increased. The 20-h Ag<sub>3</sub>PO<sub>4</sub> photoanodes presented a photocurrent density of about 4.32 mA/cm<sup>2</sup> at 1 V vs. Ag/AgCl, which is almost three times higher than that of the 1-h  $Ag_3PO_4$  photoanode (1.48 mA/cm<sup>2</sup>). This enhancement could be attributed to the boosted active surface areas and also light capture ability of the highly porous samples than non-porous samples. The higher incident photon-to-electron conversion efficiency (IPCE) profiles of the 20-h Ag<sub>3</sub>PO<sub>4</sub> photoanode than 1-h Ag<sub>3</sub>PO<sub>4</sub> photoanode, as shown in Figure 3.14, suggest the higher external quantum efficiencies (EQEs) of the porous photoanodes than non-porous ones, which have made contributions to the higher photocurrent densities in Figure 3.13a. As has been discussed before, the decrease in the amount of Ag<sup>0</sup> NPs on the surface of Ag<sub>3</sub>PO<sub>4</sub> as the reaction time increased may also have a positive effect on the photocurrent density. The Ag<sup>0</sup> NPs at Ag<sub>3</sub>PO<sub>4</sub>/electrolyte interfaces could easily attract photoelectrons due to their superior electron affinity and serve as the recombination centers for the photogenerated electron/hole pairs and, therefore, lead to the loss of photoholes entering water oxidation reaction.

It could be observed from Figure 3.13a that the dark current densities of Ag<sub>3</sub>PO<sub>4</sub> photoanodes also increased gradually together with the photocurrent densities. This phenomenon suggests that Ag<sub>3</sub>PO<sub>4</sub> also have electrocatalytic activity. Therefore, the dependences of  $J_{\text{light}} - J_{\text{dark}}$  values (*i.e.* photocurrent density minus dark current density) on the applied bias for all the Ag<sub>3</sub>PO<sub>4</sub> photoanodes have been calculated to verify the contribution of photocatalytic activity in the as-recorded photocurrent densities. Figure 3.15 shows that the  $J_{\text{light}} - J_{\text{dark}}$  values of different Ag<sub>3</sub>PO<sub>4</sub> photoanodes give the same

trend as that of Figure 3.13a, thus confirming that photocatalytic effect has indeed played a dominant role in generating the photocurrent. The slight decrease in the photocurrent density from the 20-h to 24-h Ag<sub>3</sub>PO<sub>4</sub> photoanodes might be originated from the increased interfacial charge transfer impedance between Ag<sub>3</sub>PO<sub>4</sub> surfaces and the electrolyte solution. This effect is evidenced by the EIS Nyquist plots exhibited in Figure 3.16. On one hand, the relatively low interfacial charge transfer resistance ( $R_{CT}$ ) and thus high conductivity of all these Ag<sub>3</sub>PO<sub>4</sub> photoanodes could be regarded as one reason for the high current densities observed in Figure 3.13a. On the other hand, the increased interfacial  $R_{CT}$  derived from the increased porosity induced an increase in the density of non-conductive voids in the Ag<sub>3</sub>PO<sub>4</sub> films obtained with longer reaction time might have resulted in the decreased photocurrent density from the 20-h to 24-h Ag<sub>3</sub>PO<sub>4</sub> porous photoanodes.

Besides the high photocurrent density, it is notable that the porous Ag<sub>3</sub>PO<sub>4</sub> photoanodes obtained with prolonged reaction times ( $\geq 20$  h) also demonstrated better stability than non-porous ones, which could be known clearly from a 1-h successive PEC water splitting test carried out at 1 V *vs.* Ag/AgCl (Figure 3.13b). The promotion in stability of the porous Ag<sub>3</sub>PO<sub>4</sub> photoanodes may be caused by the pores inside the Ag<sub>3</sub>PO<sub>4</sub> layer because the porous structures could help to release the stress at the near-surface region so that could help to keep them more stable than non-porous Ag<sub>3</sub>PO<sub>4</sub> according to previous reports.[19–21] SEM images of the 1-h and 20-h Ag<sub>3</sub>PO<sub>4</sub> photoanodes after experiencing the 1-h PEC water splitting tests, shown in Figure 3.17, clearly confirm that the continuous water oxidation reaction on the surface of the 1-h photoanode has caused a strong surface erosion of the originally non-porous Ag<sub>3</sub>PO<sub>4</sub> film (Figure 3.17a–b), leading to a non-negligible loss of active materials and decrease



**Figure 3.16.** Nyquist plots of electrochemical impedance spectroscopy (EIS) recorded under dark conditions of the fabricated 1-h, 20-h and 24-h Ag<sub>3</sub>PO<sub>4</sub> microcrystalline photoanodes. The insert shows the equivalent circuit which consists of a serial resistor ( $R_S$ ) and a resistor–capacitor circuit. The resistor–capacitor circuit models the parallel combination of the interfacial charge transfer resistance ( $R_{CT}$ ) and the constant phase element (CPE), attributed to the interface between the surface of photoanode and the electrolyte.

of the photocurrent density. However, the highly porous structure of the 20-h Ag<sub>3</sub>PO<sub>4</sub> photoanode (Figure 3.17c–d), on the contrary, remained essentially the same after the 1-h stability test, ensuring the maintenance of the high photocurrent density. The 20-h porous Ag<sub>3</sub>PO<sub>4</sub> photoanode, moreover, displayed considerable stability even when the stability test was further extended to 5 h, with its surface morphology, phase purity, and chemical surface state maintaining essentially unchanged (Figure 3.18). The Faradaic efficiencies of the 5-h successive water splitting reaction using the 20-h porous Ag<sub>3</sub>PO<sub>4</sub> film as the photoanode are all calculated to be ~100% based on the H<sub>2</sub> evolution amount

detected using a gas chromatography (GC) setup (Figure 3.19), which further confirm the stability of the 20-h porous Ag<sub>3</sub>PO<sub>4</sub> photoanode in relatively long-time reactions.



**Figure 3.17.** Top-view SEM images for the (a–b) 1-h as well as the (c–d) 20-h  $Ag_3PO_4$  photoanodes after experiencing 1-h continuous PEC water splitting tests at an applied bias of 1.0 V *vs.* Ag/AgCl. The insert in panel (b) shows an enlarged view of the surface of the 1-h  $Ag_3PO_4$  sample after the stability test.

Regarding the curve shape in Figure 3.13b, it can be known that the photocurrent densities of both the 20-h and 24-h  $Ag_3PO_4$  porous microcrystalline film photoanodes first increase and then decrease during the stability test (*i.e.* the brown and purple curves). Possibly, there are two reasons which might have resulted in this slight increase at first:



**Figure 3.18.** (a) Current density–time (*J*–*t*) curve recorded in the 5-h successive PEC water splitting test using the 20-h porous  $Ag_3PO_4$  photoanode at 1.0 V *vs.* Ag/AgCl, and representative SEM image (b), XRD pattern (c), and XPS spectra of Ag 3d (d), P 2p (e) and O 1s (f) region of the 20-h Ag\_3PO\_4 photoanode after the 5-h stability test. The binding energy in panels (d–f) is calibrated by using the contaminant C 1s peak at 284.8 eV as the reference.



**Figure 3.19.** Gas chromatography collection of the  $H_2$  gas generated during the 5-h successive stability test (Figure 3.18a) using the 20-h porous Ag<sub>3</sub>PO<sub>4</sub> microcrystalline photoanode at an applied bias of 1.0 V *vs.* Ag/AgCl, and the correspondingly calculated Faradaic efficiencies.

1. This increase in partial photocurrent density might be caused by photoexcitation and photoactivation of the porous  $Ag_3PO_4$  photoanodes and the subsequent participation of photoexcited holes in the anodic water oxidation reaction at the interface.[22]

2. Considering that  $Ag_3PO_4$  is an *n*-type semiconductor with both a direct bandgap of ~2.43 eV and an indirect bandgap of ~2.36 eV, it is possible that the indirect bandgap excitation of  $Ag_3PO_4$  may have caused non-negligible thermal effect induced by the lattice vibration. As a result, the slight increase of the local temperature at the surface of the photoanode might also have brought about the promotion in catalytic activity, and thus the photocurrent density.

However, in the case of the nearly non-porous 1-h Ag<sub>3</sub>PO<sub>4</sub> photoanode, in spite of these favorable effects, the non-negligible loss of active materials revealed by the SEM

images (Figure 3.17a–b) has still played a dominant role in causing the decrease in the photocurrent density, and thus no increase in the photocurrent density could be observed (*i.e.* the gray curve in Figure 3.13b). Overall, the high photocurrent density as well as improved stability manifested by the 20-h  $Ag_3PO_4$  photoanodes suggest great potential for highly efficient PEC energy conversion.

#### **3.4.** Conclusions

In this chapter, we prepared large-area  $(4.5 \times 4.5 \text{ cm}^2)$  and solution-processed Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films with tunable porosity controlled by the reaction time as novel photoanodes for PEC water splitting. Enhanced light absorption was demonstrated for the microcrystalline Ag<sub>3</sub>PO<sub>4</sub> films with an optimized porosity obtained for a prolonged reaction time ( $\geq 20$  h). This enhanced light absorption results in improved water oxidation performance of the Ag<sub>3</sub>PO<sub>4</sub>–based photoanodes. Particularly, the highly porous 20-h Ag<sub>3</sub>PO<sub>4</sub> photoanode presented a photocurrent density of around 4.32 mA/cm<sup>2</sup> at 1.0 V *vs.* Ag/AgCl, which is nearly three times higher than that of the non-porous 1-h Ag<sub>3</sub>PO<sub>4</sub> photoanode (1.48 mA/cm<sup>2</sup>). In addition, superior stability of the 20-h porous Ag<sub>3</sub>PO<sub>4</sub> photoanode was revealed by the 1-h successive PEC water splitting experiment. Based on these results, it is believed these porous Ag<sub>3</sub>PO<sub>4</sub> microcrystalline photoanodes hold great potential to support efficient solar-to-fuel energy conversion and many other PEC applications.

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# 4. Overall summaries

In this thesis, solution-processed large-area  $(4.5 \times 4.5 \text{ cm}^2) \text{ Ag}_3\text{PO}_4$  microcrystalline films are first prepared *via* a room-temperature, air-exposed immersion reaction. The multiple effects of the native defects in Ag<sub>3</sub>PO<sub>4</sub>, *i.e.* V<sub>Ag</sub>, Ag<sub>i</sub>, and surface precipitation of Ag<sup>0</sup> NPs, are first demonstrated experimentally. By varying the addition amount of the oxidation agent H<sub>2</sub>O<sub>2</sub> during the synthetic reaction, Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films with different amounts of Ag<sup>0</sup> NPs on the surface are obtained. It is found that despite the favorable superior conductivity and plasmonic effect having improved the interfacial charge transfer and light absorption of the Ag<sub>3</sub>PO<sub>4</sub> film photoanode, the defective film with the largest amount of Ag<sup>0</sup> NPs on the surface finally exhibits the lowest current density, due to the disadvantageous effects involving the decreased crystalline quality, easier decomposition of active material of Ag<sub>3</sub>PO<sub>4</sub>, decreased electrochemical surface area with Ag<sup>0</sup> NPs on the surface serving as the recombination centers. As a result, even the Ag<sub>3</sub>PO<sub>4</sub> microcrystalline film having achieved the highest photocurrent density fails to demonstrate considerable stability. These results further highlight the significance of fundamentally reducing the native defects of Ag<sub>3</sub>PO<sub>4</sub> towards enhanced water oxidation stability and efficiency of the microcrystalline film photoanodes.

Typically, the post-treatment of vacuum annealing at different temperatures (*i.e.* 100, 200, 300, 400 and 500 °C) has been investigated as a strategy for solving the problem of native defects and low light absorption of Ag<sub>3</sub>PO<sub>4</sub>. It is found that below 400 °C is the |Recovery| stage for Ag<sub>3</sub>PO<sub>4</sub> microcrystals when defect amount decreases and therefore crystallinity increases. From 400 °C, the |Recrystallization| stage and also precipitation of Ag<sub>3</sub>PO<sub>4</sub> bulk crystals start, and afterwards till 500 °C is the |Recrystallization + Grain

growth| stage for Ag<sub>3</sub>PO<sub>4</sub>, when more faceted and larger bulk crystals of Ag<sub>3</sub>PO<sub>4</sub> precipitate extensively. In this series of Ag<sub>3</sub>PO<sub>4</sub> film photoanodes, the same variation in  $[J_{\text{light}} - J_{\text{dark}}]$  vs. *E* curves are found as the light absorption profiles, *i.e.* 400 °C > 500 °C > 300 °C > 200 °C > non-annealed > 100 °C-annealed Ag<sub>3</sub>PO<sub>4</sub> film, indicating that light absorption, especially the bandgap absorption could be the dominant factor to the PEC performance. Particularly, much improved PEC water oxidation stability and efficiency of the 400 °C-annealed Ag<sub>3</sub>PO<sub>4</sub> film photoanode are demonstrated compared with the non-annealed Ag<sub>3</sub>PO<sub>4</sub> film, suggesting that the methodology of vacuum annealing could be utilized for improving the PEC performance of Ag<sub>3</sub>PO<sub>4</sub> film photoanodes because of its ability to reduce the native defects like  $V_{\text{Ag}}$ , remove the Ag<sup>0</sup> NPs on the surface, and increase the bandgap light absorption of these Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films.

Next, from the viewpoint of advanced structural design, in consideration that the usage of microcrystalline film photoanode may cause the problem of decreased active surface area despite its relatively higher stability than nanocrystalline powder-deposited photoanode, the possibility of introducing porous structures into these microcrystalline Ag<sub>3</sub>PO<sub>4</sub> film photoanodes is also investigated towards further improved photoactivity and stability. Particularly, it is demonstrated that the porous degree of the as-obtained Ag<sub>3</sub>PO<sub>4</sub> films could be tuned by changing the reaction time from 1 h to 2 h, 5 h, 9 h, 20 h and 24 h, with the usage of specific reaction parameters. These porous Ag<sub>3</sub>PO<sub>4</sub> films, in spite of leaving the native defects not recovered and Ag<sup>0</sup> NPs on the surface not removed, finally could demonstrate improved surface areas and light absorption, owing to which improved PEC performance has also been achieved. Particularly, the optimized highly porous 20-h Ag<sub>3</sub>PO<sub>4</sub> microcrystalline film photoanode shows superior water oxidation efficiency with a photocurrent density of 4.32 mA/cm<sup>2</sup> which is almost three

times higher than that of non-porous 1-h film (1.48 mA/cm<sup>2</sup>) at 1.0 V vs. Ag/AgCl, and is also higher than the 400 °C-annealed non-porous film photoanode.

Overall, it could be concluded that: a) The instability of the Ag<sub>3</sub>PO<sub>4</sub> microcrystalline films is originated from the native defects involving  $V_{Ag}$  and Ag<sup>0</sup> NPs precipitated on the surface, thus the methodology of vacuum annealing, typically at 400 °C, turns out to be more favorable for improved stability because it can fundamentally solve, or at least mitigate the various problems induced by native defects; b) Light absorption is a crucial factor to the PEC performance, whereas the surface area which largely determines the final process of PEC water oxidation – surface redox reaction is also important, thus the porous Ag<sub>3</sub>PO<sub>4</sub> microcrystalline film photoanodes can achieve higher efficiency in spite of leaving the defect and absorption issues not solved; c) Trying to make a combination of these two techniques, *i.e.* taking both the fundamental properties of the material (*e.g.*, defects and absorption) and the advanced structural design (*e.g.*, larger surface area) into concern, may provide further possibility towards more stable and efficient Ag<sub>3</sub>PO<sub>4</sub> film photoanodes for PEC water oxidation.

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