Cover Sheet

This file is the final accepted manuscript for the following paper:

ANISOtime: Traveltime Computation Software for Laterally Homogeneous, Transversely Isotropic, Spherical Media By Kensuke Konishi, Anselme F. E. Borgeaud, Kenji Kawai, Robert J. Geller Publisher: Seismological Society of America Published: 14 July 2021 Seismological Research Letters (2021) 92 (6): 3811–3820. DOI: https://doi.org/10.1785/0220200306

Access to ANISOtime software

ANISOtime can be downloaded from

https://github.com/UT-GlobalSeismology/anisotime.

Downloadable executable versions are available for Windows, macOS, and Unix/Linux; the source code can also be downloaded. A user guide can be downloaded from this site, and is also embedded in the software.

| 1 | ANISOtime: Traveltime computation software |
|----|---|
| 2 | for laterally homogeneous, transversely isotropic, |
| 3 | spherical media |
| 4 | Kensuke Konishi ¹ , Anselme F. E. Borgeaud ^{1,2} , Kenji Kawai ² , and |
| 5 | Robert J. Geller ^{*2} |
| 6 | ¹ Institute of Earth Sciences, Academia Sinica, 11529 Taipei, |
| 7 | Taiwan. |
| 8 | ² Department of Earth and Planetary Science, School of Science, |
| 9 | University of Tokyo, Hongo 7-3-1, Bunkyo-ku, Tokyo 113-0033 |
| 10 | Japan. |
| | |

May 7, 2021

Declaration of Competing Interests: The authors acknowledge there are no conflicts of interest recorded.

12 ABSTRACT

11

- 13 Software packages for computing seismic traveltimes and raypaths in an isotropic,
- ¹⁴ spherically symmetric, Earth model are well known and widely used. However,

^{*}Corresponding Author. email: bob@eps.s.u-tokyo.ac.jp

even though the theory for transversely isotropic (TI), spherically symmetric, 15 models has been known since the late 1960s, readily available programs for 16 traveltime calculations are restricted to isotropic models. We have developed 17 a new software package, ANISOtime, for computing seismic traveltimes and 18 raypaths in laterally homogeneous, transversely isotropic (TI), spherical media. 19 This package calculates traveltime tables for both immediate and subsequent 20 use. ANISOtime has both graphical user interface (GUI) and command-line 21 interface (CLI) modes. The package is available for free public download. As it 22 offers cross-platform compatibility through Java 8, it runs on Windows, macOS, 23 and Unix/Linux. 24

²⁵ INTRODUCTION

Traveltime computation is required widely in seismology. Although the Earth is laterally heterogeneous, the starting point for much research and teaching is traveltime computation for a laterally homogeneous, isotropic, model. The *TauP Toolkit* (Crotwell *et al.*, 1999) is well known as a reliable and readily usable public software package for traveltime and raypath computations for arbitrary spherically symmetric, isotropic velocity models.

The importance of anisotropy in seismological research is steadily increasing. For example, inferring anisotropy in the mantle allows inferences to be drawn about the direction of mantle flow. Inferring anisotropy is also an important research topic in exploration seismology. Use of *ANISOtime* can contribute to such research. It can also be used in introductory seismology courses to give students some exposure to basic concepts of anisotropy. A sample homework exercise is available from the GitHub page for *ANISOtime*.

As is well known (Love, 1927; Crampin, 1981), the most general anisotropic elastic medium has 21 independent elastic constants, while an isotropic elastic

medium has just two, λ (or κ) and μ . In this paper we consider transversely 41 isotropic (TI) media with a vertical symmetry axis (sometimes called VTI me-42 dia). The basic theory for such media is well known (Vlaar, 1968, 1969; Wood-43 house, 1981), but to our knowledge, no readily available public software package 44 can handle traveltime calculations for transversely isotropic media. 45

We have developed a free public software package, ANISOtime, for making 46 traveltime calculations for a spherically symmetric, TI, medium. ANISOtime 47 takes advantage of the cross-platform benefits of the Java language and has both 48 graphical user interface (GUI) and command line interface (CLI) modes. 49

THEORY 50

As noted above, the basic theory for computing traveltimes in a transversely 51 isotropic medium with a vertical symmetry axis is well known, and ANISOtime 52 uses these results. The three basic types of body-waves in such a medium are 53 called pseudo-P, pseudo-SV, and SH, respectively. The "pseudo" for the first and 54 second wave types is because they are not strictly longitudinal and transverse. 55 To make this paper self-contained, we present a derivation of the theory in 56 Text S1 of the supplemental material. This may be a useful supplementary text 57 for use in introductory courses. All variables used in the main body of this paper 58 are listed in Table 1 in order of their appearance. For purposes of the theoretical 59 derivations the ray parameter p has units of s/radian, the epicentral distance 60 Δ is in radians, and all other variables are in SI units, but for convenience the 61 density ρ is in g/cm³, the radius r is in km, the velocities are in km/s, the

- epicentral distance Δ is in degrees, and the ray parameter p is in s/degree in 63 the input to ANISOtime. 64
- 65

62

ANISOtime computes traveltime T and epicentral distance Δ for a spherical

66 model as follows:

$$\Delta(p) = \int q_{\Delta}(p, r) \, dr \tag{1}$$

$$T(p) = \int q_T(p, r) \, dr, \qquad (2)$$

where the kernels are defined in eqs. (S72)–(S77) in the supplemental material. The above integrations are computed by Simpson's rule (with the exception of the layer at the turning point, see below) for a given integral mesh $(r_1, r_2, ..., where r_i < r_{i+1})$ as follows:

$$\Delta(p) = \sum_{i} \int_{r_i}^{r_{i+1}} q_\Delta(p, r) \, dr \tag{3}$$

$$T(p) = \sum_{i} \int_{r_i}^{r_{i+1}} q_T(p, r) \, dr.$$
(4)

The computational mesh can be arbitrary. Our default mesh spacing is 1 km. 71 In order to perform the integration accurately, special care must be taken in 72 handling the integration near the turning point of a raypath, where the kernel 73 $q_{\tau} \rightarrow 0$, and the integrands become singular but integrable. In ANISOtime, the 74 integration for the interval bounded by the turning point is computed following 75 Jeffreys and Jeffreys (1956, p. 288–290), as is also done by Woodhouse (1981). 76 The details of this procedure are discussed in Text S2 of the supplemental ma-77 terial. 78

79 Earth model

In order to compute Δ and T using eqs. (1) and (2), the density ρ and five independent elastic constants for solid TI media (A, C, F, L, N) are required. ⁸² The constitutive relation for a TI medium is as follows:

$$\begin{pmatrix} \sigma_{xx} \\ \sigma_{yy} \\ \sigma_{zz} \\ \sigma_{xz} \\ \sigma_{yz} \\ \sigma_{xy} \end{pmatrix} = \begin{pmatrix} A & H & F & & \\ H & A & F & & \\ F & F & C & & \\ & & L & & \\ & & & L & \\ & & & L & \\ & & & N \end{pmatrix} \begin{pmatrix} e_{xx} \\ e_{yy} \\ e_{zz} \\ 2e_{xz} \\ 2e_{yz} \\ 2e_{yy} \end{pmatrix},$$
(5)

84 where

$$H = A - 2N. \tag{6}$$

85

83

The five independent elastic constants, A, C, F, L, and N, must be chosen so that the strain energy density is positive definite. This can be verified for any particular set of elastic constants by computing the principal minors of the strain energy tensor (the matrix in eq. 5). For an isotropic medium the relations between the above five elastic constants and λ and μ are as follows:

$$\lambda + 2\mu = A = C \tag{7}$$

$$\mu = L = N \tag{8}$$

$$\lambda = F = H. \tag{9}$$

86

The "PolynomialStructure" form (see Text S3 of the supplemental material) is one of two permissible formats to specify the input parameters to ANISOtime, ρ , V_{PV} , V_{PH} , V_{SV} , V_{SH} , and η , as cubic functions of radius. This allows an analytical computation of the turning point of a raypath by solving a cubic equation. The "PolynomialStructure" form also specifies the shear, and bulk attenuation coefficients Q_{μ} , and Q_{κ} , respectively. The attenuation coefficients are not used by *ANISOtime*, but should be included in the input parameter file so that users can use the same input file for *ANISOtime* and for the Direct Solution Method waveform computation software (DSM; Kawai et al., 2006). The standard definitions of V_{PV} , etc., are used (e.g., Panning and Romanowicz, 2006):

98

$$V_{PH} = \sqrt{A/\rho} \tag{10}$$

$$V_{PV} = \sqrt{C/\rho} \tag{11}$$

$$V_{SH} = \sqrt{N/\rho} \tag{12}$$

$$V_{SV} = \sqrt{L/\rho} \tag{13}$$

$$\eta = \frac{F}{A - 2L} \tag{14}$$

In ANISOtime, the anisotropic and isotropic Preliminary Reference Earth Mod-99 els (PREM: Dziewonski and Anderson, 1981) and AK135 model (Kennett et al., 100 1995) are embedded in the program in "PolynomialStructure" form. When us-101 ing ANISOtime in CLI mode, these models can be called using the argument 102 values "prem," "iprem," and "ak135." We also support an input model for-103 mat called "Named Discontinuity" (see Text S4 of the supplemental material) 104 in the TauP toolkit (Crotwell et al., 1999). At present, all models must have 105 an Earth-like structure with a "mantle" (i.e., a solid outer region), underlain 106 by an "outer core" (i.e., a liquid region), with an "inner core" (i.e., another 107 solid region) at the center. Each of these regions can be arbitrarily vertically 108 heterogeneous; i.e., the solid outer region ("mantle") can actually consist of a 109 crust underlain by an upper mantle, transition zone, etc. The "outer core" must 110 have zero shear modulus $(V_{SH} = V_{SV} = 0)$, while the other regions must have 111

strictly positive velocities and density. Models must specify the radius of the core-mantle boundary (CMB) and the inner-core boundary (ICB). Any model that satisfies the above conditions can be used.

¹¹⁵ Phases and distance

The rules for possible phase names follow those of the *Taup Toolkit* (Crotwell *et al.*, 1999). Symbols that describe wave types and interactions are listed in Table 2.

For most of the phases, the naming follows conventions used in global seis-119 mology. One distinction is that phase names must specify all the individual 120 branches of the raypath (together with interactions at internal boundaries). For 121 instance, the ScS2 phase, which bounces two times at the core-mantle bound-122 ary (tracing a 'W' in the mantle), is named 'ScSScS.' Another distinction is 123 for phases that reflect at internal discontinuities other than the core-mantle 124 boundary or the inner-core, which must be specified using a 'v' for a topside 125 reflection, or a hat '^' for an underside reflection, followed by the depth of the 126 internal discontinuity. For instance, an S phase with a topside reflection at the 127 670 km discontinuity in PREM is named 'Sv670S.' 128

For phases for which all branches are S phases (e.g., S, ScS, ScSScS, SS), *ANISOtime* allows the user to specify the polarization (SH or SV). The polarization is not specified in the phase name, but rather by using the options '-SV' and '-SH' in the command line (see section "CLI," below), or by switching the polarization in the GUI. The S branches for phases that include a P branch (e.g., SKS, ScP) will always have only SV polarization, and an error will result if the user tries to use the '-SH' option with such phases.

The epicentral distance computed by *ANISOtime* is the distance along the Earth's surface for the whole raypath. Although raypaths can have an epicentral distance of 360° or larger, ANISOtime does not make computations for such
phases at present.

140 ABOUT ANISOtime

ANISOtime has both command line interface (CLI) and graphical user interface (GUI) modes. Both are similar to those used by the *TauP Toolkit*. There are some features which *TauP* can handle but *ANISOtime* cannot yet. We have no plans for any further development of the program, but will try to respond if there is a strong demand from users. *ANISOtime* automatically downloads and installs updates when there is a new release.

147 INSTALLATION

Java 8, or a more recent version, must be installed in order to run ANISOtime.
Users can verify their Java environment by accessing https://www.java.com/
en/download/installed8.jsp.

ANISOtime can be downloaded from https://github.com/UT-GlobalSeismology/anisotime. Downloadable executable versions are available for Windows, macOS, and Unix/Linux; the source code can also be downloaded. A user guide can be downloaded from this site, and is also embedded in the software.

156 CLI

¹⁵⁷ When launched with arguments, *ANISOtime* runs in CLI mode. For instance, to ¹⁵⁸ compute and output an image of an SH phase raypath from a seismic source at a ¹⁵⁹ depth of 500 km at an epicentral distance of 60° propagating in the (anisotropic) ¹⁶⁰ PREM model, the arguments are "-h 500 -deg 60 -ph S -mod prem -eps -o /path/to," which are identical to the arguments in TauP (except for the -eps
and -o options). This will return the traveltime and create an eps file of the
raypath in folder "/path/to."

¹⁶⁴ ¹⁶⁵ ¹⁶⁶ ¹⁶⁷ ¹⁶⁷ To obtain the traveltime and raypath for a pseudo-SV phase, the additional ¹⁶⁸ ¹⁶⁹ ¹⁶⁹ ¹⁶⁹ ¹⁷⁰ ¹⁷⁰ ¹⁷⁰ ¹⁷¹ ¹⁷¹ ¹⁷¹ ¹⁷¹ ¹⁷¹ ¹⁷² ¹⁷¹ ¹⁷² ¹⁷¹ ¹⁷² ¹⁷¹ ¹⁷² ¹⁷² ¹⁷² ¹⁷² ¹⁷² ¹⁷² ¹⁷² ¹⁷³ ¹⁷⁴ ¹⁷⁵ ¹⁷⁵

The full list of arguments can be obtained from the command "anisotime -help," and is given in Table 3.

174 RECORD SECTION

To draw a traveltime curve or create a record section, there must be sets of traveltime and epicentral distance values. Here is an example for the computation for pseudo-P-, pseudo-SV-, and SH-waves in PREM for the epicentral distance range $30^{\circ} \leq \Delta \leq 60^{\circ}$ with interval 5° for a source depth of 0 km: $\frac{176}{180}$ % anisotime -rs 30,60,5 -h 0 -ph P,S -mod prem -o /path/to

¹⁸² Other examples can be found in the user guide.

183 **GUI**

185

¹⁸⁴ When *ANISOtime* is launched without any arguments, it launches in GUI mode.

- 186 % anisotime
- ¹⁸⁸ The GUI has 2 computational modes:
- 189 1. Epicentral distance mode

190 2. Ray parameter mode.

In either mode, the user must select seismic phases, epicentral distance, structure model and the depth of the source. Note that other parameters can also be changed. In epicentral distance mode, the epicentral distance is specified and raypaths are computed for that value, while in ray parameter mode raypaths
are computed for the specified input ray parameter. Figs. 1 and 2 show results
for epicentral distance mode and ray parameter mode, respectively.

¹⁹⁷ END USER LICENSE AGREEMENT (EULA)

The software is licensed under the GNU General Public License (GPL), Ver-198 sion 3.0 (https://www.gnu.org/licenses/gpl-3.0.en.html.) After downloading 199 ANISOtime, users must accept an End User License Agreement (EULA) be-200 fore it can be launched. The main points of the EULA are that users agree 201 to comply with the GPL in the event they use parts or all of our software in 202 other works, and that they waive all possible claims in the event of problems 203 with the software. This brief description is purely informal; the sole and bind-204 ing agreement is that in the EULA itself. The authors welcome questions and 205 bug reports to ut-globalseis@googlegroups.com and will respond if possible, but 206 make no legally binding promise to do so. 207

208 RAYPATH CATALOG

When the ray parameter p is chosen, the raypath (for a given Earth structure) 209 is determined, and so are both the epicentral distance Δ and the traveltime T 210 for the raypath. Although the traveltime T and epicentral distance Δ can be di-211 rectly computed for a particular ray parameter, most users want the traveltime 212 T for a particular epicentral distance Δ . Since we cannot obtain the traveltime 213 T or ray parameter p directly from the epicentral distance Δ , we must first 214 find a ray parameter for a target epicentral distance Δ and then compute the 215 traveltime T for the ray parameter p. In many cases, users require many pairs 216 of T and Δ . In order to reduce the computational time, ANISOtime first com-217

²¹⁸ putes sets of ray parameters p and corresponding epicentral distances Δ (i.e., ²¹⁹ a catalog) for a given structure, so that it can look for the ray parameter p²²⁰ which gives an epicentral distance Δ chosen by the user. When a user computes ²²¹ traveltimes for a new structure, the catalog for the structure will automatically ²²² be stored and used for later calculations.

223 VALIDITY CHECKS

²²⁴ Comparison to TauP

To verify the accuracy of the ANISOtime package, we first compare it to TauP 225 for the case of isotropic PREM. Fig. 3 shows the difference of computed travel-226 times for the default phases (S, P, ScS, PcP, SKS, PKP, SKiKS, and PKiKP) 227 for a source at the Earth's surface. Fig. 3a shows the traveltime differences 228 when using the default PREM of TauP (for which the depth grid intervals are 229 between 15–100 km). Excluding the region around 25° , which corresponds to 230 S and P phase triplications, thetraveltime difference for all phases is within 231 0.06 s. This relatively large discrepancy is due to the fact that the default 232 accuracy of ANISOtime is better than that for TauP. Traveltime comparisons 233 using a higher accuracy calculation for TauP are shown in Fig. 3b. The higher 234 accuracy calculation is obtained by using an input layered velocity structure 235 with finer depth sampling (6.371 km) than that for the default isotropic PREM 236 in TauP. The traveltime differences for the higher-accuracy TauP calculations 237 are within 0.0116 s of *ANISOtime* for all default phases for the entire epicentral 238 distance range, and within 0.0016 s when excluding the regions near S wave 239 triplications (around 25°). This is more than one order of magnitude smaller 240 than the errors for the default PREM of TauP. In summary, ANISOtime and 241 TauP are in good agreement for highly accurate calculations. 242

243 Comparison to an analytical solution

An analytical solution can be found for the TI medium with the elastic constants
and density defined below:

$$A(r) = A_0 r^2 \tag{15}$$

$$C(r) = C_0 r^2 \tag{16}$$

$$L(r) = L_0 r^2 \tag{17}$$

$$N(r) = N_0 r^2 \tag{18}$$

$$F(r) = F_0 r^2 \tag{19}$$

$$\rho = \text{constant.}$$
(20)

We present a derivation in Text S5 of the supplemental material. As far as 246 we know, this result is new for the spherically symmetric case, but it follows 247 straightforwardly from well known results for the Cartesian case (Shearer and 248 Chapman, 1988; Červený, 1989). We use this analytical solution to check the 249 accuracy of the numerical traveltime integrals. Note that the medium defined 250 in eqs. (15)-(20) has velocities that decrease linearly with depth. This is not a 251 physically realistic model, but it is the only one for which an analytic solution 252 is available for checking the numerical computations. 253

For the above medium, the traveltime between radius r_1 and r_2 is given by

$$T(p) = \int_{r_1}^{r_2} q_T(r, p) \, dr = q_{T0} \ln \frac{r_2}{r_1}, \tag{21}$$

where q_{T0} depends on the polarization (SH, pseudo-SV, or pseudo-P) and is given in Text S5 of the supplemental material.

²⁵⁷ We compute the difference between the traveltime for the ScS (SH), ScS

(pseudo-SV), and pseudo-PcP phases computed using ANISOtime, and that 258 computed using the analytical solution eq. (21) for the medium whose properties 259 are defined in eqs. (15)-(20) with $A_0 = 7.55 \cdot 10^{-5}, C_0 = 7.12 \cdot 10^{-5}, F_0 =$ 260 $2.63 \cdot 10^{-5}, L_0 = 2.46 \cdot 10^{-5}, N_0 = 2.61 \cdot 10^{-5}$, where A_0, C_0, F_0, L_0 , and N_0 are 261 in $Pa/m^2 = kg m^{-3} s^{-2}$, and $\rho = 10 kg/m^3$. Note that the units of $A_0 \dots$ are 262 Pa/m^2 rather than Pa, because, as shown by eqs. (15)-(19), the elastic constants 263 are obtained by multiplying by r^2 , which has units of m^2 . The relative errors 264 (defined below in eq. 22) using a constant integration mesh with 1 km spacing 265 are $2.4490 \cdot 10^{-11}$, $2.4489 \cdot 10^{-11}$, and $2.4492 \cdot 10^{-11}$, for the ScS (SH), ScS 266 (pseudo-SV), and pseudo-PcP phases, respectively. 267

relative error =
$$\frac{T_{ANISOtime} - T_{analytical}}{T_{analytical}}$$
(22)

²⁶⁸ Comparison to full-wave theory

To further test the accuracy of ANISOtime, we compare the traveltime pre-269 dicted by ANISOtime to the arrival time on waveforms computed using full-270 wave theory for realistic isotropic and transversely isotropic media. The syn-271 thetic waveforms are computed using the Direct Solution Method (Kawai et272 al., 2006) up to 2 Hz for an event at depth 571.3 km. The Earth models used 273 are PREM, isotropic PREM, and MIASP91ANI, a modified version of IASP91 274 with smoothed upper mantle discontinuities and $1\% V_{SV}$ and $3\% V_{SH}$ increases, 275 respectively, in the lowermost 250 km of the mantle (see Fig. 4). In order to 276 avoid the effects of velocity dispersion, we set $Q_{\mu} = 5000$, and $Q_{\kappa} = 57823$, i.e., 277 essentially an elastic medium, to compute the synthetics using the DSM (which 278 requires non-zero values for the anelastic parameters). Velocity dispersion re-279 duces the velocity for frequencies lower than a reference frequency (typically 280 1 Hz), and would delay the arrival of seismic pulses on the synthetics, thereby 281

inducing disagreement with ANISOtime. The waveforms are low-pass filtered
with a corner frequency of 0.25 Hz.

The synthetics and traveltime curves are shown in Figs. 5, 6, 7 and 8. Fig. 5 284 shows a record section and traveltime curves for the SH (transverse component) 285 and pseudo-SV (radial component) phases that sample the D'' layer for the 286 model in Fig. 4c. The onset times of the SH and pseudo-SV phases, and the 287 triplicated arrivals due to the D'' discontinuity, are well reproduced by ANISO-288 time. Figs. 6, and 7 show a larger set of phases on three-component synthetics 289 for isotropic PREM (Fig. 4a), and closeups of some specific phases for clarity, 290 respectively. Fig. 8 shows record sections for the S and pseudo-P phases that 291 sample the upper mantle and mantle transition zone for the (anisotropic) PREM 292 model (Fig. 4b). 293

We quantify the agreement between the full-wave synthetics and ANISO-294 time in Fig. 9. Fig. 9 shows the difference between traveltimes predicted using 295 ANISOtime, and manually-picked onset times for the direct SH and pseudo-P 296 phases and waveforms for the isotropic (Fig. 6) and anisotropic (Fig. 8) PREM 297 models. Traveltime discrepancies are within 0.2 s (except for 5 points for the SH 298 phase at small epicentral distances, where the onset time is not clear because 299 of triplicated phases), with average differences for the SH phase of -0.06 s, and 300 -0.03 s, and -0.07 s, and -0.07 s for the pseudo-P phase, for the isotropic 301 and anisotropic PREM models, respectively. We note that we only picked the 302 SH and pseudo-P phases, since these are direct phases. Other later phases (or 303 the direct SV phase) have precursors before the onset of the main phases (due 304 to internal, e.g., crustal, reflections, or S-to-P conversions); this makes it dif-305 ficult to pick precise onset times. However, we expect later pseudo-P as well 306 as pseudo-SV phases to have similar levels of agreement, since later phases are 307 combinations of S and P branches, and the equation for the traveltime of the 308

pseudo-SV phase is the same as that for the pseudo-P phase, except for a difference in sign in one term (see eq. S69 in the supplemental material). For these
later phases, the agreement with *ANISOtime* can be visually checked in Figs. 5,
7, and 8.

313 DATA AND RESOURCES

We used no data. Parameters for the three Earth models embedded in ANISO-314 time (PREM, isotropic PREM, and AK135) are available from IRIS Data Ser-315 vices http://ds.iris.edu/ds/products/emc-referencemodels/. The supplemental 316 material consists of the following five Texts. Text S1: Theory for laterally ho-317 mogeneous, transversely isotropic, media; Text S2: Integration near the turning 318 point; Text S3: "PolynomialStructure" file; Text S4: "Named Discontinuity" 319 file; Text S5: Analytical solution for a spherically symmetric, TI, medium with 320 constant velocity gradient. 321

322 ACKNOWLEDGMENTS

This research was partly supported by grants from the Japan Society for the Promotion of Science (Nos. 16K05531, 15K17744, 15H05832, and 18K03797). We thank Jeroen Ritsema, the other referee, and the associate editor for their constructive comments on the original manuscript.

327 **REFERENCES**

328

Borgeaud, A. F. E., K. Konishi, K. Kawai, and R. J. Geller (2016). Finite frequency effects on apparent S-wave splitting in the D" layer: comparison ³³¹ between ray theory and full-wave synthetics, *Geophys. J. Int.* 207, 12–28, doi:
³³² 10.1093/gji/ggw254.

Červený, V. (1989). Ray tracing in factorized anisotropic inhomogeneous
media, *Geophys. J. Int.* 99, 91–100, doi: 10.1111/j.1365-246X.1989.tb02017.x.
Crampin, S. (1981). A review of wave motion in anisotropic and
cracked elastic-media, *Wave Motion* 3, 343–391. https://doi.org/10.1016/01652125(81)90026-3

Crotwell, H. P., T. J Owens, and J. Ritsema (1999). The TauP toolkit: flexible seismic travel-time and ray-path utilities, *Seismol. Res. Lett.* **70**, 154– 160, doi: 10.1785/gssrl.70.2.154.

Dziewonski, A. M., and D. L. Anderson (1981). Preliminary reference
Earth model, *Phys. Earth Planet. Inter.* 25, 297–356, doi: 10.1016/00319201(81)90046-7.

Jeffreys, H., and B. S. Jeffreys (1956). *Methods of Mathematical Physics*, 345 3rd ed., Cambridge University Press, doi:10.1017/CBO9781139168489.

Kawai, K., N. Takeuchi, and R.J. Geller (2006). Complete synthetic seismograms up to 2 hz for transversely isotropic spherically symmetric media, *Geophys. J. Int.* 164, 411–424, doi: 10.1111/j.1365-246X.2005.02829.x.

Kennett, B. L. N., and E. R. Engdahl (1991). Travel times for global
earthquake location and phase identification, *Geophys. J. Int.* 105, 429–465,
doi: doi.org/10.1111/j.1365-246X.1991.tb06724.x.

Kennett, B. L. N., E. R. Engdahl, and R. Buland (1995). Constraints on
seismic velocities in the Earth from traveltimes, *Geophys. J. Int.* 122, 108–124,
doi: 10.1111/j.1365-246X.1995.tb03540.x.

Love, A. E. H. (1927). A Treatise on the Mathematical Theory of Elasticity 4th ed., *Cambridge University Press*.

³⁵⁷ Panning, M., and B. Romanowicz (2006). A three-dimensional radially

- anisotropic model of shear velocity in the whole mantle, Geophys. J. Int. 167,
- 359 361-379, doi: 10.1111/j.1365-246X.2006.03100.x.
- Shearer P. M., and C. H. Chapman (1988). Ray tracing in anisotropic
 media with a linear gradient, *Geophys. J.* 94, 575–580, doi: 10.1111/j.1365246X.1988.tb02277.x.
- Vlaar, N. J. (1968). Ray theory for an anisotropic inhomogeneous elastic
 medium, Bull. Seism. Soc. Am. 58, 2053–2072.
- Vlaar, N. J. (1969). Rays and travel times in a spherical anisotropic earth,
- ³⁶⁶ Bull. Seism. Soc. Am. **59**, 1051–1060.
- ³⁶⁷ Woodhouse, J. H. (1981). A note on the calculation of travel times in a
- transversely isotropic Earth model, Phys. Earth Planet. Inter. 25, 357–359,
- doi: 10.1016/0031-9201(81)90047-9.
- 370 K. Konishi, and A.F.E. Borgeaud
- 371 Institute of Earth Sciences
- 372 Academia Sinica
- 128 Academia Road, Section 2, Nangang, Taipei 11529
- 374 Taiwan
- ³⁷⁵ A.F.E. Borgeaud, K. Kawai, and R.J. Geller
- 376 Department of Earth and Planetary Science
- 377 School of Science
- 378 University of Tokyo
- 379 Hongo 7-3-1, Bunkyo-ku, Tokyo 113-0033
- 380 Japan

Table 1: Variables used in main body of paper (in order of appearance).

| Variable | Meaning | Units | Where defined |
|----------------------------------|--|---------------|----------------------|
| T | traveltime | s | before eq. (1) |
| Δ | epicentral distance | radian | before eq. (1) |
| p | ray parameter for spherical Earth | s/radian | eq. (S70) |
| r | radius in spherical polar coordinates | km | before eq. (1) |
| $q_{\Delta}(p,r)$ | kernel for epicentral distance | radian/km | eqs. (S75)–(S77) |
| $q_T(p,r)$ | kernel for traveltime | s/km | eqs. $(S72)-(S74)$ |
| A, C, F, L, N | elastic constants for VTI medium | N/m^2 | eqs. $(5)-(9)$ |
| $V_{PH}, V_{PV}, V_{SH}, V_{SV}$ | $\begin{cases} \text{velocity-like quantities} \\ \text{used as input to } ANISO time \end{cases}$ | $\rm km/s$ | eqs. (10) – (13) |
| η | $\begin{cases} \text{quantity used as input to } ANISO time \\ \text{to fully specify the five elastic constants} \end{cases}$ | dimensionless | eq. (14) |

Note: Variables used only in the supplemental material are not included in the above table

Table 2: Description of symbols for seismic phases.

| Symbols | Wave type |
|---------|--|
| P | <i>P</i> -wave, upgoing or downgoing, in the mantle |
| p | upgoing P -wave from a seismic source |
| S | S-wave, upgoing or downgoing, in the mantle |
| s | upgoing S -wave from a seismic source |
| K | <i>P</i> -wave in the outer core |
| Ι | <i>P</i> -wave in the inner core |
| J | S-wave in the inner core |
| | Interactions |
| С | topside reflection off the core-mantle boundary |
| i | topside reflection off the inner-core/outer-core boundary |
| ^ | underside reflection, used primarily for crustal and mantle interfaces |
| v | topside reflection, used primarily for crustal and mantle interfaces |
| diff | appended to P or S to represent a diffracted wave along the core-mantle boundary |

³⁸¹ Author's note to copy editor:

- ³⁸² Because the upper half and lower half of the table show different things, please
- 383 do not delete any of the horizontal lines in this table, even though this differs
- ³⁸⁴ slightly from the usual conventions for typesetting tables.

Table 3: List of arguments for ANISOtime CLI.

| Parameter | Meaning |
|--------------------|---|
| -dec | Number of decimal places for output |
| -deg | Epicentral distance Δ [deg] |
| -h | Depth of source [km] (default:0) |
| mod | Structure: prem, iprem, ak135, |
| -mod | or path to "PolynomialStructure" or "Named Discontinuity" file (default:prem). |
| -help | Prints the usage. This option has the highest priority |
| -ph,phase | Seismic phases (default:P,PCP,PKiKP,S,ScS,SKiKS) |
| -р | Ray parameter [s/deg] |
| version | Shows version information. This option has the 2nd highest priority |
| -SH | Computes traveltime for SH (default:SH) |
| -SV | Computes traveltime for SV (default:SH) |
| -dD | Parameter for catalog creation $(d\Delta)$ |
| -dR | Integral interval [km] (default:10.0) |
| -eps | Output path figure |
| -0 | Output file for a record section. If it already exists, an error will be raised |
| delta | Show only epicentral distances |
| rayp | Show only ray parameters |
| time | Show only traveltimes |
| taup | Use the same output format as <i>taup_time</i> |
| -rc,read-catalog | Path of a catalog for which traveltimes are computed |
| -rs,record-section | start,end(,interval) [deg] Computes a table of a record section for the range |

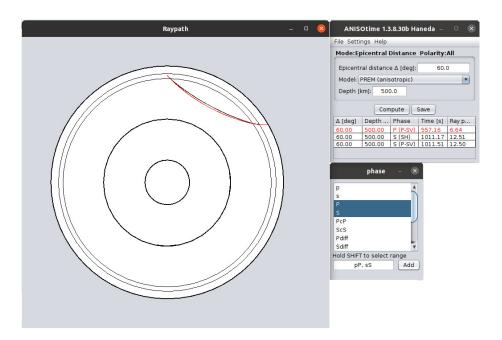


Figure 1: Sample figure made in epicentral distance mode. P- and S-waves with an epicentral distance of 60° are computed. The figure has been edited to remove the gray background in the raypath window.

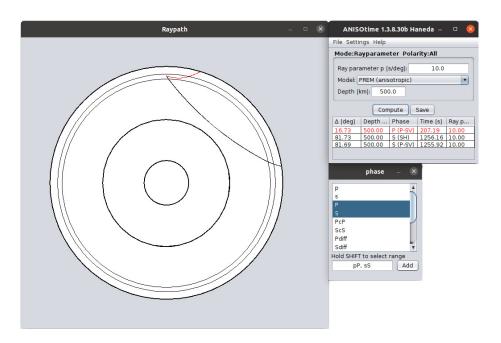


Figure 2: Sample figure made in ray parameter mode. P- and S-waves with a ray parameter of p = 10 s/degree are computed. The figure has been edited to remove the gray background in the raypath window.

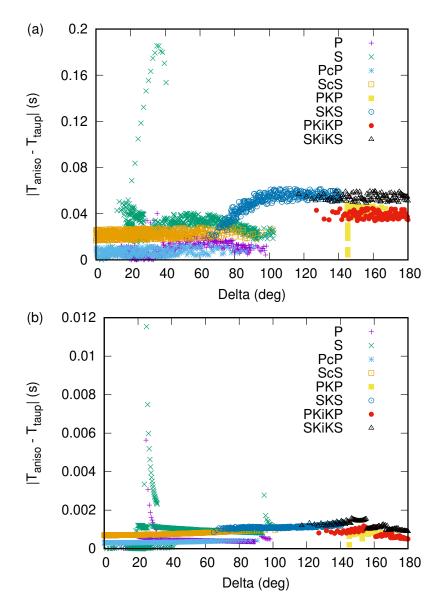


Figure 3: Absolute difference between traveltimes computed using ANISOtimeand TauP for all default phases for a source at the Earth's surface. The model used to compute traveltimes with ANISOtime is isotropic PREM "PolynomialStructure." The model used to compute traveltimes with TauP are a) the default (isotropic) PREM included in the TauP package; b) isotropic PREM, but with a finer depth sampling, resulting in increased accuracy as compared to traveltimes computed with the TauP default model (see text for details).

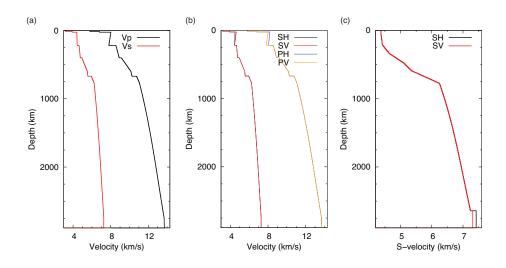


Figure 4: Models used to test *ANISOtime*. (a): isotropic PREM; (b) anisotropic PREM, with SH and PH velocities faster than SV and PV velocities, respectively, between 24 and 220 km depth; (c) MIASP91ANI, a modified version of IASP91 (Kennett and Engdahl, 1991) with smoothed discontinuities in the upper mantle (Borgeaud *et al.*, 2016) with SH velocity increased by 3%, and SV velocity increased by 1% in the lowermost 250 km of the mantle.

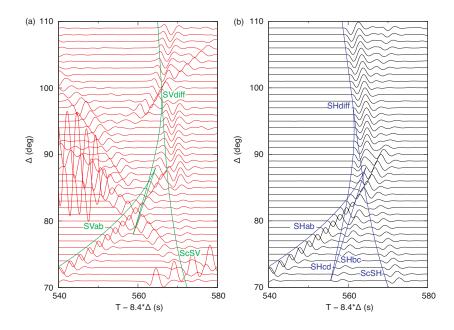


Figure 5: Waveforms for the radial (a) and transverse (b) components computed using the DSM for the anisotropic model in Fig. 4(c) for a source at 571.3 km depth. The traces are normalized using a time window around the S phase. The horizontal axis shows the reduced traveltime for a slowness of 8.4 s/degree. The traveltime curves for SV and SH waves (see labels), which were computed using ANISOtime, are in good general agreement with the synthetics.

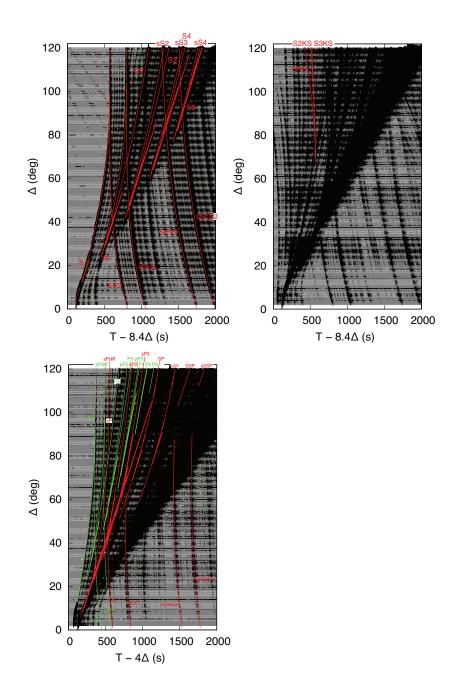


Figure 6: Waveforms for the transverse (a), radial (b), and vertical (c) components computed using the DSM for isotropic PREM (Fig. 4a) for a source at 571.3 km depth. The traces are self-normalized at each epicentral distance.

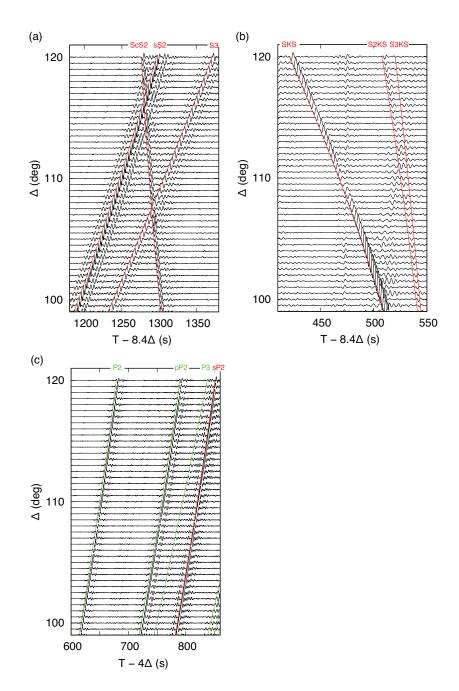


Figure 7: Closeups of the record sections in Fig. 6 showing epicentral distance range from 100 to 120 degrees, for the transverse (a), radial (b), and vertical (c) components.

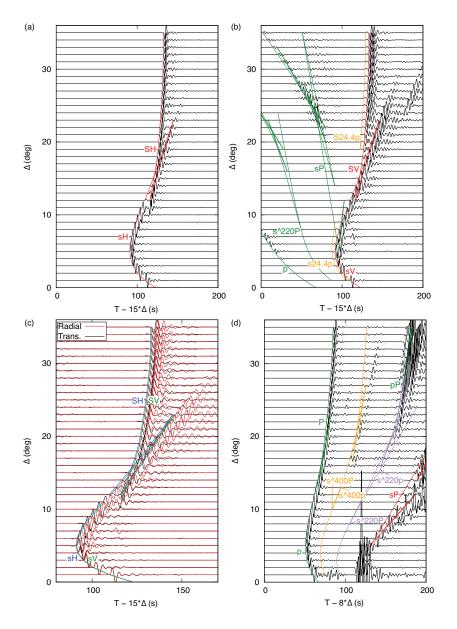


Figure 8: Waveforms computed using the DSM for anisotropic PREM (Fig. 4b) for a source at 571.3 km depth: a) transverse component; b) radial component; c) comparison of transverse and radial components; d) vertical component. The traces are normalized using time windows around the S phase (panels a, b, and c), and around the pseudo-P phase (panel d).

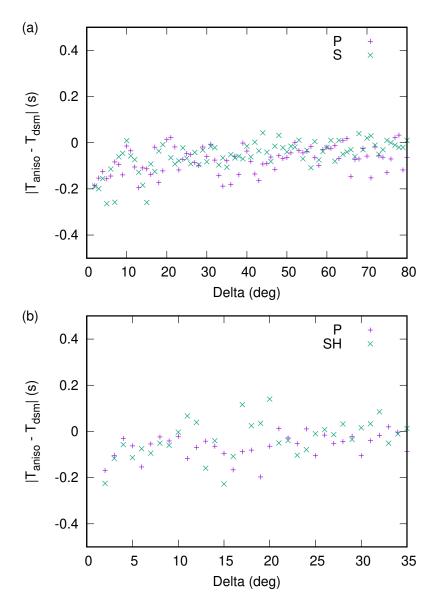


Figure 9: Differences between traveltimes computed using *ANISOtime* and onset times picked on synthetic waveforms for a) the isotropic PREM model (Fig. 6), and b) the anisotropic PREM model (Fig. 8) for the pseudo-P and SH phases.