

# ELASTIC-PLASTIC ANALYSIS OF BEAMS WITH UNIFORM CROSS-SECTION UNDER COMBINED LOADINGS (2)

—Bending and Its Combination with Torsion—

組合せ荷重を受けるはりの弾塑性解析 (2)—純曲げおよび曲げとねじり—

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## 1. Introduction

The preceding work<sup>1)</sup> in which we treated the combined loading of axial force and torsion of uniform beams is extended to include the pure bending and/or coexistence of the bending mode of deformation. When the bending predominates, it is appropriate to incorporate the quadratic shape function with respect to the coordinates  $x, y$  which are being taken in the cross-section of the beam or bar. Expository examples are concerned with the pure bending and its combination with torsion. The secondary stress inherent to the elastic-plastic deformation as well as the development of plastic enclave are studied in detail.

## 2. Finite element formulation

### (1) displacement function

We take the  $x$  and  $y$  axes in the plane of the cross-section and  $z$  axis to the axial direction of the beam similarly as in the preceding paper (see Fig. 1). We base our formulation on the following expressions of displacements  $u, v, w$  which have been shown to be obtained by the semi-inverse method<sup>2)</sup> for the uniform beam under the combined loadings of axial tensile (or compressive) force, bending and twisting moments:

$$\left. \begin{aligned} u &= f_1(x, y) + \frac{1}{2} \kappa_y z^2 - \theta y z \\ v &= f_2(x, y) - \frac{1}{2} \kappa_x z^2 + \theta z x \\ w &= f_3(x, y) - \kappa_y z x + \kappa_x y z + \varepsilon_0 z \end{aligned} \right\} (1)$$

where  $\varepsilon_0, \theta, \kappa_x$  and  $\kappa_y$  denote the longitudinal strain in the central layer, the angle of twist, and the curvatures in the  $yz$  and  $zx$  planes respectively. The symmetry of the transverse cross-section is assumed and the bending moments  $M_x$  and  $M_y$  are applied respectively in the planes  $yz$  and  $zx$  of symmetry. It is known that the functions  $f_1$  and  $f_2$  of Eq. (1) are given for the elastic deformation

$$\left. \begin{aligned} f_1(x, y) &= -\nu \kappa_x x y + \frac{\nu}{2} \kappa_y (x^2 - y^2) - \nu \varepsilon_0 x \\ f_2(x, y) &= -\frac{\nu}{2} \kappa_x (y^2 - x^2) + \nu \kappa_y x y - \nu \varepsilon_0 y \end{aligned} \right\} (2)$$

The function  $f_3$  represents the warping of cross-section due to the twisting moment  $T$ . Note from Eq. (1) that we are concerned with the deformation which is not dependent on the axial coordinate  $z$ . Therefore, we can assume only unit length along  $z$ -axis.

### (2) triangular element and shape function

Generally, the displacements  $u, v$  and  $w$  within each element shown in Fig. 1 are expressed as Eq. (1), and we assumed in the preceding paper

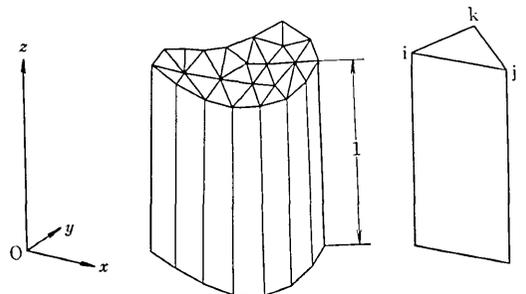


Fig. 1 Triangular element.

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that functions  $f_1, f_2$  and  $f_3$  are all linear with respect to the  $x, y$  coordinates. It was found, however, that this assumption lead to the results seriously in error when applied to the problems where the bending deformations prevail. Further, Eq. (2) indicates that the quadratic expressions of  $f_1$  and  $f_2$  give the exact solution for the elastic problem which accompanies the bending modes of deformation. Thus the shape function adopted in the present formulation is, for the triangular element, in terms of the area coordinates  $\zeta_i, \zeta_j$  and  $\zeta_k$  as follows:

$$\left. \begin{aligned} u &= \zeta_i(\zeta_i - \zeta_j - \zeta_k)u_i + \zeta_j(\zeta_j - \zeta_k - \zeta_i)u_j \\ &\quad + \zeta_k(\zeta_k - \zeta_i - \zeta_j)u_k + 4\zeta_j\zeta_k u_l + 4\zeta_k\zeta_i u_m \\ &\quad + 4\zeta_i\zeta_j u_n + \theta y(1-z) - \frac{1}{2}\kappa_y(1-z^2) \\ v &= \zeta_i(\zeta_i - \zeta_j - \zeta_k)v_i + \zeta_j(\zeta_j - \zeta_k - \zeta_i)v_j \\ &\quad + \zeta_k(\zeta_k - \zeta_i - \zeta_j)v_k + 4\zeta_j\zeta_k v_l + 4\zeta_k\zeta_i v_m \end{aligned} \right\}$$

$$\left. \begin{aligned} &+ 4\zeta_i\zeta_j v_n - \theta x(1-z) + \frac{1}{2}\kappa_x(1-z^2) \\ w &= \zeta_i w_i + \zeta_j w_j + \zeta_k w_k \\ &\quad + \kappa_y x(1-z) - \kappa_x y(1-z) - \varepsilon_0(1-z) \end{aligned} \right\} \quad (3)$$

The suffixes denote quantities associated with the corresponding corner nodes  $i, j, k$  and/or mid-side nodes  $l, m, n$ .

(3) stiffness equation

The strain components within each element are derived from the displacement functions of Eq. (3) as follows

$$\{\varepsilon\} = [B]\{d\} \quad (4)$$

where  $[\varepsilon] = [\varepsilon_x \ \varepsilon_y \ \varepsilon_z \ \gamma_{yz} \ \gamma_{zx} \ \gamma_{xy}]$   
 $[d] = [u_i \ v_i \ w_i \ u_j \ v_j \ w_j \ u_k \ v_k \ w_k \ u_l \ v_l \ w_l \ u_m \ v_m \ w_m \ u_n \ v_n \ w_n \ \varepsilon_0 \ \theta \ \kappa_x \ \kappa_y]$   
 $[B] = \zeta_i[B_i] + \zeta_j[B_j] + \zeta_k[B_k]$   
 and  $[B_i], [B_j], [B_k]$  are

$$[B_i] = \begin{bmatrix} 3b_i & 0 & 0 & -b_j & 0 & 0 & -b_k & 0 & 0 & 0 & 0 & 4b_k & 0 & 4b_j & 0 & 0 & 0 & 0 \\ 0 & 3c_i & 0 & 0 & -c_j & 0 & 0 & -c_k & 0 & 0 & 0 & 4c_k & 0 & 4c_j & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & y_i & -x_i \\ 0 & 0 & c_i & 0 & 0 & c_j & 0 & c_k & 0 & 0 & 0 & 0 & 0 & 0 & 0 & x_i & -1 & 0 \\ 0 & 0 & b_i & 0 & 0 & b_j & 0 & b_k & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -y_i & 0 & 1 \\ 3c_i & 3b_i & 0 & -c_j & -b_j & 0 & -c_k & -b_k & 0 & 0 & 0 & 4c_k & 4b_k & 4c_j & 4b_j & 0 & 0 & 0 & 0 \end{bmatrix}$$

$$[B_j] = \begin{bmatrix} -b_i & 0 & 0 & 3b_j & 0 & 0 & -b_k & 0 & 0 & 4b_k & 0 & 0 & 0 & 4b_i & 0 & 0 & 0 & 0 & 0 \\ 0 & -c_i & 0 & 0 & 3c_j & 0 & 0 & -c_k & 0 & 0 & 4c_k & 0 & 0 & 4c_i & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & y_j & -x_j \\ 0 & 0 & c_i & 0 & 0 & c_j & 0 & c_k & 0 & 0 & 0 & 0 & 0 & 0 & 0 & x_j & -1 & 0 & 0 \\ 0 & 0 & b_i & 0 & 0 & b_j & 0 & b_k & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -y_j & 0 & 1 & 0 \\ -c_i & -b_i & 0 & 3c_j & 3b_j & 0 & -c_k & -b_k & 0 & 4c_k & 4b_k & 0 & 0 & 4c_i & 4b_i & 0 & 0 & 0 & 0 \end{bmatrix}$$

$$[B_k] = \begin{bmatrix} -b_i & 0 & 0 & -b_j & 0 & 0 & 3b_k & 0 & 0 & 4b_j & 0 & 4b_i & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & -c_i & 0 & 0 & -c_j & 0 & 0 & 3c_k & 0 & 0 & 4c_j & 0 & 4c_i & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & y_k & -x_k \\ 0 & 0 & c_i & 0 & 0 & c_j & 0 & c_k & 0 & 0 & 0 & 0 & 0 & 0 & 0 & x_k & -1 & 0 & 0 \\ 0 & 0 & b_i & 0 & 0 & b_j & 0 & b_k & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -y_k & 0 & 1 & 0 \\ -c_i & -b_i & 0 & -c_j & -b_j & 0 & 3c_k & 3b_k & 0 & 4c_j & 4b_j & 4c_i & 4b_i & 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix}$$

with

$$\begin{bmatrix} a_i & a_j & a_k \\ b_i & b_j & b_k \\ c_i & c_j & c_k \end{bmatrix} = \frac{1}{2A} \begin{bmatrix} x_j y_k - x_k y_j & x_k y_i - x_i y_k & x_i y_j - x_j y_i \\ y_j - y_k & y_k - y_i & y_i - y_j \\ x_k - x_j & x_i - x_k & x_j - x_i \end{bmatrix}$$

where  $x_i, y_i, \dots, y_k$  are the coordinates of corner nodes and  $A$  is the area of the element triangle.

By using the stress-strain matrix  $[D]^{1),3)}$ , the associated stress field is expressed as

$$\{\sigma\} = [D]\{\varepsilon\} = [D][B]\{d\} \quad (5)$$

where  $[\sigma] = [\sigma_x \ \sigma_y \ \sigma_z \ \tau_{yz} \ \tau_{zx} \ \tau_{xy}]$

The strain energy  $U$  in each element can be expressed by Eqs. (4) and (5) as



$\Delta\sigma_y$  and  $\Delta\tau_{xy}$  were found to exist, but they are considerably small, compared with  $\Delta\sigma_x$ .

The stress increment of the element generally reverses its sign as the elastic-plastic boundary traverses the element. Therefore, it is difficult to elucidate the general tendency of the secondary stress (not the stress increment at each stage of deformation), particularly by the rather coarse element division used in this article. So, it is only remarked here that the absolute value of secondary stresses  $\sigma_x$ ,  $\sigma_y$ , and  $\tau_{xy}$  is the order of 3 kg/mm<sup>2</sup>, whereas the primary stress  $\sigma_x$  in axial direction is between 23.1 and 25.3 kg/mm<sup>2</sup> at the limit, i. e. fully plastic state of deformation.

The secondary stresses do not appear in the incompressible elastic-plastic material. Fig. 4 shows the secondary stress  $\sigma_x$  as well as its increment vanish, as the Poisson's ratio  $\nu$  approaches to 0.5.

(2) combined loading of bending and twist

For this combination, the numerical solution of Steele<sup>4)</sup> is available for the square bar of the plastic-rigid material. In order to compare with Steele's solution, computation was carried out for two cases where the value of strain ratio  $d\theta/d\kappa_x$

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$\sigma_x$  : stage 4  
 $\Delta\sigma_x$  : stage 3→4

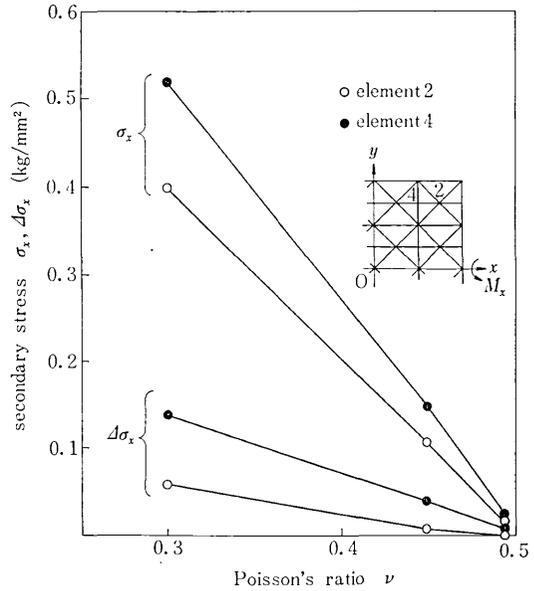


Fig. 4 Relation between secondary stress  $\sigma_x$  and Poisson's ratio  $\nu$ .

is kept constant ( $d\theta/d\kappa_x = \sqrt{3}$  and  $\sqrt{3}/2$  respectively) during the whole process of loading.

Fig. 5 depicts the development of plastic region for the case of strain ratio  $d\theta/d\kappa_x = \sqrt{3}$ . The plastic yielding starts at the extreme boundary

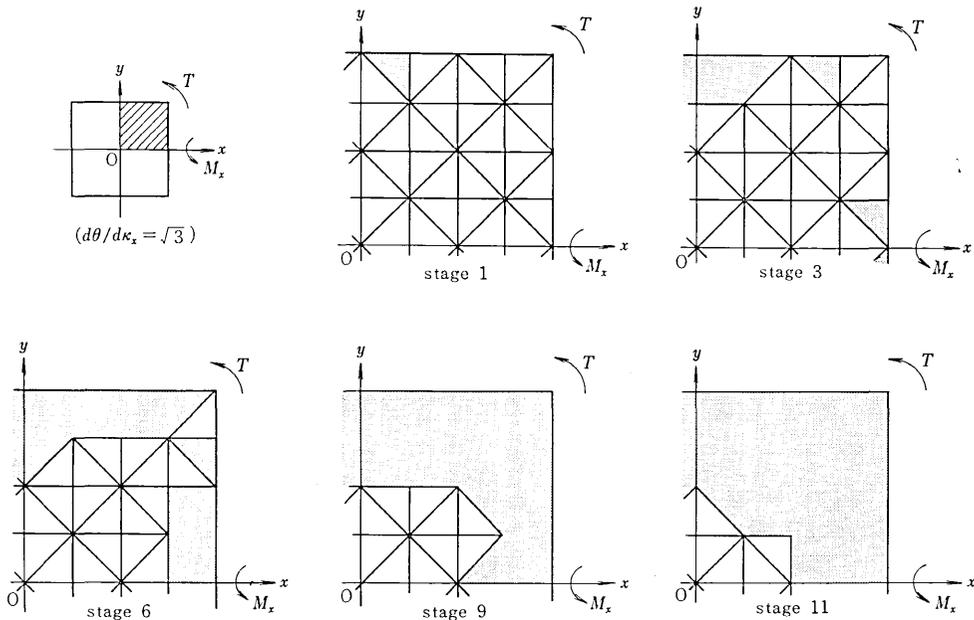


Fig. 5 Development of plastic region (combined loading of bending and twist)

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 element on  $y$  axis where both stresses due to bending and twist predominate. Subsequent yielding occurs in the element at the extreme of  $x$ , and then the plastic region expands to cover the whole cross-section as the loads increase.

Fig. 6 shows the axial stress distribution  $\sigma_x$  in the ratio to the yield stress  $\sigma_Y$  of the material at the fully plastic state of deformation. The stress at the centroid of the element is given. The relation between bending and twisting moments  $M, T$  at the fully plastic state is illustrated in

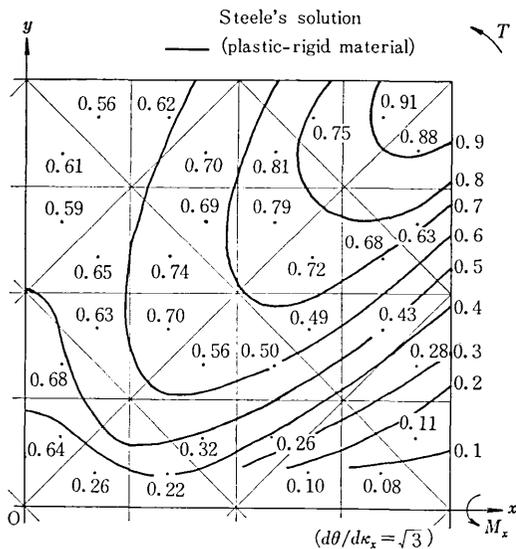


Fig. 6 Distribution of  $\sigma_x/\sigma_Y$ .

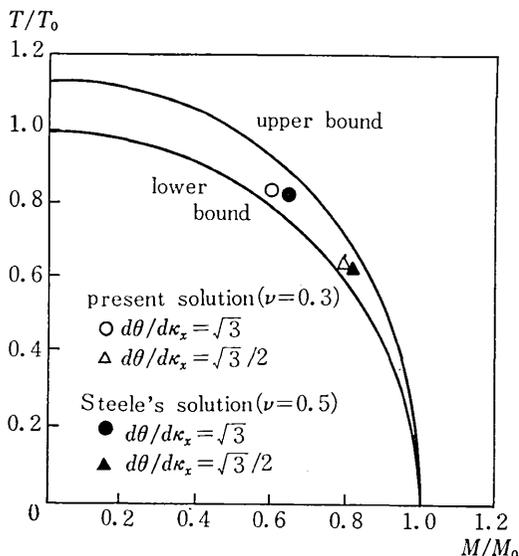


Fig. 7 Interaction curve (combined loading of bending and twist).

Fig. 7 where a comparison with the interaction curves of plastic-rigid analysis is also made.

It should be emphasized that the present formulation facilitates the analysis at every stages of deformation, not merely at the limiting fully plastic stage. Moreover, it is of interest to note that the  $M-T$  relation at the limiting stage obtained for the compressible elastic-plastic material lies between the upper and lower bounds of interaction curve for the plastic-rigid material.

As for the features of the secondary stress, it can only be said that its maximum absolute value is  $0.33 \text{ kg/mm}^2$  at the fully plastic state of the material whose yield stress is  $25 \text{ kg/mm}^2$ .

#### 4. Conclusions

The finite element solution procedure is formulated for the elastic-plastic deformation of uniform beams subjected to combined loadings which consist of axial tensile (or compressive) force, bending and twisting moments. In the case where the bending deformation predominates, the use of the quadratic displacement function is imperative. Present formulation facilitates the pursuit of the development of the plastic region. Further, it has become possible to make evaluation of the secondary stress which could not be assessed so far. An extension of the present formulation is attempted towards the problems which accompany the shear load and/or non-uniformity of cross-section, i. e. the cases where the deformation and stress are dependent on the axial coordinate  $z$ .

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#### References

- 1) Y. Yamada and K. Takatsuka, "Elasto-Plastic Analysis of Beams with Uniform Cross-Section under Combined Loadings", Seisan Kenkyu, 23-12 (1971), pp. 531-535.
- 2) Y. Yamada, "Sosei-Rikigaku (Theory of Plasticity)", Nikkanogyo Shinbunsha (1965)
- 3) Y. Yamada, T. Kawai, N. Yoshimura and T. Sakurai, "Analysis of the Elastic-Plastic Problems by the Matrix Displacement Method", AFFDL-TR-68-150 (1968), pp. 1271-1299.
- 4) M. C. Steele, "The Plastic Bending and Twisting of Square Section Members", J. Mech. Phys. Solids, 3-2 (1954), pp. 156-166.