

Polynomial relations for q -characters
via the ODE/IM correspondence

(ODE/IM対応を用いた q 指標の多項
式関係)

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POLYNOMIAL RELATIONS FOR Q -CHARACTERS VIA THE ODE/IM CORRESPONDENCE

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ABSTRACT. Let $U_q(\mathfrak{b})$ be the Borel subalgebra of a quantum affine algebra of type $X_n^{(1)}$ ($X = A, B, C, D$). Guided by the ODE/IM correspondence in quantum integrable models, we propose conjectural polynomial relations among the q -characters of certain representations of $U_q(\mathfrak{b})$.

Keywords: Borel subalgebra, q -character, Baxter's Q -operator, ODE/IM correspondence

1. INTRODUCTION

Let $U_q(\mathfrak{g})$ be a quantum affine algebra of type $X_n^{(1)}$ ($X = A, B, C, D$), and let $U_q(\mathfrak{b})$ be its Borel subalgebra. In this paper we shall consider the problem of finding polynomial relations satisfied by the q -characters of the fundamental modules in the sense of [1] and related modules. This problem is intimately related with that of functional equations for Baxter's Q -operators in quantum integrable models. In order to motivate the present study let us review this connection.

In quantum integrable systems, one is interested in the spectra of a commutative family of transfer matrices. The latter are constructed from the universal R matrix of a quantum affine algebra, by taking the trace of the first component over some finite dimensional representation called 'auxiliary space'. When the auxiliary spaces are the Kirillov-Reshetikhin (KR) modules, the corresponding transfer matrices satisfy an important family of polynomial identities known as the T -system [2, 3]. (For a recent survey on this topic, see [4].) Subsequently the T -system has been formulated and proved [5, 6] as identities of q -characters. It has been shown further [7] that the T -system is actually the defining relations of the Grothendieck ring of finite dimensional modules of quantum affine algebras. Notice that, from the construction by trace, the q -characters and the transfer matrices are both ring homomorphisms defined on the Grothendieck ring. Since the q -characters are injective [8], identities for q -characters imply the same identities for transfer matrices.

Baxter's Q -operators were first introduced in the study of the 8-vertex model [9]. Since then they have been recognized as a key tool in classical and quantum integrable systems, and there is now a vast literature on this subject. In the seminal paper [10], Bazhanov, Lukyanov and Zamolodchikov revealed that the Q -operators can also be obtained from the universal R matrix, provided the auxiliary space is chosen to be a (generically infinite-dimensional) representation of the Borel subalgebra. The work [10] for $U_q(\widehat{\mathfrak{sl}}_2)$ has been extended by several authors [11, 12, 13, 14, 15] to higher rank and supersymmetric cases.

In view of the results mentioned above, it is natural to ask whether one can find polynomial relations, analogous to the T -system, for the q -characters of $U_q(\mathfrak{b})$ (and hence for the Q -operators as well). The goal of this paper is to propose candidates of such identities.

Our idea is to use the so-called ODE/IM correspondence which relates the eigenvalues of Q -operators and certain ordinary linear differential equations. The reader is referred e.g. to the review [16] on this topic. In [17], the correspondence is discussed for general non-twisted affine Lie algebras using scalar (pseudo-)differential operators. In this paper we reformulate their results in terms of first order systems. The general setting (to be explained in Section 3) is as follows.

Let ${}^L\mathfrak{g}$ denote the Langlands dual Lie algebra of \mathfrak{g} . Consider the following ${}^L\mathfrak{g}$ -valued first order linear differential operator

$$(1.1) \quad \mathcal{L} = \frac{d}{dx} - \frac{\ell}{x} + \sum_{i=1}^n e_i + (x^{Mh^\vee} - E)e_0.$$

Here e_0, \dots, e_n are the Chevalley generators of ${}^L\mathfrak{g}$, ℓ is a generic element of the Cartan subalgebra ${}^L\mathfrak{h} \subset {}^L\mathfrak{g}$, h^\vee is the dual Coxeter number of \mathfrak{g} , and $M > 0$, $E \in \mathbb{C}$ are parameters. For each fundamental representation $V^{(a)}$ of ${}^L\mathfrak{g}$, there is a basis $\{\chi_J^{(a)}(x, E)\}$ of $V^{(a)}$ -valued solutions to the equation $\mathcal{L}\chi = 0$ characterized by the behavior

$$\chi_J^{(a)}(x, E) = \mathbf{u}_J^{(a)} x^{\lambda_J^{(a)}} (1 + O(x)) \quad (x \rightarrow 0),$$

where $\lambda_J^{(a)}$, $\mathbf{u}_J^{(a)}$ are the eigenvalues of ℓ and the corresponding eigenvectors. Hence this basis is labeled by weight vectors of $V^{(a)}$. There is also a canonical solution $\psi^{(a)}(x, E)$ which has the fastest decay at $x \rightarrow +\infty$. Due to the special choice (1.1) of \mathcal{L} , the canonical solutions are shown to satisfy a system of relations similar to the Plücker relations, called the ψ -system in [17] (see Subsection 3.2 below). Introduce the connection coefficients $Q_J^{(a)}(E) \in \mathbb{C}$ by

$$\psi^{(a)}(x, E) = \sum_J Q_J^{(a)}(E) \chi_J^{(a)}(x, E).$$

Then the ψ -system implies a set of polynomial relations among the connection coefficients $Q_J^{(a)}(E)$.

We expect the following to be true:

- (1) For each connection coefficient $Q_J^{(a)}(E)$, there is associated a formal power series $\mathcal{Q}_{J,z}^{(a)}$. Up to a simple multiplicative factor, the latter is the q -character of an irreducible highest ℓ -weight module of $U_q(\mathfrak{b})$. In particular, for the highest or lowest weights of $V^{(a)}$ the corresponding modules are the fundamental modules [1] of $U_q(\mathfrak{b})$.
- (2) With the identification

$$Q_J^{(a)}(E) \longleftrightarrow \mathcal{Q}_{J,z}^{(a)}, \quad E \longleftrightarrow z,$$

the polynomial relations for $Q_J^{(a)}(E)$ implied by the ψ -system hold true also for $\mathcal{Q}_{J,z}^{(a)}$.

We adopt (1), (2) as our working hypothesis for finding relations among the q -characters. A few remarks are in order here. We consider separately the cases related to the spin representation (i.e., $(\mathfrak{g}, a) = (C_n^{(1)}, n), (D_n^{(1)}, n-1), (D_n^{(1)}, n)$) and use the letter $\mathcal{R}_{\varepsilon,z}^{(a)}$ to denote the counterpart of $\mathcal{Q}_{J,z}^{(a)}$. We expect that $\mathcal{Q}_{J,z}^{(a)}$ or $\mathcal{R}_{\varepsilon,z}^{(a)}$ corresponds to the q -character of an irreducible highest ℓ -weight module of $U_q(\mathfrak{b})$ whose highest ℓ -weight is given by formulas (4.16), (4.19)–(4.20), (4.23)–(4.24). (In the case $C_n^{(1)}$, there are problems with this interpretation, as explained in Remark 5.9.) Also the relations for the $\mathcal{Q}_{J,z}^{(a)}$'s or $\mathcal{R}_{\varepsilon,z}^{(a)}$'s are not exactly the same as those for $Q_J^{(a)}(E)$'s, but it is necessary to fine-tune the coefficients by some power functions independent of z . The details will be given in Section 5 below. That it is natural to consider all $\mathcal{Q}_{J,z}^{(a)}$ corresponding to general weights of $V^{(a)}$ is a viewpoint suggested by the work [14] for type A algebras.

Now let us come to the content of the present work. As a first step, we give explicit candidates for $\mathcal{Q}_{i,z}^{(1)}$ associated with each weight of the vector representation $V^{(1)}$. This is done by taking suitable limits of the known q -characters of KR modules given by tableaux sums. It is known [1] that for the highest and lowest weights this procedure indeed gives the irreducible q -characters. As the next step, we define the formal series $\mathcal{Q}_{J,z}^{(a)}$ for a general node a of the Dynkin diagram. We define them by Casorati determinants whose entries are $\mathcal{Q}_{i,z}^{(1)}$ with suitable shifts of parameters. We then give candidates of polynomial relations among the $\mathcal{Q}_{J,z}^{(a)}$ expected from the ψ -system. In the cases related to spin representations, however, we do not have explicit candidates the $\mathcal{R}_{\varepsilon,z}^{(a)}$'s in general, so we only write down the candidates for the relations. The resulting polynomial relations

are given in Proposition 5.1, Theorem 5.2 (type $A_n^{(1)}$), Conjecture 5.3 (type $B_n^{(1)}$), Conjecture 5.5-5.8 (type $C_n^{(1)}$) and Conjecture 5.10-5.11 (type $D_n^{(1)}$), respectively.

For the type A algebras, one can check that $\mathcal{Q}_{i,z}^{(1)}$ (not necessarily the highest or lowest ones) are indeed irreducible q -characters of modules given in [13]. The polynomial relations corresponding to the ψ -system can be summarized as a single identity

$$(1.2) \quad \det \left(\mathcal{Q}_{\nu, q^{-2\mu+n+2}z}^{(1)} x_{\nu}^{-\mu+\nu} \right)_{\mu, \nu=1}^{n+1} = 1,$$

where $x_{\nu} = e^{\epsilon_{\nu}}$, $\{\epsilon_{\nu}\}_{\nu=1}^{n+1}$ being an orthonormal basis related to the simple roots by $\alpha_{\nu} = \epsilon_{\nu} - \epsilon_{\nu+1}$. This relation carries the name ‘Wronskian identity’ (as for its status, see the remark at the beginning of Appendix C). We shall give a direct proof that (1.2) is satisfied for q -characters in Appendix C.

For the other types of algebras, the situation is less satisfactory. At the moment we do not know the irreducibility of modules corresponding to the $\mathcal{Q}_{i,z}^{(1)}$ given by our procedure (except those corresponding to the highest or lowest vectors). For the spin representations of $C_n^{(1)}$ and $D_n^{(1)}$, explicit formulas for the $\mathcal{R}_{\varepsilon,z}^{(a)}$ ’s are missing. More seriously, we have not been able to prove the proposed identities for q -characters by computational methods. Instead, we support our working hypothesis by performing the following checks.

- (i) $\mathfrak{g} = B_2^{(1)}$, proof by hand
- (ii) $\mathfrak{g} = B_3^{(1)}$, computer check
- (iii) $\mathfrak{g} = D_4^{(1)}$, computer check
- (iv) $\mathfrak{g} = B_n^{(1)}, C_n^{(1)}, D_n^{(1)}$, proof in the limit to ordinary characters

The main results of the present paper consist in formulating the conjectured relations, and in performing the checks mentioned above.

The text is organized as follows. In Section 2, we prepare basic definitions concerning the Borel subalgebra $U_q(\mathfrak{b})$ of a quantum affine algebra $U_q(\mathfrak{g})$, and collect necessary facts about their representations. In Section 3, we give an account of the ψ -system in the ODE/IM correspondence and indicate how to derive them using the formulation by first order systems. We note that in [17] the ψ -system for algebras other than A type is mentioned as conjectures. In Section 4, we introduce the series $\mathcal{Q}_{i,z}^{(1)}$ as limits of the q -characters of KR modules. We also define these series for the other nodes of the Dynkin diagram. Section 5 is devoted to the proposals for polynomial relations. By comparing with the relations for the connection coefficients, we write down the relations for each type of algebras $A_n^{(1)}, B_n^{(1)}, C_n^{(1)}$ and $D_n^{(1)}$. In Section 6, we give a summary of our work.

The text is followed by four appendices. Appendix A gives a list of realizations of the dual Lie algebras ${}^L\mathfrak{g}$. In Appendix B, we give slightly more details on the procedure to obtain ψ -system by taking $\mathfrak{g} = A_2^{(1)}$ as an example. In Appendix C, we give a proof of the Wronskian identity for type $A_n^{(1)}$. In Appendix D, we prove the identities for type $B_n^{(1)}$ in the limit to ordinary characters. In Appendix E, the same is done for type $C_n^{(1)}$ and $D_n^{(1)}$.

2. PRELIMINARIES

In this section we introduce our notation on quantum affine algebras and their Borel subalgebras, and collect necessary facts that will be used later. Throughout this paper, we assume that q is a nonzero complex number which is not a root of unity.

2.1. Quantum Borel algebras. Let \mathfrak{g} be an affine Lie algebra associated with a generalized Cartan matrix $C = (c_{ij})_{0 \leq i, j \leq n}$ of non-twisted type. Let $D = \text{diag}(d_0, \dots, d_n)$ be the unique diagonal matrix such that DC is symmetric and $d_0 = 1$. Set $I = \{1, 2, \dots, n\}$, and let $\hat{\mathfrak{g}}$ denote the simple Lie algebra with the Cartan matrix $(c_{ij})_{i, j \in I}$. Let $\{\alpha_i\}_{i \in I}$, $\{\alpha_i^{\vee}\}_{i \in I}$ and $\{\omega_i\}_{i \in I}$ be the simple roots, simple coroots and the fundamental weights of $\hat{\mathfrak{g}}$, respectively. We set $P = \oplus_{i \in I} \mathbb{Z}\omega_i$, $Q = \oplus_{i \in I} \mathbb{Z}\alpha_i$.

Set $q_i = q^{d_i}$. We shall use the standard notation

$$[k]_i = \frac{q_i^k - q_i^{-k}}{q_i - q_i^{-1}}, \quad [n]_i! = \prod_{k=1}^n [k]_i, \quad \begin{bmatrix} n \\ k \end{bmatrix}_i = \frac{[n]_i!}{[k]_i! [n-k]_i!}.$$

The quantum affine algebra $U_q(\mathfrak{g})$ is the \mathbb{C} -algebra defined by generators $E_i, F_i, K_i^{\pm 1}$ ($i = 0, \dots, n$) and the relations

$$\begin{aligned} K_i K_i^{-1} &= 1 = K_i^{-1} K_i, \quad K_i K_j = K_j K_i, \\ K_i E_j K_i^{-1} &= q_i^{c_{ij}} E_j, \quad K_i F_j K_i^{-1} = q_i^{-c_{ij}} F_j, \\ [E_i, F_j] &= \delta_{ij} \frac{K_i - K_i^{-1}}{q_i - q_i^{-1}}, \\ \sum_{r=0}^{1-c_{ij}} \begin{bmatrix} 1-c_{ij} \\ r \end{bmatrix}_i (-1)^r E_i^r E_j E_i^{1-c_{ij}-r} &= 0, \quad i \neq j, \\ \sum_{r=0}^{1-c_{ij}} \begin{bmatrix} 1-c_{ij} \\ r \end{bmatrix}_i (-1)^r F_i^r F_j F_i^{1-c_{ij}-r} &= 0, \quad i \neq j. \end{aligned}$$

We do not write the formulas defining the Hopf structure on $U_q(\mathfrak{g})$ since we are not going to use them.

As is well known [18, 19], $U_q(\mathfrak{g})$ is isomorphic to the \mathbb{C} -algebra with generators $x_{i,r}^{\pm}$ ($i \in I, r \in \mathbb{Z}$), $k_i^{\pm 1}$ ($i \in I$), $h_{i,r}$ ($i \in I, r \in \mathbb{Z} \setminus \{0\}$) and central elements $c^{\pm 1/2}$, with the following defining relations

$$\begin{aligned} k_i k_i^{-1} &= 1 = k_i^{-1} k_i, \quad c^{1/2} c^{-1/2} = 1 = c^{-1/2} c^{1/2}, \\ k_i k_j &= k_j k_i, \quad k_i h_{j,r} = h_{j,r} k_i, \\ k_i x_{j,r}^{\pm} k_i^{-1} &= q_i^{\pm c_{ij}} x_{j,r}^{\pm}, \\ [h_{j,r}, x_{j,s}^{\pm}] &= \pm \frac{1}{r} [rc_{ij}]_i c^{\mp |r|/2} x_{j,r+s}^{\pm}, \\ x_{i,r+1}^{\pm} x_{j,s}^{\pm} - q_i^{\pm c_{ij}} x_{j,s}^{\pm} x_{i,r+1}^{\pm} &= q_i^{\pm c_{ij}} x_{i,r}^{\pm} x_{j,s+1}^{\pm} - x_{j,s+1}^{\pm} x_{i,r}^{\pm}, \\ [x_{i,r}^+, x_{j,s}^-] &= \delta_{ij} \frac{c^{(r-s)/2} \phi_{i,r+s}^+ - c^{-(r-s)/2} \phi_{i,r+s}^-}{q_i - q_i^{-1}}, \\ \sum_{\pi \in \mathfrak{S}_{1-c_{ij}}} \sum_{k=0}^{1-c_{ij}} (-1)^k \begin{bmatrix} 1-c_{ij} \\ k \end{bmatrix}_i x_{i,r_{\pi(1)}}^{\pm} \cdots x_{i,r_{\pi(k)}}^{\pm} x_{j,s}^{\pm} x_{i,r_{\pi(k+1)}}^{\pm} \cdots x_{i,r_{\pi(1-c_{ij})}}^{\pm} &= 0, \quad i \neq j, \end{aligned}$$

for all integers r_j , where \mathfrak{S}_m is the symmetric group on m letters, and the $\phi_{i,r}^{\pm}$ are given by

$$\sum_{r=0}^{\infty} \phi_{i,\pm r}^{\pm} u^{\pm r} = k_i^{\pm 1} \exp \left(\pm (q_i - q_i^{-1}) \sum_{s=1}^{\infty} h_{i,\pm s} u^{\pm s} \right).$$

By definition, the Borel subalgebra $U_q(\mathfrak{b})$ is the Hopf subalgebra of $U_q(\mathfrak{g})$ generated by $E_i, K_i^{\pm 1}$ ($0 \leq i \leq n$). It is known [20] that $U_q(\mathfrak{b})$ is isomorphic to the algebra with generators $E_i, K_i^{\pm 1}$ ($0 \leq i \leq n$) and the defining relations

$$\begin{aligned} K_i K_j &= K_j K_i, \quad K_i E_j K_i^{-1} = q_i^{c_{ij}} E_j, \\ \sum_{r=0}^{1-c_{ij}} \begin{bmatrix} 1-c_{ij} \\ r \end{bmatrix}_i (-1)^r E_i^r E_j E_i^{1-c_{ij}-r} &= 0, \quad i \neq j. \end{aligned}$$

The algebra $U_q(\mathfrak{b})$ contains $x_{i,m}^+, x_{i,r}^-, k_i^{\pm 1}$ and $\phi_{i,r}^+$, where $i \in I, m \geq 0$ and $r > 0$.

2.2. Category \mathcal{O} of $U_q(\mathfrak{b})$. In this subsection we recall basics about $U_q(\mathfrak{b})$ -modules in category \mathcal{O} . For more details see [1].

Denote by \mathfrak{t} the subalgebra of $U_q(\mathfrak{b})$ generated by $\{k_i^{\pm 1}\}_{i \in I}$. For a $U_q(\mathfrak{b})$ -module V and $\lambda \in P$, set

$$V_\lambda = \{v \in V \mid k_i v = q_i^{\lambda(\alpha_i^\vee)} v \quad (\forall i \in I)\}.$$

When $V_\lambda \neq 0$, it is called the weight space of weight λ . A $U_q(\mathfrak{b})$ -module V is said to be of type 1 if $V = \bigoplus_{\lambda \in P} V_\lambda$.

A series of complex numbers $\Psi = (\Psi_{i,r})_{i \in I, r \in \mathbb{Z}_{\geq 0}}$ is called an ℓ -weight if $\Psi_{i,0} \neq 0$ for all $i \in I$. We denote by \mathfrak{t}_ℓ^* the set of ℓ -weights. For a $U_q(\mathfrak{b})$ -module V and $\Psi \in \mathfrak{t}_\ell^*$, the subspace

$$V(\Psi) = \{v \in V \mid \exists p > 0, \forall i \in I, \forall m \geq 0, (\phi_{i,m}^+ - \Psi_{i,m})^p v = 0\}$$

is called the ℓ -weight space of ℓ -weight Ψ .

A $U_q(\mathfrak{b})$ -module V is said to be a highest ℓ -weight module of highest ℓ -weight $\Psi \in \mathfrak{t}_\ell^*$ if there exists a vector $v \in V$ such that $V = U_q(\mathfrak{b})v$ and

$$E_i v = 0 \quad (i \in I), \quad \phi_{i,r}^+ v = \Psi_{i,r} v \quad (i \in I, r \geq 0).$$

For each $\Psi \in \mathfrak{t}_\ell^*$ there exists a unique simple $U_q(\mathfrak{b})$ -module of highest ℓ -weight Ψ . We denote it by $L(\Psi)$.

A highest ℓ -weight module is of type 1 if its highest ℓ -weight Ψ satisfies

$$(2.1) \quad \Psi_{i,0} = q_i^{p_i} \quad \text{for some } p_i \in \mathbb{Z} \quad (i \in I).$$

For any non-zero complex numbers $c_i \in \mathbb{C}^\times$, the map $E_i \mapsto E_i, K_i \mapsto c_i K_i$ ($i = 0, \dots, n$) gives rise to an automorphism of $U_q(\mathfrak{b})$. After twisting by such an automorphism, any highest ℓ -weight module can be brought to one satisfying the condition (2.1). We denote by $\mathfrak{t}_{\ell,P}^*$ the set of ℓ -weights satisfying (2.1).

Set $D(\lambda) = \lambda - Q^+$, $Q^+ = \sum_{i \in I} \mathbb{Z}_{\geq 0} \alpha_i$. A $U_q(\mathfrak{b})$ -module V of type 1 is said to be an object in category \mathcal{O} if

- (i) for all $\lambda \in P$ we have $\dim V_\lambda < \infty$,
- (ii) there exist a finite number of elements $\lambda_1, \dots, \lambda_s \in P$ such that the weights of V are contained in $\bigcup_{j=1, \dots, s} D(\lambda_j)$.

In what follows we shall identify $\Psi \in \mathfrak{t}_\ell^*$ with their generating series,

$$\Psi = (\Psi_1(u), \dots, \Psi_n(u)), \quad \Psi_i(u) = \sum_{r \geq 0} \Psi_{i,r} u^r.$$

Simple objects in category \mathcal{O} are classified by the following theorem.

Theorem 2.1 ([1]). *Suppose that $\Psi \in \mathfrak{t}_{\ell,P}^*$. Then the simple module $L(\Psi)$ is an object in category \mathcal{O} if and only if $\Psi_i(u)$ is a rational function of u for any $i \in I$.*

In particular, for $i \in I$ and $z \in \mathbb{C}^\times$, the simple modules $L_{i,z}^\pm = L(\Psi)$ defined by the highest ℓ -weight

$$\Psi_j(u) = \begin{cases} (1 - zu)^{\pm 1} & j = i, \\ 1 & j \neq i, \end{cases}$$

are objects in category \mathcal{O} . These modules are called the fundamental representations [1].

It is known [21, 22] that finite dimensional simple $U_q(\mathfrak{g})$ -modules remain simple when restricted to $U_q(\mathfrak{b})$. According to the classification of the former [23, 24], the simple module $L(\Psi)$ is finite dimensional if its highest ℓ -weight has the form

$$\Psi_i(u) = q_i^{\deg P_i} \frac{P_i(q_i^{-1}u)}{P_i(q_i u)} \quad (\forall i \in I),$$

where $P_i(u)$ is a polynomial such that $P_i(0) = 1$. In the case where

$$P_j(u) = \begin{cases} (1 - q_i^{-m+1}zu)(1 - q_i^{-m+3}zu) \cdots (1 - q_i^{m-1}zu) & (j = i), \\ 1 & (j \neq i), \end{cases}$$

with some $i \in I$, $m \in \mathbb{Z}_{>0}$ and $z \in \mathbb{C}^\times$, the module $L(\Psi)$ is called a Kirillov-Reshetikhin(KR) module. We denote it by $W_{m,z}^{(i)}$.

2.3. Characters and q -characters. Let $\mathbb{Z}^{t_{\ell,P}^*}$ denote the set of maps $t_{\ell,P}^* \rightarrow \mathbb{Z}$. For $\Psi \in t_{\ell,P}^*$, define $[\Psi] \in \mathbb{Z}^{t_{\ell,P}^*}$ by $[\Psi](\Psi') = \delta_{\Psi,\Psi'}$. For a $U_q(\mathfrak{b})$ -module V of type 1 in category \mathcal{O} , its q -character $\chi_q(V)$ is defined as an element of $\mathbb{Z}^{t_{\ell,P}^*}$,

$$\chi_q(V) = \sum_{\Psi \in t_{\ell,P}^*} \dim V(\Psi) \cdot [\Psi].$$

Similarly let \mathbb{Z}^P denote the set of maps $P \rightarrow \mathbb{Z}$, and define $e^\lambda \in \mathbb{Z}^P$ by $e^\lambda(\mu) = \delta_{\lambda,\mu}$. The ordinary character $\chi(V)$ is an element of \mathbb{Z}^P ,

$$\chi(V) = \sum_{\lambda \in P} \dim V_\lambda \cdot e^\lambda.$$

We have a natural map $\varpi : \mathbb{Z}^{t_{\ell,P}^*} \rightarrow \mathbb{Z}^P$ which sends $[\Psi]$ to e^λ such that $\Psi_{i,0} = q_i^{\lambda(\alpha_i^\vee)}$. Under ϖ the q -character specializes to the ordinary character,

$$\chi(V) = \varpi(\chi_q(V)).$$

For an object V in category \mathcal{O} with highest ℓ -weight Ψ , its normalized q -character $\tilde{\chi}_q(V)$ is defined by

$$\tilde{\chi}_q(V) = [\Psi^{-1}] \cdot \chi_q(V).$$

In [1], the authors constructed the fundamental representations $L_{i,z}^\pm$ as certain limits of KR modules, and showed that the q -characters of the latter tend to those of $L_{i,z}^\pm$. Precisely, we have the following theorem.

Theorem 2.2 ([1]). (i) For all $1 \leq i \leq n$ we have

$$\tilde{\chi}_q(L_{i,z}^-) = \lim_{m \rightarrow \infty} \tilde{\chi}_q(W_{m,q_i^{-m}z}^{(i)}),$$

where $\lim_{m \rightarrow \infty} q^{-m}$ is understood as 0.

(ii) Let \mathcal{N}_i denote the multiplicity of α_i in the maximal root of \mathfrak{g} . If $\mathcal{N}_i = 1$, then we have

$$\tilde{\chi}_q(L_{i,z}^+) = \chi(L_{i,z}^+) = \chi(L_{i,z}^-).$$

3. ψ -SYSTEM

In this section, we reformulate the ψ -systems given in [17].

3.1. ${}^L\mathfrak{g}$ -connection. From now on, let \mathfrak{g} be an affine Lie algebra of type $X_n^{(1)}$ ($X = A, B, C, D$), and let ${}^L\mathfrak{g}$ denote its Langlands dual algebra. Let h^\vee be the dual Coxeter number of \mathfrak{g} (see Table 1).

\mathfrak{g}	$A_n^{(1)}$	$B_n^{(1)}$	$C_n^{(1)}$	$D_n^{(1)}$
${}^L\mathfrak{g}$	$A_n^{(1)}$	$A_{2n-1}^{(2)}$	$D_{n+1}^{(2)}$	$D_n^{(1)}$
h^\vee	$n+1$	$2n-1$	$n+1$	$2n-2$
$\dim V^{(1)}$	$n+1$	$2n$	$2n+2$	$2n$

TABLE 1.

Denote by e_j, f_j, h_j ($0 \leq j \leq n$) the Chevalley generators of ${}^L\mathfrak{g}$. We set $e = \sum_{j=1}^n e_j$. Fix also an element $\ell \in {}^L\mathfrak{h}$ from the Cartan subalgebra of ${}^L\mathfrak{g}$, and let $\zeta \in \mathbb{C}^\times$. We consider the following ${}^L\mathfrak{g}$ -valued connection (cf. e.g. [25]):

$$(3.1) \quad \mathcal{L} = \frac{d}{dx} - \frac{\ell}{x} + e + p(x, E)\zeta e_0,$$

$$(3.2) \quad p(x, E) = x^{Mh^\vee} - E \quad (M > 0, E \in \mathbb{C}).$$

Take an element $h_\rho \in {}^L\mathfrak{h}$ such that $[h_\rho, e_j] = e_j$ ($1 \leq j \leq n$) and $[h_\rho, e_0] = -(h^\vee - 1)e_0$. The choice (3.2) ensures the following symmetry property for $\mathcal{L} = \mathcal{L}(x, E; \zeta)$:

$$(3.3) \quad \omega^{k(h_\rho - 1)} \mathcal{L}(x, E; e^{2\pi i k} \zeta) \omega^{-kh_\rho} = \mathcal{L}(\omega^k x, \Omega^k E; \zeta),$$

where $k \in \mathbb{C}$ and

$$(3.4) \quad \omega = e^{\frac{2\pi i}{h^\vee(M+1)}}, \quad \Omega = \omega^{h^\vee M}.$$

On any finite dimensional ${}^L\mathfrak{g}$ -module V , (3.1) defines a first order system of differential equations $\mathcal{L}(x, E; 1)\phi(x, E) = 0$. Quite generally, for a ${}^L\mathfrak{g}$ -module V , we denote by V_k the ${}^L\mathfrak{g}$ -module obtained by twisting V by the automorphism $e_j \mapsto \exp(2\pi i k \delta_{j,0})e_j$, $f_j \mapsto \exp(-2\pi i k \delta_{j,0})f_j$. The operator $\mathcal{L}(x, E; e^{2\pi i k})$ represents the action of $\mathcal{L}(x, E; 1)$ on V_k . For a V -valued solution $\phi(x, E)$ we set

$$(3.5) \quad \phi_k(x, E) = \omega^{-kh_\rho} \phi(\omega^k x, \Omega^k E).$$

Then the symmetry (3.3) implies $\mathcal{L}(x, E; e^{2\pi i k})\phi_k(x, E) = 0$.

With each node a of the Dynkin diagram of ${}^L\mathfrak{g}$ is associated a fundamental module $V^{(a)}$ of ${}^L\mathfrak{g}$. We summarize our convention about them and some facts which will be used later. We omit the proofs since they can be found without difficulty.

Remark 3.1. In [17], scalar (pseudo-)differential equations are considered. Using the realization of ${}^L\mathfrak{g}$ given in Appendix A and rewriting the equation $\mathcal{L}\phi = 0$ for the highest component of ϕ , one obtains the formulas [17, (3.18)–(3.21)] (for simplicity we have taken $K = 1$ there).

The module $V^{(1)}$ is called the vector representation of ${}^L\mathfrak{g}$. Its explicit realization is given in Appendix A. If k is an integer, it is obvious that $V_k = V$ for any V . The vector representation for $\mathfrak{g} = C_n^{(1)}$ has the additional property

$$(3.6) \quad V_{\frac{1}{2}}^{(1)} \simeq V^{(1)} \quad (\mathfrak{g} = C_n^{(1)}).$$

For general a , we distinguish the following two cases,

$$\begin{aligned} (\text{NS}) : & \quad (\mathfrak{g}, a) \neq (C_n^{(1)}, n), (D_n^{(1)}, n-1), (D_n^{(1)}, n), \\ (\text{S}) : & \quad (\mathfrak{g}, a) = (C_n^{(1)}, n), (D_n^{(1)}, n-1), (D_n^{(1)}, n). \end{aligned}$$

We shall refer to them as the non-spin case and the spin case, respectively. Here and after we set

$$t = \begin{cases} 2 & (\mathfrak{g} = C_n^{(1)}), \\ 1 & (\text{otherwise}). \end{cases}$$

In the non-spin case, we have

$$(3.7) \quad V^{(a)} = \bigwedge^a V_{\frac{a-1}{2t}}^{(1)} \quad \text{for (NS)}.$$

In the case $\mathfrak{g} = B_n^{(1)}$, we have in addition

$$(3.8) \quad V_{\frac{1}{2}}^{(a)} \simeq V^{(2n-a)} \quad (\mathfrak{g} = B_n^{(1)}, a = 1, \dots, n),$$

where $V^{(a)}$ for $a > n$ stands for the right-hand side of (3.7).

In the case $(\mathfrak{g}, a) = (C_n^{(1)}, n)$, $V^{(n)}$ is the spin representation of the subalgebra $L\hat{\mathfrak{g}} = \mathfrak{o}(2n+1) \subseteq D_{n+1}^{(2)}$. Likewise, in the cases $(\mathfrak{g}, a) = (D_n^{(1)}, n-1)$ or $(D_n^{(1)}, n)$, $V^{(n-1)}$ and $V^{(n)}$ are the spin representations of the subalgebra $L\hat{\mathfrak{g}} = \mathfrak{o}(2n) \subseteq D_n^{(1)}$.

Let us consider the solutions at the irregular singularity $x = \infty$. It is convenient to use a gauge transformed form of \mathcal{L} ,

$$(3.9) \quad x^{Mh_\rho} \mathcal{L} x^{-Mh_\rho} = \frac{d}{dx} + \Lambda x^M - \frac{\ell + Mh_\rho}{x} - Ee_0 x^{-M(h_\rho-1)},$$

where $\Lambda = e + e_0$. Let $\mu^{(a)}$ be the eigenvalue of Λ on $V^{(a)}$ which has the largest real part. This eigenvalue is multiplicity free, and is given explicitly as follows:

$$\mu^{(a)} = \begin{cases} \frac{\sin \frac{\pi a}{th^\vee}}{\sin \frac{\pi}{th^\vee}} & \text{for (NS),} \\ \frac{1}{2 \sin \frac{\pi}{th^\vee}} & \text{for (S).} \end{cases}$$

Let $\mathbf{u}^{(a)}$ be an eigenvector of Λ corresponding to $\mu^{(a)}$. From the representation (3.9) it follows that there is a unique $V^{(a)}$ -valued solution $\psi^{(a)}(x, E)$ which satisfies the following in a sector containing the positive real axis $x > 0$:

$$\psi^{(a)}(x, E) = e^{-\frac{\mu^{(a)}}{M+1} x^{M+1}} x^{-Mh_\rho} \left(\mathbf{u}^{(a)} + o(1) \right) \quad (x \rightarrow \infty).$$

We call $\psi^{(a)}(x, E)$ the canonical solution. In view of the relation (3.7) and the formula for $\mu^{(a)}$ given above, we have

$$(3.10) \quad \psi^{(a)} = \psi_{-\frac{a-1}{2t}}^{(1)} \wedge \psi_{-\frac{a-3}{2t}}^{(1)} \wedge \cdots \wedge \psi_{\frac{a-1}{2t}}^{(1)} \quad \text{for (NS).}$$

3.2. ψ -system. In this subsection we state the ψ -system for the canonical solutions $\psi^{(a)}(x, E)$ introduced above. Let us consider them case by case.

$\mathfrak{g} = A_n^{(1)}$ ($L\mathfrak{g} = A_n^{(1)}$). For $a = 1, \dots, n$, we have the embedding of $L\mathfrak{g}$ -modules

$$(3.11) \quad \iota : \bigwedge^2 V_{\frac{1}{2}}^{(a)} \hookrightarrow V^{(a-1)} \otimes V^{(a+1)}$$

where $V^{(0)} = V^{(n+1)} = \mathbb{C}$. On the space $V^{(a-1)} \otimes V^{(a+1)}$, the functions $\phi = \iota(\psi_{-\frac{1}{2}}^{(a)} \wedge \psi_{\frac{1}{2}}^{(a)})$ and $\phi = \psi^{(a-1)} \otimes \psi^{(a+1)}$ both satisfy the equation $\mathcal{L}\phi = 0$ and have the behavior

$$\phi(x, E) = O \left(\exp \left(-\frac{2\mu^{(a)} \cos \frac{\pi}{h^\vee} x^{M+1}}{M+1} \right) \right) \quad (x \rightarrow \infty)$$

in a sector containing $x > 0$. Since such a solution is unique upto a constant multiple, we conclude that (after adjusting the constant multiple)

$$(3.12) \quad \iota(\psi_{-\frac{1}{2}}^{(a)} \wedge \psi_{\frac{1}{2}}^{(a)}) = \psi^{(a-1)} \otimes \psi^{(a+1)} \quad (a = 1, \dots, n).$$

In particular, denoting by $\{\mathbf{u}_j\}_{j=1}^{n+1}$ the standard basis of $V^{(1)}$ we have

$$(3.13) \quad \psi^{(n+1)} = \mathbf{u}_1 \wedge \cdots \wedge \mathbf{u}_{n+1}.$$

We call the relations (3.12), (3.13) the ψ -system for $A_n^{(1)}$.

The ψ -system for the other types can be deduced by the same argument, using the relevant embeddings of representations. We obtain relations of the following form.

$$(3.14) \quad \begin{aligned} &\underline{\mathfrak{g} = B_n^{(1)} (L\mathfrak{g} = A_{2n-1}^{(2)})}. \\ &\iota(\psi_{-\frac{1}{2}}^{(a)} \wedge \psi_{\frac{1}{2}}^{(a)}) = \psi^{(a-1)} \otimes \psi^{(a+1)} \quad (a = 1, \dots, n-1), \end{aligned}$$

$$(3.15) \quad \iota(\psi_{-\frac{1}{4}}^{(n)} \wedge \psi_{\frac{1}{4}}^{(n)}) = \psi_{-\frac{1}{4}}^{(n-1)} \otimes \psi_{\frac{1}{4}}^{(n-1)}.$$

Here ι stands for the embedding of ${}^L\mathfrak{g}$ (B.1) or

$$\bigwedge^2 V_{\frac{1}{4}}^{(n)} \hookrightarrow V_{-\frac{1}{4}}^{(n-1)} \otimes V_{\frac{1}{4}}^{(n-1)},$$

which follows from (B.1) and (3.8).

$$\underline{\mathfrak{g}} = C_n^{(1)} \quad ({}^L\mathfrak{g} = D_{n+1}^{(2)}).$$

$$(3.16) \quad \iota(\psi_{-\frac{1}{4}}^{(a)} \wedge \psi_{\frac{1}{4}}^{(a)}) = \psi_{-\frac{1}{4}}^{(a-1)} \otimes \psi_{\frac{1}{4}}^{(a+1)} \quad (a = 1, \dots, n-2),$$

$$(3.17) \quad \iota(\phi^{(n-1)}) = \psi_{-\frac{1}{2}}^{(n)} \wedge \psi_{\frac{1}{2}}^{(n)},$$

$$(3.18) \quad \iota(\phi^{(n)}) = \psi_{-\frac{1}{4}}^{(n)} \otimes \psi_{\frac{1}{4}}^{(n)}.$$

Here we have set

$$\begin{aligned} \phi^{(n-1)} &= \psi_{-\frac{n-2}{4}}^{(1)} \wedge \psi_{-\frac{n-4}{4}}^{(1)} \wedge \dots \wedge \psi_{\frac{n-2}{4}}^{(1)}, \\ \phi^{(n)} &= \psi_{-\frac{n-1}{4}}^{(1)} \wedge \psi_{-\frac{n-3}{4}}^{(1)} \wedge \dots \wedge \psi_{\frac{n-1}{4}}^{(1)}, \end{aligned}$$

and ι stands for an analog of (B.1) or the embeddings

$$\begin{aligned} \bigwedge^{n-1} V_{\frac{n-2}{4}}^{(1)} &\simeq V^{(n-1)} \hookrightarrow \bigwedge^2 V_{\frac{1}{2}}^{(n)}, \\ \bigwedge^n V_{\frac{n-1}{4}}^{(1)} &\hookrightarrow V_{-\frac{1}{4}}^{(n)} \otimes V_{\frac{1}{4}}^{(n)}, \end{aligned}$$

where (3.6) is taken into account.

$$\underline{\mathfrak{g}} = D_n^{(1)} \quad ({}^L\mathfrak{g} = D_n^{(1)}).$$

$$(3.19) \quad \iota(\psi_{-\frac{1}{2}}^{(a)} \wedge \psi_{\frac{1}{2}}^{(a)}) = \psi_{-\frac{1}{2}}^{(a-1)} \otimes \psi_{\frac{1}{2}}^{(a+1)} \quad (a = 1, \dots, n-3),$$

$$(3.20) \quad \iota(\psi^{(n-2)}) = \psi_{-\frac{1}{2}}^{(n-1)} \wedge \psi_{\frac{1}{2}}^{(n-1)}, \quad \iota(\psi^{(n-2)}) = \psi_{-\frac{1}{2}}^{(n)} \wedge \psi_{\frac{1}{2}}^{(n)},$$

$$(3.21) \quad \iota(\phi^{(n-1)}) = \psi^{(n-1)} \otimes \psi^{(n)}.$$

We have set

$$\phi^{(n-1)} = \psi_{-\frac{n-2}{2}}^{(1)} \wedge \psi_{-\frac{n-4}{2}}^{(1)} \wedge \dots \wedge \psi_{\frac{n-2}{2}}^{(1)}$$

and ι stands for the embedding (B.1) or

$$\begin{aligned} \bigwedge^{n-2} V_{\frac{n-3}{2}}^{(1)} &\simeq V^{(n-2)} \hookrightarrow \bigwedge^2 V_{\frac{1}{2}}^{(n-1)}, \quad \bigwedge^{n-2} V_{\frac{n-3}{2}}^{(1)} \simeq V^{(n-2)} \hookrightarrow \bigwedge^2 V_{\frac{1}{2}}^{(n)}, \\ \bigwedge^{n-1} V_{\frac{n-2}{2}}^{(1)} &\hookrightarrow V^{(n-1)} \otimes V^{(n)}. \end{aligned}$$

3.3. Connection coefficients. Now we introduce the connection coefficients $Q_j^{(a)}(E)$.

First let us consider the vector representation $V^{(1)}$. We choose ℓ generic. Set $\ell = \text{diag}(\ell_1, \dots, \ell_N)$ ($N = \dim V^{(1)}$, see Appendix A), and let \mathbf{u}_j ($1 \leq j \leq N$) be the corresponding eigenvector of ℓ . Then there is a unique $V^{(1)}$ -valued solution $\chi_j^{(1)}(x, E)$ characterized by the expansion at the origin,

$$\chi_j^{(1)}(x, E) = x^{\ell_j} \mathbf{u}_j (1 + O(x)) \quad (x \rightarrow 0).$$

From the symmetry (3.3) of $\mathcal{L}(x, E)$ we find that

$$(3.22) \quad \chi_{j,k}^{(1)}(x, E) = \omega^{k\lambda_j} \chi_j^{(1)}(x, E) \quad (k \in (1/t)\mathbb{Z}),$$

where we have set

$$(3.23) \quad \ell - h_\rho = \text{diag}(\lambda_1, \dots, \lambda_N).$$

Define $Q_j^{(1)}(E) \in \mathbb{C}$ by

$$(3.24) \quad \psi^{(1)}(x, E) = \sum_{j=1}^N Q_j^{(1)}(E) \chi_j^{(1)}(x, E).$$

From (3.22) and (3.24), we have for $k \in (1/t)\mathbb{Z}$

$$\psi_k^{(1)}(x, E) = \sum_{j=1}^N Q_{j,k}^{(1)}(E) \omega^{k\lambda_j} \chi_j^{(1)}(x, E),$$

where $Q_{j,k}^{(1)}(E) = Q_j^{(1)}(\Omega^k E)$. For a sequence $J = (j_1, \dots, j_a)$, introduce the notation

$$Q_J^{(a)}(E) = \det \left(Q_{j_m, \frac{2l-a-1}{2t}}^{(1)}(E) \cdot \omega^{\lambda_{j_m}(l-j)} \right)_{1 \leq l, m \leq a},$$

$$\chi_J^{(a)}(x, E) = \chi_{j_1, -\frac{a-1}{2t}}^{(1)}(x, E) \wedge \dots \wedge \chi_{j_a, -\frac{a-1}{2t}}^{(1)}(x, E).$$

It follows from (3.10) that in the non-spin case

$$(3.25) \quad \psi^{(a)}(x, E) = \sum_J Q_J^{(a)}(E) \chi_J^{(a)}(x, E),$$

where the sum is taken over all $J = (j_1, \dots, j_a)$, $1 \leq j_1 < \dots < j_a \leq N$.

One can similarly define the connection coefficients in the spin case as well.

4. SERIES $\Omega_{J,z}^{(a)}$, $\mathcal{R}_{\varepsilon,z}^{(a)}$

It has been shown in [1] that the fundamental modules $L_{i,z}^\pm$ of the Borel subalgebra arise as certain limits of the KR modules. In this section, associated with each weight space of the vector representation of $L\mathfrak{g}$ we obtain a formal power series $\Omega_{i,z}^{(1)}$, which is obtained as certain limit of the q -character of KR module $W_{m,z}^{(1)}$. By Theorem 2.2, up to simple overall multipliers, those formal series corresponding to the highest or lowest weights are the irreducible q -characters $\chi_q(L_{1,z}^\pm)$. We expect that in general the $\Omega_{i,z}^{(1)}$ are also proportional to irreducible q -characters of $U_q(\mathfrak{b})$. (See however Remark 5.9.)

We shall use the following elements of $\mathbb{Z}^{t^*,P}$:

$$(4.1) \quad \mathcal{Y}_{i,z} = [(1, \dots, (1 - zu)^{-1}, \dots, 1)], \quad e^{\omega_i} = [(1, \dots, q_i^{i\text{-th}}, \dots, 1)],$$

$$(4.2) \quad Y_{i,z} = [(1, \dots, q_i^{\frac{1 - q_i^{-1}zu}{1 - q_i zu}}, \dots, 1)] = e^{\omega_i} \mathcal{Y}_{i,q_i z} \mathcal{Y}_{i,q_i^{-1}z}^{-1},$$

$$(4.3) \quad A_{i,z} = Y_{i,q_i^{-1}z} Y_{i,q_i z} \prod_{\substack{1 \leq j \leq n, \\ c_{ji} = -1}} Y_{j,z}^{-1} \prod_{\substack{1 \leq j \leq n, \\ c_{ji} = -2}} Y_{j,q_j^{-1}z}^{-1} Y_{j,q_j z}^{-1} \prod_{\substack{1 \leq j \leq n, \\ c_{ji} = -3}} Y_{j,q_j^{-2}z}^{-1} Y_{j,z}^{-1} Y_{j,q_j^2 z}^{-1}.$$

Highest ℓ -weights are monomials in $\mathcal{Y}_{i,z}^{\pm 1}$ and $e^{\pm \omega_i}$. Abusing the notation, for a monomial $M = [\Psi]$ we shall also write $L(M)$ for $L(\Psi)$.

4.1. q -character of KR module $W_{m,z}^{(1)}$. For the KR module $W_{m,z}^{(1)}$, explicit formulas for the q -characters are known as sums over tableaux [3]. We recall these formulas below.

We first introduce some notations. In this subsection, for a sequence $\mathbf{k} = (k_1, \dots, k_u)$, we set $|\mathbf{k}| = \sum_{a=1}^u k_a$ and $K_a = \sum_{b=a}^u k_b$ for any $1 \leq a \leq u$. Then the q -character of the KR module $W_{m,z}^{(1)}$ of quantum affine algebra $U_q(\mathfrak{g})$ is given as:

$$\underline{\mathfrak{g}} = A_n^{(1)}:$$

$$\chi_q(W_{m,z}^{(1)}) = \prod_{j=1}^m Y_{1,q^{m+1-2j}z} \sum_{\substack{k=(k_1,\dots,k_{n+1}), \\ |k|=m, k_a \geq 0(\forall a)}} \prod_{b=1}^n \prod_{j=1}^{K_{b+1}} A_{b,q^{m+1+b-2j}z}^{-1},$$

$$\underline{\mathfrak{g}} = B_n^{(1)}:$$

$$\begin{aligned} \chi_q(W_{m,z}^{(1)}) &= \prod_{j=1}^m Y_{1,q^{m+1-2j}z} \\ &\times \left\{ \sum_{\substack{k=(k_1,\dots,k_n,k_{\bar{n}},\dots,k_{\bar{1}}), \\ |k|=m, k_a \geq 0(\forall a)}} \prod_{b=1}^n \prod_{j=1}^{K_{b+1}} A_{b,q^{m+1+b-2j}z}^{-1} \prod_{j=1}^{K_{\bar{b}}} A_{b,q^{m+2n-b-2j}z}^{-1} \right. \\ &\quad \left. + \sum_{\substack{k=(k_1,\dots,k_n,k_{\bar{n}},\dots,k_{\bar{1}}), \\ |k|=m-1, k_a \geq 0(\forall a)}} \prod_{b=1}^n \prod_{j=1}^{K_{b+1}+1} A_{b,q^{m+1+b-2j}z}^{-1} \prod_{j=1}^{K_{\bar{b}}} A_{b,q^{m+2n-b-2j}z}^{-1} \right\}, \end{aligned}$$

where we set $k_{n+1} = k_{\bar{n}}$.

$$\underline{\mathfrak{g}} = C_n^{(1)}:$$

$$\begin{aligned} \chi_q(W_{m,z}^{(1)}) &= \prod_{j=1}^m Y_{1,q^{\frac{m+1-2j}{2}}z} \sum_{\substack{k=(k_1,\dots,k_n,2k_0,k_{\bar{n}},\dots,k_{\bar{1}}), \\ |k|=m, k_b \geq 0(\forall b)}} \prod_{b=1}^{n-1} \prod_{j=1}^{K_{b+1}} A_{b,q^{\frac{m+1+b-2j}{2}}z}^{-1} \\ &\times \prod_{b=1}^{n-1} \prod_{j=1}^{K_{\bar{b}}} A_{b,q^{\frac{m+2n+3-b-2j}{2}}z}^{-1} \prod_{j=1}^{K_{\bar{n}}} A_{n,q^{\frac{m+n+2-2j}{2}}z}^{-1} \prod_{j=1}^{k_0} A_{n,q^{\frac{m+n+2-2(K_{\bar{n}}+2j)}{2}}z}^{-1}, \end{aligned}$$

where in this case $|k| = \sum_{i=1}^n k_i + 2k_0 + \sum_{i=1}^n k_{\bar{i}}$, and for $a \preceq n$, $K_a = \sum_{i=a}^n k_i + 2k_0 + \sum_{i=1}^n k_{\bar{i}}$,

$$\underline{\mathfrak{g}} = D_n^{(1)}:$$

$$\chi_q(W_{m,z}^{(1)}) = \chi_q\left(\left(W_{m,z}^{(1)}\right)^+\right) + \chi_q\left(\left(W_{m,z}^{(1)}\right)^-\right),$$

where

$$\begin{aligned} \chi_q\left(\left(W_{m,z}^{(1)}\right)^+\right) &= \prod_{j=1}^m Y_{1,q^{m+1-2j}z} \sum_{\substack{k=(k_1,\dots,k_{n-1},k_n,k_{\frac{n-1}{2}},\dots,k_{\bar{1}}), \\ k_a \geq 0(\forall a), |k|=m}} \prod_{b=1}^{n-2} \prod_{j=1}^{K_{\bar{b}}} A_{b,q^{m+2n-1-b-2j}z}^{-1} \\ &\times \prod_{j=1}^{K_{\frac{n-1}{2}}} A_{n,q^{m+n-2j}z}^{-1} \prod_{b=1}^{n-1} \prod_{j=1}^{K_{b+1}} A_{b,q^{m+1+b-2j}z}^{-1}, \\ \chi_q\left(\left(W_{m,z}^{(1)}\right)^-\right) &= \prod_{j=1}^m Y_{1,q^{m+1-2j}z} \sum_{\substack{k=(k_1,\dots,k_{n-1},k_{\bar{n}},k_{\frac{n-1}{2}},\dots,k_{\bar{1}}), \\ k_a \geq 0(\forall a \neq \bar{n}), k_{\bar{n}} \geq 1, |k|=m}} \prod_{b=1}^{n-1} \prod_{j=1}^{K_{\bar{b}}} A_{b,q^{m+2n-1-b-2j}z}^{-1} \\ &\times \prod_{j=1}^{K_{\bar{n}}} A_{n,q^{m+n-2j}z}^{-1} \prod_{b=1}^{n-2} \prod_{j=1}^{K_{b+1}} A_{b,q^{m+1+b-2j}z}^{-1}. \end{aligned}$$

4.2. The limiting procedure. The series $\mathcal{Q}_{i,z}^{(1)}$ is defined as a limit of the q -character of KR module $W_{m,z}^{(1)}$. Let us illustrate on the case $\mathfrak{g} = A_2^{(1)}$ an example the procedure for taking the limit.

Example 4.1 ($\mathfrak{g} = A_2^{(1)}$). We consider first the case $\mathfrak{g} = A_2^{(1)}$. Following [3], we write

$$[1]_z = Y_{1,z}, \quad [2]_z = Y_{1,q^2z}^{-1} Y_{2,qz}, \quad [3]_z = Y_{2,q^3z}^{-1}.$$

Then the q -character $\chi_q(W_{m,z}^{(1)})$ of the KR module $W_{m,z}^{(1)}$ is presented as a sum over the tableaux

$$\sum_{\substack{k_1+k_2+k_3=m, \\ k_1, k_2, k_3 \geq 0}} \overbrace{\begin{array}{|c|c|c|c|c|c|} \hline 1 & \cdots & 1 & 2 & \cdots & 2 \\ \hline \end{array}}^{k_1} \overbrace{\begin{array}{|c|c|c|c|c|c|} \hline 2 & \cdots & 2 & 3 & \cdots & 3 \\ \hline \end{array}}^{k_2} \overbrace{\begin{array}{|c|c|c|c|c|c|} \hline 3 & \cdots & 3 & 3 & \cdots & 3 \\ \hline \end{array}}^{k_3}$$

where the k -th box from the right carries the parameter $q^{m+1-2k}z$. This can be rewritten further as

$$\begin{aligned} & \prod_{j=1}^m Y_{1,q^{m+1-2j}z} \sum_{\substack{k_1+k_2+k_3=m, \\ k_1, k_2, k_3 \geq 0}} \prod_{j=1}^{k_2+k_3} A_{1,q^{m+2-2j}z}^{-1} \prod_{j=1}^{k_3} A_{2,q^{m+3-2j}z}^{-1} \\ &= \sum_{\substack{k_1+k_2+k_3=m, \\ k_1, k_2, k_3 \geq 0}} e^{(k_1-k_2)\omega_1} e^{(k_2-k_3)\omega_2} \\ & \times \left[\left(\frac{(1-q^{-m}z)(1-q^{m-2k_3+2}z)}{(1-q^{m-2(k_2+k_3)}z)(1-q^{m-2(k_2+k_3)+2}z)} \right), \left(\frac{(1-q^{m-2(k_3+k_2)+1}z)(1-q^{m+3}z)}{(1-q^{m-2k_3+1}z)(1-q^{m-2k_3+3}z)} \right) \right]. \end{aligned}$$

Let us consider the limit $m \rightarrow \infty$. There are three possibilities to obtain meaningful answers,

$$\begin{aligned} k_1 &\rightarrow \infty, & k_2, k_3 &: \text{finite}, & q^{-k_1}z &\rightarrow 0, \\ k_2 &\rightarrow \infty, & k_1, k_3 &: \text{finite}, & q^{-k_2}z &\rightarrow 0, \\ k_3 &\rightarrow \infty, & k_1, k_2 &: \text{finite}, & q^{-k_3}z &\rightarrow 0. \end{aligned}$$

Writing $e^{\alpha_i} = x_i/x_{i+1}$ and defining for $i = 1, 2, 3$

$$\mathcal{Q}_{i,z}^{(1)} = \prod_{j=1}^{i-1} (1 - x_j/x_i) \times \lim_{\substack{k_i \rightarrow \infty \\ q^{-k_i}z \rightarrow 0}} x_i^{-m} (x_1 x_2 x_3)^{\frac{m}{3}} \chi_q(W_{m,q^{-m}z}^{(1)}),$$

we obtain the result

$$\begin{aligned} \mathcal{Q}_{1,z}^{(1)} &= y_{1,z} \sum_{k_2, k_3 \geq 0} \prod_{j=1}^{k_2+k_3} A_{1,q^{2-2j}z}^{-1} \prod_{j=1}^{k_3} A_{2,q^{3-2j}z}^{-1}, \\ \mathcal{Q}_{2,z}^{(1)} &= y_{1,q^2z}^{-1} y_{2,qz} \sum_{k_3 \geq 0} \prod_{j=1}^{k_3} A_{2,q^{3-2j}z}^{-1}, \\ \mathcal{Q}_{3,z}^{(1)} &= y_{2,q^3z}^{-1}. \end{aligned}$$

Up to simple multipliers, they are the irreducible q -characters

$$\begin{aligned} \chi_q(L_{1,z}^-) &= \mathcal{Q}_{1,z}^{(1)}, \quad \chi_q(L(y_{1,q^2z}^{-1} y_{2,qz})) = \frac{\mathcal{Q}_{2,z}^{(1)}}{1 - e^{-\alpha_1}}, \\ \chi_q(L_{2,q^3z}^+) &= \frac{\mathcal{Q}_{3,z}^{(1)}}{(1 - e^{-\alpha_1 - \alpha_2})(1 - e^{-\alpha_2})}. \end{aligned}$$

Using the explicit construction of modules [13], it can be checked that the second one is the q -character of the simple module whose highest ℓ -weight corresponds to the monomial $y_{1,q^2z}^{-1} y_{2,qz}$.

4.3. **Series $\Omega_{i,z}^{(1)}$.** In order to discuss the general case, let us prepare some notation. For $\mathfrak{g} = X_n^{(1)}$ ($X = A, B, C, D$), introduce a parametrization of $\{\alpha_i, \omega_i\}$ by orthonormal vectors $\{\epsilon_i\}$, and an index set \mathcal{J} with a partial ordering \prec as follows.

$$\begin{aligned}
A_n^{(1)} : \quad & \alpha_i = \epsilon_i - \epsilon_{i+1} \quad (1 \leq i \leq n), \\
& \omega_i = \epsilon_1 + \cdots + \epsilon_i - \frac{i}{n+1} \sum_{j=1}^{n+1} \epsilon_j \quad (1 \leq i \leq n), \\
& \mathcal{J} = \{1, 2, \dots, n+1\}, \quad 1 \prec 2 \prec \cdots \prec n+1, \\
B_n^{(1)} : \quad & \alpha_i = \epsilon_i - \epsilon_{i+1} \quad (1 \leq i \leq n-1), \quad \alpha_n = \epsilon_n, \\
& \omega_i = \epsilon_1 + \cdots + \epsilon_i \quad (1 \leq i \leq n-1), \\
& \omega_n = \frac{1}{2}(\epsilon_1 + \cdots + \epsilon_n), \\
& \mathcal{J} = \{1, \dots, n, \bar{n}, \dots, \bar{1}\}, \quad 1 \prec \cdots \prec n \prec \bar{n} \prec \cdots \prec \bar{1}, \\
C_n^{(1)} : \quad & \alpha_i = \frac{1}{\sqrt{2}}(\epsilon_i - \epsilon_{i+1}) \quad (1 \leq i \leq n-1), \quad \alpha_n = \sqrt{2}\epsilon_n, \\
& \omega_i = \frac{1}{\sqrt{2}}(\epsilon_1 + \cdots + \epsilon_i) \quad (1 \leq i \leq n), \\
& \mathcal{J} = \{1, \dots, n, 0, \bar{n}, \dots, \bar{1}\}, \quad 1 \prec \cdots \prec n \prec 0 \prec \bar{n} \prec \cdots \prec \bar{1}, \\
D_n^{(1)} : \quad & \alpha_i = \epsilon_i - \epsilon_{i+1} \quad (1 \leq i \leq n-1), \quad \alpha_n = \epsilon_{n-1} + \epsilon_n, \\
& \omega_i = \epsilon_1 + \cdots + \epsilon_i \quad (1 \leq i \leq n-2), \\
& \omega_{n-1} = \frac{1}{2}(\epsilon_1 + \cdots + \epsilon_{n-1} - \epsilon_n), \\
& \omega_n = \frac{1}{2}(\epsilon_1 + \cdots + \epsilon_{n-1} + \epsilon_n), \\
& \mathcal{J} = \{1, \dots, n, \bar{n}, \dots, \bar{1}\}, \quad 1 \prec \cdots \prec n-1 \prec \frac{n}{\bar{n}} \prec \overline{n-1} \prec \cdots \prec \bar{1}.
\end{aligned}$$

Define also x_i and $f_{j,k}$ for $j, k \in \mathcal{J}$ by

$$\begin{aligned}
A_n^{(1)} : \quad & x_i = e^{\epsilon_i} \quad (1 \leq i \leq n+1), \quad f_{j,k} = \frac{1}{1 - x_k/x_j}, \\
B_n^{(1)} : \quad & x_i = e^{\epsilon_i} = x_{\bar{i}}^{-1} \quad (1 \leq i \leq n), \quad f_{j,k} = \frac{1 + \delta_{k,\bar{j}}/x_j}{1 - x_k/x_j}, \\
C_n^{(1)} : \quad & x_i = e^{\frac{1}{\sqrt{2}}\epsilon_i} = x_{\bar{i}}^{-1} \quad (1 \leq i \leq n), \quad x_0 = 1, \\
& f_{j,k} = \frac{1}{(1 - x_k/x_j)(1 + \delta_{k,0}x_k/x_j)}, \\
D_n^{(1)} : \quad & x_i = e^{\epsilon_i} = x_{\bar{i}}^{-1} \quad (1 \leq i \leq n), \quad f_{j,k} = \frac{1 - \delta_{k,\bar{j}}x_k/x_j}{1 - x_k/x_j}.
\end{aligned}$$

Taking the same procedure as $\mathfrak{g} = A_2^{(1)}$, we give the definitions of the series $\mathcal{Q}_{i,z}^{(1)}$. Precisely, for each $i \in \mathcal{J}$, we define

$$\mathcal{Q}_{i,z}^{(1)} = \prod_{j \prec i} f_{i,j}^{-1} \begin{cases} \lim_{\substack{k_i \rightarrow \infty, \\ q^{-k_i} z \rightarrow 0}} \left(x_i \prod_{j=1}^{n+1} x_j^{-\frac{1}{n+1}} \right)^{-m} \chi_q \left(W_{m,q^{-m}z}^{(1)} \right), & A_n^{(1)}, \\ \lim_{\substack{k_i \rightarrow \infty, \\ q^{-k_i} z \rightarrow 0}} x_i^{-m} \chi_q \left(W_{m,q^{-m}z}^{(1)} \right), & B_n^{(1)}, \\ \begin{cases} \lim_{\substack{k_i \rightarrow \infty, \\ q^{-\frac{k_i}{2}} z \rightarrow 0}} x_i^{-m} \chi_q \left(W_{m,q^{-\frac{m}{2}}z}^{(1)} \right) & i \neq 0, \\ \lim_{\substack{k_0 \rightarrow \infty, \\ q^{-\frac{k_0}{2}} z \rightarrow 0}} \chi_q \left(W_{m,q^{-\frac{m}{2}}z}^{(1)} \right) & i = 0, \end{cases} & C_n^{(1)}, \\ \begin{cases} \lim_{\substack{k_i \rightarrow \infty, \\ q^{-k_i} z \rightarrow 0}} x_i^{-m} \chi_q \left(W_{m,q^{-m}z}^{(1)} \right) & i \neq n, \bar{n}, \\ \lim_{\substack{k_n \rightarrow \infty, \\ q^{-k_n} z \rightarrow 0}} x_n^{-m} \chi_q \left(\left(W_{m,q^{-m}z}^{(1)} \right)^+ \right) & i = n, \\ \lim_{\substack{k_{\bar{n}} \rightarrow \infty, \\ q^{-k_{\bar{n}}} z \rightarrow 0}} x_{\bar{n}}^{-m} \chi_q \left(\left(W_{m,q^{-m}z}^{(1)} \right)^- \right) & i = \bar{n}, \end{cases} & D_n^{(1)}. \end{cases}$$

In the following, in the sum of the form \sum_{k_i, \dots, k_j} , unless mentioned explicitly, k_i, \dots, k_j run over all non-negative integers, and we use the abbreviation $K_l = \sum_{\mu=l}^j k_\mu$ for $i \prec l \prec j$. In the case $\mathfrak{g} = C_n^{(1)}$ and $l \prec 0$, we set $K_l = \sum_{j=l}^n k_j + 2k_0 + \sum_{j=1}^n k_{\bar{j}}$.

We give below the explicit formula for $\mathcal{Q}_{i,z}^{(1)}$ for each $i \in \mathcal{J}$.

case $A_n^{(1)}$: For $1 \leq i \leq n+1$,

$$\mathcal{Q}_{i,z}^{(1)} = \Phi_{i,z} \sum_{k_{i+1}, \dots, k_{n+1}} \prod_{j=i}^n \prod_{l=1}^{K_{j+1}} A_{j,q^{j+1-2l}z}^{-1},$$

where

$$(4.4) \quad \Phi_{i,z} = \mathcal{Y}_{i-1,q^i z}^{-1} \mathcal{Y}_{i,q^{i-1}z}.$$

case $B_n^{(1)}$: For $1 \leq i \leq n$,

$$\begin{aligned} \mathcal{Q}_{i,z}^{(1)} &= \Phi_{i,z} \\ &\times \left\{ \sum_{k_{i+1}, \dots, k_n, k_{\bar{n}}, \dots, k_{\bar{1}}} \prod_{j=i}^n \prod_{l=1}^{K_{j+1}} A_{j,q^{j+1-2l}z}^{-1} \prod_{j=1}^n \prod_{l=1}^{K_{\bar{j}}} A_{j,q^{2n-j-2l}z}^{-1} \right. \\ &\quad \left. + \sum_{k_{i+1}, \dots, k_n, k_{\bar{n}}, \dots, k_{\bar{1}}} \prod_{j=i}^n \prod_{l=1}^{1+K_{j+1}} A_{j,q^{j+1-2l}z}^{-1} \prod_{j=1}^n \prod_{l=1}^{K_{\bar{j}}} A_{j,q^{2n-j-2l}z}^{-1} \right\}, \\ \mathcal{Q}_{\bar{i},z}^{(1)} &= \Phi_{\bar{i},z} \sum_{k_{\bar{i}-1}, \dots, k_{\bar{1}}} \prod_{j=1}^{i-1} \prod_{l=1}^{K_{\bar{j}}} A_{j,q^{2n-j-2l}z}^{-1}, \end{aligned}$$

where

$$(4.5) \quad \Phi_{i,z} = \mathcal{Y}_{i-1,q^i z}^{-1} \mathcal{Y}_{i,q^{i-1}z}, \quad \Phi_{\bar{i},z} = \mathcal{Y}_{i-1,q^{2n-i-1}z}^{-1} \mathcal{Y}_{i,q^{2n-i}z}.$$

case $C_n^{(1)}$:

$$\begin{aligned}
\Omega_{i,z}^{(1)} &= \Phi_{i,z} \sum_{k_{i+1}, \dots, k_n, k_0, k_{\bar{n}}, \dots, k_{\bar{1}}} \prod_{j=i}^{n-1} \prod_{l=1}^{K_{j+1}} A_{j,q}^{-1} \frac{1+j-2l}{2} z \\
&\quad \times \prod_{j=1}^{n-1} \prod_{l=1}^{K_{\bar{j}}} A_{j,q}^{-1} \frac{2n+3-j-2l}{2} z \prod_{l=1}^{K_{\bar{n}}} A_{n,q}^{-1} \frac{n+2-2l}{2} z \prod_{l=1}^{k_0} A_{n,q}^{-1} \frac{n+2-2(K_{\bar{n}}+2l)}{2} z \quad (1 \leq i \leq n), \\
\Omega_{0,z}^{(1)} &= \widehat{\Omega}_{0,z}^{(1)} + \widehat{\Omega}_{0,z}^{(1)}, \\
\widehat{\Omega}_{0,z}^{(1)} &= \Phi_{0,z} \sum_{\substack{k_{\bar{n}}, \dots, k_{\bar{1}}, \\ K_{\bar{n}}: \text{even}}} \prod_{j=1}^{n-1} \prod_{l=1}^{K_{\bar{j}}} A_{j,q}^{-1} \frac{2n+3-j-l}{2} z \prod_{l=1}^{\frac{K_{\bar{n}}}{2}} A_{n,q}^{-1} \frac{n+4}{2} - 2l z, \\
\widehat{\Omega}_{0,z}^{(1)} &= x_n^{-1} \Phi_{\bar{0},z} \sum_{\substack{k_{\bar{n}}, \dots, k_{\bar{1}}, \\ K_{\bar{n}}: \text{odd}}} \prod_{j=1}^{n-1} \prod_{l=1}^{K_{\bar{j}}} A_{j,q}^{-1} \frac{2n+3-j-l}{2} z \prod_{l=1}^{\frac{K_{\bar{n}}-1}{2}} A_{n,q}^{-1} \frac{n+2}{2} - 2l z, \\
\Omega_{i,z}^{(1)} &= \Phi_{\bar{i},z} \sum_{k_{\bar{i}-1}, \dots, k_{\bar{1}}} \prod_{j=1}^{i-1} \prod_{l=1}^{K_{\bar{j}}} A_{j,q}^{-1} \frac{2n+3-j-l}{2} z \quad (1 \leq i \leq n),
\end{aligned}$$

where

$$(4.6) \quad \Phi_{i,z} = y_{i-1,q^{\frac{i}{2}}z}^{-1} y_{i,q^{\frac{i-1}{2}}z}, \quad \Phi_{\bar{i},z} = y_{i-1,q^{\frac{2n+2-i}{2}}z}^{-1} y_{i,q^{\frac{2n+3-i}{2}}z} \quad (1 \leq i \leq n-1),$$

$$(4.7) \quad \Phi_{n,z} = y_{n-1,q^{\frac{n}{2}}z}^{-1} y_{n,q^{\frac{n-2}{2}}z} y_{n,q^{\frac{n}{2}}z}, \quad \Phi_{\bar{n},z} = y_{n-1,q^{\frac{n+2}{2}}z}^{-1} y_{n,q^{\frac{n+2}{2}}z} y_{n,q^{\frac{n+4}{2}}z}^{-1},$$

$$(4.8) \quad \Phi_{0,z} = y_{n,q^{\frac{n}{2}}z}^{-1} y_{n,q^{\frac{n+2}{2}}z}, \quad \Phi_{\bar{0},z} = y_{n-1,q^{\frac{n}{2}}z}^{-1} y_{n-1,q^{\frac{n+2}{2}}z} y_{n,q^{\frac{n-2}{2}}z}^{-1} y_{n,q^{\frac{n+4}{2}}z}^{-1}.$$

case $D_n^{(1)}$:

$$\begin{aligned}
\Omega_{i,z}^{(1)} &= \Phi_{i,z} \left\{ \sum_{k_{i+1}, \dots, k_n, \dots, k_{\bar{1}}} \prod_{j=1}^{n-2} \prod_{l=1}^{K_{\bar{j}}} A_{j,q^{2n-1-j-2l}z}^{-1} \prod_{l=1}^{K_{\bar{n}-1}} A_{n,q^{n-2l}z}^{-1} \prod_{j=i}^{n-1} \prod_{l=1}^{K_{j+1}} A_{j,q^{j+1-2l}z}^{-1} \right. \\
&\quad \left. + \sum_{\substack{k_{i+1}, \dots, k_{\bar{n}}, \dots, k_{\bar{1}}, \\ k_{\bar{n}} \geq 1}} \prod_{j=1}^{n-1} \prod_{l=1}^{K_{\bar{j}}} A_{j,q^{2n-1-j-2l}z}^{-1} \prod_{l=1}^{K_{\bar{n}}} A_{n,q^{n-2l}z}^{-1} \prod_{j=i}^{n-2} \prod_{l=1}^{K_{j+1}} A_{j,q^{j+1-2l}z}^{-1} \right\} \\
&\quad (1 \leq i \leq n-1), \\
\Omega_{n,z}^{(1)} &= \Phi_{n,z} \sum_{k_{\bar{n}-1}, \dots, k_{\bar{1}}} \prod_{j=1}^{n-2} \prod_{l=1}^{K_{\bar{j}}} A_{j,q^{2n-1-j-2l}z}^{-1} \prod_{l=1}^{K_{\bar{n}-1}} A_{n,q^{n-2l}z}^{-1}, \\
\Omega_{i,z}^{(1)} &= \Phi_{\bar{i},z} \sum_{k_{\bar{i}-1}, \dots, k_{\bar{1}}} \prod_{j=1}^{i-1} \prod_{l=1}^{K_{\bar{j}}} A_{j,q^{2n-1-j-2l}z}^{-1} \quad (1 \leq i \leq n),
\end{aligned}$$

where

$$(4.9) \quad \Phi_{i,z} = y_{i-1,q^i z}^{-1} y_{i,q^{i-1}z}, \quad \Phi_{\bar{i},z} = y_{i-1,q^{2n-2-i}z}^{-1} y_{i,q^{2n-1-i}z} \quad (1 \leq i \leq n-2),$$

$$(4.10) \quad \Phi_{n-1,z} = y_{n-2,q^{n-1}z}^{-1} y_{n-1,q^{n-2}z} y_{n,q^{n-2}z}, \quad \Phi_{\bar{n}-1,z} = y_{n-2,q^{n-1}z}^{-1} y_{n-1,q^n z} y_{n,q^n z}^{-1},$$

$$(4.11) \quad \Phi_{n,z} = y_{n-1,q^n z}^{-1} y_{n,q^{n-2}z}, \quad \Phi_{\bar{n},z} = y_{n-1,q^{n-2}z}^{-1} y_{n,q^n z}.$$

We give two more examples of explicit formulas.

Example 4.2 ($\mathfrak{g} = B_2^{(1)}$).

$$\begin{aligned}
Q_{1,z}^{(1)} &= y_{1,z} \left\{ \sum_{k_2, k_{\bar{2}}, k_{\bar{1}} \geq 0} \prod_{l=1}^{K_{\bar{1}}} A_{1,q^{3-2l}z}^{-1} \prod_{l=1}^{K_{\bar{2}}} A_{2,q^{2-2l}z}^{-1} \prod_{l=1}^{K_{\bar{2}}} A_{2,q^{3-2l}z}^{-1} \prod_{l=1}^{K_2} A_{1,q^{2-2l}z}^{-1} \right. \\
&\quad \left. + \sum_{k_2, k_{\bar{2}}, k_{\bar{1}} \geq 0} \prod_{l=1}^{K_{\bar{1}}} A_{1,q^{3-2l}z}^{-1} \prod_{l=1}^{K_{\bar{2}}} A_{2,q^{2-2l}z}^{-1} \prod_{l=1}^{K_{\bar{2}}+1} A_{2,q^{3-2l}z}^{-1} \prod_{l=1}^{K_2+1} A_{1,q^{2-2l}z}^{-1} \right\}, \\
Q_{2,z}^{(1)} &= y_{1,q^2z}^{-1} y_{2,qz} \left\{ \sum_{k_{\bar{2}}, k_{\bar{1}} \geq 0} \prod_{l=1}^{K_{\bar{1}}} A_{1,q^{3-2l}z}^{-1} \prod_{l=1}^{K_{\bar{2}}} A_{2,q^{2-2l}z}^{-1} \prod_{l=1}^{K_{\bar{2}}} A_{2,q^{3-2l}z}^{-1} \right. \\
&\quad \left. + \sum_{k_{\bar{2}}, k_{\bar{1}} \geq 0} \prod_{l=1}^{K_{\bar{1}}} A_{1,q^{3-2l}z}^{-1} \prod_{l=1}^{K_{\bar{2}}} A_{2,q^{2-2l}z}^{-1} \prod_{l=1}^{K_{\bar{2}}+1} A_{2,q^{3-2l}z}^{-1} \right\}, \\
Q_{2,z}^{(1)} &= y_{1,qz} y_{2,q^2z}^{-1} \sum_{k \geq 0} \prod_{l=1}^k A_{1,q^{3-2l}z}^{-1}, \\
Q_{1,z}^{(1)} &= y_{1,q^3z}^{-1}.
\end{aligned}$$

In this case $N_1 = 1$, hence by Theorem 2.2, we have

$$\chi_q(L_{1,z}^-) = Q_{1,z}^{(1)}, \quad \chi_q(L_{1,q^3z}^+) = \frac{Q_{1,z}^{(1)}}{(1 - e^{-\alpha_1})(1 - e^{-\alpha_1 - \alpha_2})(1 - e^{-\alpha_1 - 2\alpha_2})}.$$

Example 4.3 ($\mathfrak{g} = C_2^{(1)}$).

$$\begin{aligned}
Q_{1,z}^{(1)} &= y_{1,z} \sum_{k_2, k_0, k_{\bar{2}}, k_{\bar{1}} \geq 0} \prod_{l=1}^{K_{\bar{1}}} A_{1,q^{3-l}z}^{-1} \prod_{l=1}^{K_{\bar{2}}} A_{2,q^{2-l}z}^{-1} \prod_{l=1}^{k_0} A_{2,q^{2-K_{\bar{2}}-2l}z}^{-1} \prod_{l=1}^{K_2} A_{1,q^{1-l}z}^{-1}, \\
Q_{2,z}^{(1)} &= y_{1,qz}^{-1} y_{2,z} y_{2,qz} \sum_{k_0, k_{\bar{2}}, k_{\bar{1}} \geq 0} \prod_{l=1}^{K_{\bar{1}}} A_{1,q^{3-l}z}^{-1} \prod_{l=1}^{K_{\bar{2}}} A_{2,q^{2-l}z}^{-1} \prod_{l=1}^{k_0} A_{2,q^{2-K_{\bar{2}}-2l}z}^{-1}, \\
Q_{0,z}^{(1)} &= \widehat{Q}_{0,z}^{(1)} + \widehat{Q}_{0,z}^{(1)}, \\
\widehat{Q}_{0,z}^{(1)} &= y_{2,qz} y_{2,q^2z}^{-1} \sum_{\substack{k_{\bar{2}}, k_{\bar{1}} \geq 0, \\ K_{\bar{2}}: \text{even}}} \prod_{l=1}^{K_{\bar{1}}} A_{1,q^{3-l}z}^{-1} \prod_{l=1}^{\frac{K_{\bar{2}}}{2}} A_{2,q^{3-2l}z}^{-1}, \\
\widehat{Q}_{0,z}^{(1)} &= x_2^{-1} y_{1,qz}^{-1} y_{1,q^2z} y_{2,z} y_{2,q^3z}^{-1} \sum_{\substack{k_{\bar{2}}, k_{\bar{1}} \geq 0, \\ K_{\bar{2}}: \text{odd}}} \prod_{l=1}^{K_{\bar{1}}} A_{1,q^{3-l}z}^{-1} \prod_{l=1}^{\frac{K_{\bar{2}}-1}{2}} A_{2,q^{2-2l}z}^{-1}, \\
Q_{2,z}^{(1)} &= y_{1,q^2z} y_{2,q^2z}^{-1} y_{2,q^3z}^{-1} \sum_{k_{\bar{1}} \geq 0} \prod_{l=1}^{k_{\bar{1}}} A_{1,q^{3-l}z}^{-1}, \\
Q_{1,z}^{(1)} &= y_{1,q^3z}^{-1}.
\end{aligned}$$

Note that the weight 0 has multiplicity 2. Correspondingly $Q_{0,z}^{(1)}$ is a sum of two terms $\widehat{Q}_{0,z}^{(1)}$ and $\widehat{Q}_{0,z}^{(1)}$. These terms cannot be separated in the process of taking the limit. Since they have highest ℓ -weights whose ratio is not a monomial of the $A_{i,z}$'s, $Q_{0,z}^{(1)}$ cannot be an irreducible q -character. By Theorem 2.2 (i), we have $\chi_q(L_{1,z}^-) = Q_{1,z}^{(1)}$. Since $N_1 = 2$, Theorem 2.2 (ii) is not applicable.

4.4. **Series $\mathcal{Q}_{J,z}^{(a)}$.** In this subsection we give the definition of $\mathcal{Q}_{J,z}^{(a)}$ where $J \in \mathcal{J}^a$ and $(\mathfrak{g}, a) \neq (C_n^{(1)}, n), (D_n^{(1)}, n-1), (D_n^{(1)}, n)$. We recall that $q_1 = q^{1/2}$ for $C_n^{(1)}$ and $q_1 = q$ in the other cases.

We fix our convention about the indices as follows. For $J = (j_1, \dots, j_a) \in \mathcal{J}^a$, we denote by \tilde{J} the underlying set $\{j_1, \dots, j_a\} \subseteq \mathcal{J}$. Set further

$$(4.12) \quad \bar{J} = (\bar{j}_a, \dots, \bar{j}_1),$$

$$(4.13) \quad J^* = (i_1, \dots, i_b), \text{ with } J^* = \mathcal{J} \setminus \tilde{J} \text{ and } i_1 < \dots < i_b.$$

We say J is increasing if $j_1 < \dots < j_a$.

For an element $J = (j_1, \dots, j_a) \in \mathcal{J}^a$, we define $\mathcal{Q}_{J,z}^{(a)}$ by

$$(4.14) \quad \mathcal{Q}_{J,z}^{(a)} = \det \left(\mathcal{Q}_{j_\nu, q_1^{-a-1+2\mu}z}^{(1)} x_{j_\nu}^{\mu-1} \right)_{\mu, \nu=1}^a \\ = \begin{vmatrix} \mathcal{Q}_{j_1, q_1^{1-a}z}^{(1)} x_{j_1}^0 & \mathcal{Q}_{j_2, q_1^{1-a}z}^{(1)} x_{j_2}^0 & \dots & \mathcal{Q}_{j_a, q_1^{1-a}z}^{(1)} x_{j_a}^0 \\ \mathcal{Q}_{j_1, q_1^{3-a}z}^{(1)} x_{j_1}^1 & \mathcal{Q}_{j_2, q_1^{3-a}z}^{(1)} x_{j_2}^1 & \dots & \mathcal{Q}_{j_a, q_1^{3-a}z}^{(1)} x_{j_a}^1 \\ \vdots & \vdots & \ddots & \vdots \\ \mathcal{Q}_{j_1, q_1^{a-1}z}^{(1)} x_{j_1}^{a-1} & \mathcal{Q}_{j_2, q_1^{a-1}z}^{(1)} x_{j_2}^{a-1} & \dots & \mathcal{Q}_{j_a, q_1^{a-1}z}^{(1)} x_{j_a}^{a-1} \end{vmatrix}.$$

We set $\mathcal{Q}_{\emptyset, z}^{(0)} = 0$.

Note that if the entries of $J \in \mathcal{J}^a$ are not distinct (i.e. $\#J < a$), then $\mathcal{Q}_{J,z}^{(a)} = 0$.

Define

$$\bar{\mathcal{Q}}_J^{(a)} = \varpi(\mathcal{Q}_{J,z}^{(a)}).$$

Hence by the definition of $\mathcal{Q}_{J,z}^{(a)}$, it is easy to get

$$(4.15) \quad \bar{\mathcal{Q}}_J^{(a)} = \prod_{\substack{i \in J, j \in J^*, \\ i < j}} f_{i,j} \prod_{\substack{i, j \in J, \\ i < j}} ((x_j - x_i) f_{i,j}).$$

For an increasing element $J = (j_1, \dots, j_a) \in \mathcal{J}^a$, set

$$(4.16) \quad \Phi_{J,z} = \prod_{k=1}^a \Phi_{j_k, q_1^{a+1-2k}z},$$

where $\Phi_{j,z}$ are defined in (4.4)–(4.11) with $\mathcal{Y}_{i,z}$ being given in (4.1). Let $L(\Phi_{J,z})$ be the unique irreducible $U_q(\mathfrak{b})$ -module with the highest ℓ -weight $\Phi_{J,z}$. Except in the case $\mathfrak{g} = C_n^{(1)}$, we expect that the q -character of $L(\Phi_{J,z})$ is given by

$$\chi_q(L(\Phi_{J,z})) = \prod_{\substack{i \in J, j \in J^*, \\ i > j}} f_{j,i} \prod_{\substack{i, j \in J, \\ i < j}} (-x_i^{-1}) \cdot \mathcal{Q}_{J,z}^{(a)}.$$

In particular, from (4.15) we expect that

$$\chi(L(\Phi_{J,z})) = \prod_{\substack{(i,j) \in (J \times J^*) \cup (J^* \times J), \\ i < j}} f_{i,j} \times \begin{cases} 1 & (\mathfrak{g} = A_n^{(1)}), \\ \prod_{\substack{i, \bar{i} \in J, \\ i < \bar{i}}} (1 + x_i^{-1}) & (\mathfrak{g} = B_n^{(1)}), \\ \prod_{\substack{i, \bar{i} \in J, \\ i < \bar{i}}} (1 - x_i^{-2}) & (\mathfrak{g} = D_n^{(1)}). \end{cases}$$

In the case of $\mathfrak{g} = C_n^{(1)}$, we expect the same to be true for $\mathcal{Q}_{J,z}^{(a)}$ if $j_1, \dots, j_a \neq 0$. As noted already in Example 4.3, the series $\mathcal{Q}_{0,z}^{(1)}$ does not correspond to an irreducible q -character. We expect rather that $\widehat{\mathcal{Q}}_{0,z}^{(1)}, \widehat{\mathcal{Q}}_{0,z}^{(1)}$ correspond to irreducible q -characters. See also Remark 5.9.

4.5. **Series $\mathcal{R}_{\varepsilon,z}^{(n)}$ for the spin node: case $C_n^{(1)}$.** In this subsection, we introduce another set of series $\mathcal{R}_{\varepsilon,z}^{(n)}$ for the spin node. Unlike the series $\mathcal{Q}_{J,z}^{(a)}$, in general we do not know the explicit formulas for them. We define $\mathcal{R}_{\varepsilon,z}^{(n)}$ by the q -characters of the irreducible $U_q(\mathfrak{b})$ module $L(M_{\varepsilon,z})$, and give a rule to determine the highest ℓ -weight $M_{\varepsilon,z}$ which is a monomial in $\mathbb{C}[\mathcal{Y}_{i,z}^{\pm 1}]_{1 \leq i \leq n, z \in \mathbb{C}^\times}$.

Exhibiting the n -dependence explicitly, let P_n be the weight lattice of the simple Lie algebra of type C_n . Set

$$E_n = \{\varepsilon = (\varepsilon_1, \dots, \varepsilon_n) \mid \varepsilon_i = \pm 1, 1 \leq i \leq n\}.$$

We define two weight functions

$$(4.17) \quad w_1 : E_n \longrightarrow P_n, \quad \varepsilon \longmapsto \frac{1}{2} \sum_{i=1}^n \varepsilon_i \epsilon_i,$$

$$(4.18) \quad w_2 : \{J \mid J \in \mathcal{J}^a (1 \leq a \leq \#\mathcal{J}), J \subseteq \mathcal{J}\} \longrightarrow P_n, \quad (j_1, \dots, j_a) \longmapsto \sum_{k=1}^a \operatorname{sgn}(j_k) \epsilon_{j_k},$$

where $\epsilon_0 = 0$, and for $j \in \mathcal{J}$ we set

$$\operatorname{sgn}(j) = \begin{cases} + & j \preceq n, \\ 0 & j = 0, \\ - & j \succeq \bar{n}. \end{cases}$$

Borrowing an idea from [3], we introduce monomials $M_{\varepsilon,z}$ in $\mathbb{C}[\mathcal{Y}_{i,z}^{\pm 1}]_{1 \leq i \leq n, z \in \mathbb{C}^\times}$ inductively as follows. Define two operators τ^y, τ_c^z by

$$\tau^y : \mathcal{Y}_{i,z} \rightarrow \mathcal{Y}_{i+1,z}, \quad \tau_c^z : \mathcal{Y}_{i,z} \rightarrow \mathcal{Y}_{i,q^c z}.$$

For $n = 1$ we set $M_{+,z} = \mathcal{Y}_{1,z}$ and $M_{-,z} = \mathcal{Y}_{1,q^2 z}^{-1}$. In the general case we set

$$(4.19) \quad M_{(++)z} = \tau^y(M_{(+\xi),z}), \quad M_{(+-)z} = \mathcal{Y}_{1,q^{\frac{n}{2}} z} \tau^y(M_{(-\xi),z}),$$

$$(4.20) \quad M_{(-+)z} = \mathcal{Y}_{1,q^{\frac{n+2}{2}} z}^{-1} \tau_1^z \tau^y(M_{(+\xi),z}), \quad M_{(--)z} = \tau_1^z \tau^y(M_{(-\xi),z}),$$

where $\xi \in E_{n-2}$.

Now let $\varepsilon \in E_n$ with $w_1(\varepsilon) = \frac{1}{2} \left(\sum_{k=1}^s \epsilon_{\varphi_k} - \sum_{k=1}^t \epsilon_{\psi_k} \right)$, and consider the simple module $L(M_{\varepsilon,z})$. We define $\mathcal{R}_{\varepsilon,z}^{(n)}$ and $\bar{\mathcal{R}}_{\varepsilon}^{(n)}$ through the q -character and the character as follows.

$$\begin{aligned} \mathcal{R}_{\varepsilon,z}^{(n)} &= \chi_q(L(M_{\varepsilon,z})) \prod_{1 \leq k \leq l \leq t} (1 - x_{\psi_k} x_{\psi_l}) \prod_{\substack{1 \leq k \leq s, 1 \leq l \leq t, \\ \varphi_k > \psi_l}} (1 - x_{\varphi_k}^{-1} x_{\psi_l}), \\ \bar{\mathcal{R}}_{\varepsilon}^{(n)} &= \chi(L(M_{\varepsilon,z})) \prod_{1 \leq k \leq l \leq t} (1 - x_{\psi_k} x_{\psi_l}) \prod_{\substack{1 \leq k \leq s, 1 \leq l \leq t, \\ \varphi_k > \psi_l}} (1 - x_{\varphi_k}^{-1} x_{\psi_l}). \end{aligned}$$

For the latter we have the following guess.

Conjecture 4.4. For $\varepsilon = (\varepsilon_1, \dots, \varepsilon_n) \in E_n$, we have

$$\chi(L(M_{\varepsilon,z})) = \prod_{1 \leq i \leq n} \frac{1}{1 - x_i^{-2\varepsilon_i}} \prod_{1 \leq i < j \leq n} \frac{1}{1 - x_i^{-\varepsilon_i} x_j^{-\varepsilon_j}}.$$

In the special case $n = 2$, one can obtain the series for $C_2^{(1)}$ from that of $B_2^{(1)}$. Let $\mathcal{Q}_{i,z}^{(1)} (1 \leftrightarrow 2)$ stand for the series for $B_2^{(1)}$ with $a = 1$ given in Example 4.3, wherein we interchange $\mathcal{Y}_{1,z}$ with $\mathcal{Y}_{2,z}$ and $A_{1,z}$ with $A_{2,z}$. Then we have

$$(4.21) \quad \mathcal{R}_{(++)z}^{(2)} = \mathcal{Q}_{1,z}^{(1)} (1 \leftrightarrow 2), \quad \mathcal{R}_{(+-)z}^{(2)} = \mathcal{Q}_{2,z}^{(1)} (1 \leftrightarrow 2),$$

$$(4.22) \quad \mathcal{R}_{(-+)z}^{(2)} = \mathcal{Q}_{2,z}^{(1)} (1 \leftrightarrow 2), \quad \mathcal{R}_{(--)z}^{(2)} = \mathcal{Q}_{1,z}^{(1)} (1 \leftrightarrow 2).$$

4.6. **Series $\mathcal{R}_{\varepsilon,z}^{(n-1)}$, $\mathcal{R}_{\varepsilon,z}^{(n)}$ for the spin nodes: case $D_n^{(1)}$.** For the two spin nodes of $D_n^{(1)}$, we follow the procedure for $C_n^{(1)}$.

We use the same weight functions (4.17), (4.18) as for $C_n^{(1)}$. We also introduce two subsets of E_n ,

$$E_{n,\varsigma} = \{\varepsilon \in E_n \mid \#\{\varepsilon_i = -\} = \varsigma \pmod{2}\}, \quad (\varsigma = 0, 1).$$

We define monomials $M_{\varepsilon,z}$ inductively by

$$(4.23) \quad M_{(++)z} = \tau^y(M_{(+)z}), \quad M_{(+-)z} = y_{1,q^{n-2}z} \tau^y(M_{(-)z}),$$

$$(4.24) \quad M_{(-+)z} = y_{n,q^n z} \tau_z^y \tau^y(M_{(+)z}), \quad M_{(--)z} = \tau_z^y \tau^y(M_{(-)z}),$$

with $\xi \in E_{n-2}$ and the initial values

$$\begin{aligned} M_{(++)z} &= y_{2,z}, & M_{(--)z} &= y_{2,q^2 z}^{-1}, \\ M_{(+-)z} &= y_{1,z}, & M_{(-+)z} &= y_{1,q^2 z}^{-1}. \end{aligned}$$

For $\varepsilon \in E_{n,\varsigma}$ ($\varsigma = 0, 1$) with $w_1(\varepsilon) = \frac{1}{2} \left(\sum_{k=1}^s \varepsilon_{\varphi_k} - \sum_{k=1}^t \varepsilon_{\psi_k} \right)$, we define

$$\begin{aligned} \mathcal{R}_{\varepsilon,z}^{(n-\varsigma)} &= \chi_q(L(M_{\varepsilon,z})) \prod_{1 \leq k < l \leq t} (1 - x_{\psi_k} x_{\psi_l}) \prod_{\substack{1 \leq k \leq s, 1 \leq l \leq t, \\ \varphi_k > \psi_l}} (1 - x_{\varphi_k}^{-1} x_{\psi_l}), \\ \bar{\mathcal{R}}_{\varepsilon}^{(n-\varsigma)} &= \chi(L(M_{\varepsilon,z})) \prod_{1 \leq k < l \leq t} (1 - x_{\psi_k} x_{\psi_l}) \prod_{\substack{1 \leq k \leq s, 1 \leq l \leq t, \\ \varphi_k > \psi_l}} (1 - x_{\varphi_k}^{-1} x_{\psi_l}). \end{aligned}$$

Conjecture 4.5. For $\varepsilon = (\varepsilon_1, \dots, \varepsilon_n) \in E_{n,\varsigma}$ ($\varsigma = 0, 1$), we have

$$\chi(L(M_{\varepsilon,z})) = \prod_{1 \leq i < j \leq n} \frac{1}{1 - x_i^{-\varepsilon_i} x_j^{-\varepsilon_j}}.$$

Although we do not have formulas for the series $\mathcal{R}_{\varepsilon,z}^{(n-\varsigma)}$ ($\varsigma = 0, 1$) for general n , in the special case $D_4^{(1)}$, similarly to $\mathcal{Q}_{i,z}^{(1)}$, we define $\mathcal{R}_{\varepsilon,z}^{(3)}$ (resp. $\mathcal{R}_{\varepsilon,z}^{(4)}$) as certain limits of the q -characters of the KR modules $W_{m,z}^{(3)}$ (resp. $W_{m,z}^{(4)}$). More precisely, one can obtain $\mathcal{R}_{\varepsilon,z}^{(3)}$ (resp. $\mathcal{R}_{\varepsilon,z}^{(4)}$) from $\mathcal{Q}_{i,z}^{(1)}$ by exchanging $y_{1,z}$ with $y_{3,z}$ (resp. $y_{4,z}$) and $A_{1,z}$ with $A_{3,z}$ (resp. $A_{4,z}$). For example, we have

$$\begin{aligned} \mathcal{Q}_{2,z}^{(1)} &= y_{2,q^5 z}^{-1} y_{1,q^4 z} \sum_{k \geq 0} \prod_{j=1}^k A_{1,q^{6-2j} z}^{-1}, & \mathcal{Q}_{1,z}^{(1)} &= y_{1,q^6 z}^{-1}, \\ \mathcal{R}_{(---+),z}^{(3)} &= y_{2,q^5 z}^{-1} y_{3,q^4 z} \sum_{k \geq 0} \prod_{j=1}^k A_{3,q^{6-2j} z}^{-1}, & \mathcal{R}_{(----),z}^{(3)} &= y_{3,q^6 z}^{-1}, \\ \mathcal{R}_{(---+),z}^{(4)} &= y_{2,q^5 z}^{-1} y_{4,q^4 z} \sum_{k \geq 0} \prod_{j=1}^k A_{4,q^{6-2j} z}^{-1}, & \mathcal{R}_{(----),z}^{(4)} &= y_{4,q^6 z}^{-1}, \end{aligned}$$

and so on.

5. POLYNOMIAL RELATIONS

In this section, we shall give the main results of this paper. We propose polynomial relations among the series defined in Section 4, which are expected from the ψ -system obtained in Section 3.

We prepare some notation which we will use below.

For an element $j \in \mathcal{J}$, an element $J = (j_1, \dots, j_a) \in \mathcal{J}^a$ and a positive integer k with $1 \leq k \leq a$, we set

$$\lambda_J = \sum_{l=1}^a \lambda_{j_l}, \quad x_J = \prod_{l=1}^a x_{j_l},$$

$$J - j_k = (j_1, \dots, \hat{j}_k, \dots, j_a), \quad (j, J) = (j, j_1, \dots, j_a).$$

Let $\{v_i\}_{i \in \tilde{\mathcal{J}}}$ be the standard basis of the vector representation $V^{(1)}$ of ${}^L\mathfrak{g}$ (cf. Appendix A). Here $\tilde{\mathcal{J}} = \{1, \dots, n, 0, \bar{0}, \bar{n}, \dots, \bar{1}\}$ for $C_n^{(1)}$ and $\tilde{\mathcal{J}} = \mathcal{J}$ for the other cases. For $J = (j_1, \dots, j_a) \in \tilde{\mathcal{J}}^a$, we set $v_J = v_{j_1} \wedge v_{j_2} \wedge \dots \wedge v_{j_a}$.

5.1. Plücker-type relations. We begin by writing down the ψ -system and the corresponding relations for $\mathcal{Q}_{J,z}^{(a)}$ for all nodes excepting the last (or the second last for $D_n^{(1)}$). In fact they are just the consequence of the fact that the $\mathcal{Q}_{J,z}^{(a)}$ are defined by determinants of the $\mathcal{Q}_{i,z}^{(1)}$'s. So they are rather definitions, and we write them down just for uniformity reasons. As we shall see later, the only non-trivial relations come from the last (or the second last) node.

For comparison, we start with the ψ -system. Let $J_1 = (i_1, \dots, i_a)$, $J_2 = (j_1, \dots, j_a)$ be two elements of $\tilde{\mathcal{J}}^a$. The embedding ι (B.1) of ${}^L\mathfrak{g}$ -modules is explicitly given by

$$\iota : \bigwedge^2 V_{\frac{1}{2t}}^{(a)} \hookrightarrow V^{(a-1)} \otimes V^{(a+1)}$$

$$v_{J_1} \wedge v_{J_2} \mapsto \frac{1}{2} \left(\sum_{k=1}^a (-1)^{k-1} v_{J_1 - i_k} \otimes v_{(i_k, J_2)} - \sum_{k=1}^a (-1)^{k-1} v_{J_2 - j_k} \otimes v_{(j_k, J_1)} \right).$$

Therefore (3.12), (3.14), (3.16) and (3.19) imply

$$(5.1) \quad \mathcal{Q}_{J_1, -\frac{1}{2t}}^{(a)}(E) \mathcal{Q}_{J_2, \frac{1}{2t}}^{(a)}(E) - \mathcal{Q}_{J_1, \frac{1}{2t}}^{(a)}(E) \mathcal{Q}_{J_2, -\frac{1}{2t}}^{(a)}(E) \omega^{\lambda_{J_1} - \lambda_{J_2}}$$

$$= \sum_{k=1}^a (-1)^k \mathcal{Q}_{J_2 - j_k}^{(a-1)}(E) \mathcal{Q}_{(j_k, J_1)}^{(a+1)}(E) \omega^{-\lambda_{j_k}}.$$

The corresponding relations for $\mathcal{Q}_{J,z}^{(a)}$ read as follows.

Proposition 5.1. *For two increasing elements $J_1 = (i_1, i_2, \dots, i_a)$ and $J_2 = (j_1, j_2, \dots, j_a)$ of \mathcal{J}^a , we have*

$$(5.2) \quad \mathcal{Q}_{J_1, q_1^{-1}z}^{(a)} \mathcal{Q}_{J_2, q_1 z}^{(a)} - \mathcal{Q}_{J_1, q_1 z}^{(a)} \mathcal{Q}_{J_2, q_1^{-1}z}^{(a)} \frac{x_{J_1}}{x_{J_2}} = \sum_{k=1}^a (-1)^k \mathcal{Q}_{J_2 - j_k, z}^{(a-1)} \mathcal{Q}_{(j_k, J_1), z}^{(a+1)} x_{j_k}^{-1}.$$

This is an immediate consequence of the Sylvester identity (see e.g. [26, pp.108]). Note that (5.1) and (5.2) has the same form under the identification

$$\mathcal{Q}_J^{(a)}(E) \leftrightarrow \mathcal{Q}_{J,z}^{(a)}, \quad E \leftrightarrow z, \quad \Omega^{\frac{1}{2t}} \leftrightarrow q_1, \quad \omega^{\lambda_i} \leftrightarrow x_i.$$

5.2. Wronskian identity for $A_n^{(1)}$. For type $A_n^{(1)}$, the non-trivial relation is the following ‘Wronskian identity’.

Theorem 5.2. *Let $\mathfrak{g} = A_n^{(1)}$. We have*

$$\det \left(\mathcal{Q}_{\nu, q^{n+2-2\mu}z}^{(1)} x_{\nu}^{-\mu+\nu} \right)_{\mu, \nu=1}^{n+1} = 1.$$

Note that the left-hand side is just $\mathcal{Q}_{J,z}^{(n+1)} (-1)^{(n+1)n/2} \times \prod_{b=1}^n x_b^{n+1-b}$ with $J = (1, 2, \dots, n+1)$. We prove Theorem 5.2 in Appendix B.

5.3. Polynomial relations related to the last node: case $B_n^{(1)}$. In this subsection we let $\mathfrak{g} = B_n^{(1)}$. We give conjectural relations which correspond to the last identity (3.15) of the ψ -system. We first give identities for the connection coefficients $Q_J^{(a)}(E)$.

In this subsection, for an element $j \in \mathcal{J}$ and $J = (j_1, \dots, j_a) \in \mathcal{J}^a$ we define

$$c(j) = \begin{cases} (-1)^j & j \preceq n \\ (-1)^{\bar{j}-1} & j \succeq \bar{n} \end{cases}, \quad c(J) = \prod_{k=1}^a c(j_k).$$

We also denote by $\sigma(J, J^*)$ the signature of the permutation $(1, \dots, n, \bar{n}, \dots, \bar{1}) \mapsto (J, J^*)$.

The bilinear form $\langle \cdot, \cdot \rangle : V \times V_{\frac{1}{2}} \rightarrow \mathbb{C}$ given by $\langle v_i, v_j \rangle = c(i)\delta_{i,\bar{j}}$ is $L_{\mathfrak{g}}$ -invariant. From this it is easy to see that we have the following $L_{\mathfrak{g}}$ -module isomorphism

$$V_{\frac{1}{2}}^{(a)} \simeq V^{(2n-a)} \quad (a = 1, \dots, n) \\ v_J \mapsto c(J)\sigma(J, J^*)v_{J^*},$$

which implies

$$(5.3) \quad Q_{J, \frac{1}{2}}^{(a)}(E) = c(J)\sigma(J, J^*)Q_{J^*}^{(2n-a)}(E).$$

Let $J_1 = (i_1, \dots, i_n)$, $J_2 = (j_1, \dots, j_n)$ be two elements of \mathcal{J}^n . Then the relation (5.1) together with (5.3) imply

$$(5.4) \quad Q_{J_1^*}^{(n)}(E)Q_{J_2, \frac{1}{2}}^{(n)}(E)c(\overline{J_1^*})\sigma(\overline{J_1^*}, \overline{J_1})\omega^{-\lambda_{J_1}} - Q_{J_2^*}^{(n)}(E)Q_{J_1, \frac{1}{2}}^{(n)}(E)c(\overline{J_2^*})\sigma(\overline{J_2^*}, \overline{J_2})\omega^{-\lambda_{J_2}} \\ = \sum_{k=1}^n (-1)^k Q_{J_2 - j_k}^{(n-1)}(E)Q_{J_1^* - \bar{j}_k, \frac{1}{2}}^{(n-1)}(E)c(\overline{(j_k, J_1)^*})\sigma(\overline{(j_k, J_1)^*}, \overline{(j_k, J_1)})\omega^{-\lambda_{J_1} - \lambda_{j_k}}.$$

For a given element $i \in \mathcal{J}$ and $J \in \mathcal{J}^a$, let

$$(5.5) \quad X_J = \prod_{\substack{\mu, \nu \in J, \\ \mu < \nu}} (-x_{\mu}^{\frac{1}{2}\delta_{\mu, \nu}-1}), \quad x_{i, J} = \prod_{\substack{j \in J, \\ i < j}} (-x_i^{1-\frac{1}{2}\delta_{j, i}}) \prod_{\substack{j \in J^*, \\ j < i}} (x_j^{\frac{1}{2}\delta_{i, j}-1}).$$

The counterpart of (5.4) for the series $Q_{J, z}^{(a)}$ is given by the following conjecture.

Conjecture 5.3. *Let $J_1 = (i_1, \dots, i_n)$, $J_2 = (j_1, \dots, j_n) \in \mathcal{J}^n$ be increasing. Then we have*

$$Q_{J_1^*, z}^{(n)}Q_{J_2, qz}^{(n)}X_{J_1^*}X_{J_2}x_{J_1}^{-\frac{1}{2}} - Q_{J_2^*, z}^{(n)}Q_{J_1, qz}^{(n)}X_{J_2^*}X_{J_1}x_{J_2}^{-\frac{1}{2}} \\ = - \sum_{k=1}^n Q_{J_2 - j_k, z}^{(n-1)}Q_{J_1^* - \bar{j}_k, qz}^{(n-1)}X_{J_2 - j_k}X_{J_1^* - \bar{j}_k} \frac{x_{j_k, J_1}^{-\frac{1}{2}}}{x_{j_k, J_2}} x_{J_1}^{-\frac{1}{2}} x_{j_k}^{-\frac{1}{2}}.$$

At the moment we do not have a proof of these identities in general. When $n = 2$, there are 15 pairs of choices of J_1 and J_2 , and we have verified them by direct computations. To save space we do not write the proofs here. We have also checked them for $n = 3$ by Mathematica 5.0 up to certain degree, counting the degree of $A_{a, q^k z}$ to be k . In Appendix D, we show that the identity holds when specialized to the characters.

Remark 5.4. It is worth noting that, although the connection coefficients satisfy the linear relation (5.3), there are no analogs for the series $Q_{J, z}^{(a)}$. For example, in the simplest case $B_2^{(1)}$, a direct computation gives $Q_{(2\bar{2}\bar{1}), z}^{(3)} = -x_2 y_{1, q^4 z}^{-1} y_{2, q^2 z}^{-1} y_{2, q^3 z} - y_{1, q^2 z}^{-1} y_{2, q^3 z} y_{2, q^4 z}^{-1}$, while $Q_{\bar{1}, z}^{(1)} = y_{1, q^3 z}^{-1}$, and there is no function $f(x_1, x_2)$ of x_1, x_2 such that $Q_{(2\bar{2}\bar{1}), z}^{(3)} = f(x_1, x_2)Q_{\bar{1}, z}^{(1)}$. So only the quadratic relations survive.

5.4. **Polynomial relations related to spin node: case $C_n^{(1)}$.** In this subsection, $\mathfrak{g} = C_n^{(1)}$. We shall give conjectural relations for the series $\mathcal{Q}_{J,z}^{(a)}$ and $\mathcal{R}_{\varepsilon,z}^{(n)}$.

For an increasing element $J \in \mathcal{J}^a$ we can write

$$\begin{aligned} J &= \{\alpha_1, \dots, \alpha_m, \bar{\beta}_1, \dots, \bar{\beta}_r, \sigma_1, \dots, \sigma_u, \bar{\sigma}_1, \dots, \bar{\sigma}_u\} \cup (J \cap \{0\}), \\ J^* &= \{\bar{\alpha}_1, \dots, \bar{\alpha}_m, \beta_1, \dots, \beta_r, \eta_1, \dots, \eta_v, \bar{\eta}_1, \dots, \bar{\eta}_v\} \cup (J^* \cap \{0\}), \end{aligned}$$

where $\alpha_i, \beta_i, \sigma_i, \eta_i$ are mutually distinct elements of $\{1, \dots, n\}$. Let $\sigma = \{\sigma_1, \dots, \sigma_u\}$, $\eta = \{\eta_1, \dots, \eta_v\}$. Note that

$$v = \begin{cases} u+2 & 0 \in J \text{ and } J \in \mathcal{J}^{n-1} \\ u+1 & 0 \notin J \text{ and } J \in \mathcal{J}^{n-1}, \text{ or } 0 \in J \text{ and } J \in \mathcal{J}^n. \\ u & 0 \notin J \text{ and } J \in \mathcal{J}^n \end{cases}$$

Let $J \in \mathcal{J}^a$ be increasing. We introduce the following condition for $\varepsilon, \varepsilon' \in E_n$:

Condition C_J :

- $w_1(\varepsilon) + w_1(\varepsilon') = w_2(J)$,
- if $w_1(\varepsilon) - w_1(\varepsilon') = \sum_{k=1}^s \varepsilon_{\gamma_k} - \sum_{k=1}^t \varepsilon_{\delta_k}$, then $\sigma \subseteq \gamma$.

Here w_1, w_2 are the weight functions defined by (4.17), (4.18) and $\gamma = \{\gamma_1, \dots, \gamma_s\}$, $\delta = \{\delta_1, \dots, \delta_t\}$. It is easy to see $\sigma \cup \eta = \gamma \cup \delta$.

In the following Conjectures 5.5–5.8, $\mathfrak{g} = C_n^{(1)}$, σ, η are disjoint subsets of $\{1, 2, \dots, n\}$, the sum is taken over all $\varepsilon, \varepsilon' \in E_n$ satisfying Condition C_J where γ, δ are determined there.

Conjecture 5.5. Suppose $J \in \mathcal{J}^{n-1}$ is increasing and $0 \notin J$. Then we have

$$\begin{aligned} & (-1)^{\frac{u(u-1)}{2}} \mathcal{Q}_{J,z}^{(n-1)} x_{\sigma}^{-2u} x_{\eta}^{2-u} \prod_{\substack{i,j \in J, \\ i \prec j}} (-x_i^{-1}) \prod_{\substack{1 \leq k \leq u, 1 \leq l \leq u+1, \\ \sigma_k \succ \eta_l}} \left(-\frac{x_{\sigma_k}}{x_{\eta_l}} \right) \\ &= \sum (-1)^{\frac{t(t+1)}{2}} \mathcal{R}_{\varepsilon,qz}^{(n)} \mathcal{R}_{\varepsilon',q^{-1}z}^{(n)} x_{\gamma}^{2-2u} x_{\delta}^{1-t} \prod_{\substack{1 \leq k \leq s, 1 \leq l \leq t, \\ \gamma_k \succ \delta_l}} \left(-\frac{x_{\gamma_k}}{x_{\delta_l}} \right). \end{aligned}$$

Conjecture 5.6. Suppose $J \in \mathcal{J}^{n-1}$ is increasing and $0 \in J$. Then we have

$$\begin{aligned} & (-1)^{\frac{u(u-1)}{2}} \mathcal{Q}_{J,z}^{(n-1)} x_{\sigma}^{-2u} x_{\eta}^{2-u} \prod_{\substack{i,j \in J, \\ i \prec j}} (-x_i^{-1}) \prod_{\substack{1 \leq k \leq u, 1 \leq l \leq u+2, \\ \sigma_k \succ \eta_l}} \left(-\frac{x_{\sigma_k}}{x_{\eta_l}} \right) \\ &= \sum (-1)^{\frac{t(t-1)}{2}} \mathcal{R}_{\varepsilon,qz}^{(n)} \mathcal{R}_{\varepsilon',q^{-1}z}^{(n)} x_{\gamma}^{2-2u} x_{\delta}^{2-t} \prod_{\substack{1 \leq k \leq s, 1 \leq l \leq t, \\ \gamma_k \succ \delta_l}} \left(-\frac{x_{\gamma_k}}{x_{\delta_l}} \right). \end{aligned}$$

Conjecture 5.7. Suppose $J \in \mathcal{J}^n$ is increasing and $0 \notin J$. Then we have

$$\begin{aligned} & (-1)^{\frac{u(u-1)}{2}} \mathcal{Q}_{J,z}^{(n)} x_{\sigma}^{-2u} x_{\eta}^{1-u} \prod_{\substack{i,j \in J, \\ i \prec j}} (-x_i^{-1}) \prod_{\substack{1 \leq k \leq u, 1 \leq l \leq u, \\ \sigma_k \succ \eta_l}} \left(-\frac{x_{\sigma_k}}{x_{\eta_l}} \right) \\ &= \sum (-1)^{\frac{s(s-1)}{2}} \mathcal{R}_{\varepsilon,q^{\frac{1}{2}}z}^{(n)} \mathcal{R}_{\varepsilon',q^{-\frac{1}{2}}z}^{(n)} x_{\gamma}^{1-2u} x_{\delta}^{-t} \prod_{\substack{1 \leq k \leq s, 1 \leq l \leq t, \\ \gamma_k \succ \delta_l}} \left(-\frac{x_{\gamma_k}}{x_{\delta_l}} \right). \end{aligned}$$

Conjecture 5.8. Suppose $J \in \mathcal{J}^n$ is increasing and $0 \in J$. Then we have

$$\begin{aligned} & (-1)^{\frac{u(u-1)}{2}} Q_{J,z}^{(n)} x_{\sigma}^{-2u} x_{\eta}^{1-u} \prod_{\substack{i,j \in J, \\ i < j}} (-x_i^{-1}) \prod_{\substack{1 \leq k \leq u, 1 \leq l \leq u+1, \\ \sigma_k > \eta_l}} \left(-\frac{x_{\sigma_k}}{x_{\eta_l}} \right) \\ &= \sum (-1)^{\frac{s(s-1)}{2}} \mathcal{R}_{\varepsilon, q^{\frac{1}{2}} z}^{(n)} \mathcal{R}_{\varepsilon', q^{-\frac{1}{2}} z}^{(n)} x_{\gamma}^{1-2u} x_{\delta}^{1-t} \prod_{\substack{1 \leq k \leq s, 1 \leq l \leq t, \\ \gamma_k > \delta_l}} \left(-\frac{x_{\gamma_k}}{x_{\delta_l}} \right). \end{aligned}$$

For $n = 2$ one can verify these relations using the relation between $C_2^{(1)}$ and $B_2^{(1)}$ (see (4.21), (4.22)). For general n , we show in Appendix D the validity of the identities when specialized to the ordinary characters $\bar{Q}_J^{(a)}$ and $\bar{\mathcal{R}}_{\varepsilon}^{(n)}$.

Remark 5.9. In the conjectured identities given above, $\widehat{Q}_{0,z}^{(1)}$ and $\widehat{Q}_{\bar{0},z}^{(1)}$ enter only through the sum $Q_{0,z}^{(1)} = \widehat{Q}_{0,z}^{(1)} + \widehat{Q}_{\bar{0},z}^{(1)}$. However, since the weight 0 of $V^{(1)}$ has multiplicity 2, it is more natural to consider $\widehat{Q}_{0,z}^{(1)}$ and $\widehat{Q}_{\bar{0},z}^{(1)}$ separately. The ψ -system also suggests that there are identities involving them separately. This is indeed the case for $n = 2$ where we have, for example,

$$\begin{aligned} & Q_{1,q^{-\frac{1}{2}}z}^{(1)} \widehat{Q}_{0,q^{\frac{1}{2}}z}^{(1)} - x_1 \widehat{Q}_{0,q^{-\frac{1}{2}}z}^{(1)} Q_{1,q^{\frac{1}{2}}z}^{(1)} = -\frac{x_1}{x_2} \mathcal{R}_{(++)^{-\frac{1}{2}}z}^{(2)} \mathcal{R}_{(+-)^{\frac{1}{2}}z}^{(2)}, \\ & Q_{1,q^{-\frac{1}{2}}z}^{(1)} \widehat{Q}_{0,q^{\frac{1}{2}}z}^{(1)} - x_1 \widehat{Q}_{0,q^{-\frac{1}{2}}z}^{(1)} Q_{1,q^{\frac{1}{2}}z}^{(1)} = -x_1 \mathcal{R}_{(+-)^{-\frac{1}{2}}z}^{(2)} \mathcal{R}_{(++)^{\frac{1}{2}}z}^{(2)}. \end{aligned}$$

On the other hand, if we define $Q_{J,z}^{(a)}$ as determinants treating 0 and $\bar{0}$ independently, then when J contains both 0 and $\bar{0}$ we have $\bar{Q}_J^{(a)} = 0$ because $x_0 = x_{\bar{0}} = 1$. Hence $Q_{J,z}^{(a)}$ can not be explained as a q -character, so we must modify the working hypothesis. At the moment we do not know how to do that.

5.5. Polynomial relations related to spin nodes: case $D_n^{(1)}$. Finally we consider the case $\mathfrak{g} = D_n^{(1)}$. As in the case of $C_n^{(1)}$, for an increasing element $J \in \mathcal{J}^a$ we can write

$$\begin{aligned} J &= \{\alpha_1, \dots, \alpha_m, \bar{\beta}_1, \dots, \bar{\beta}_r, \sigma_1, \dots, \sigma_u, \bar{\sigma}_1, \dots, \bar{\sigma}_v\}, \\ J^* &= \{\bar{\alpha}_1, \dots, \bar{\alpha}_m, \beta_1, \dots, \beta_r, \eta_1, \dots, \eta_v, \bar{\eta}_1, \dots, \bar{\eta}_v\}, \end{aligned}$$

with mutually distinct $\alpha_i, \beta_i, \sigma_i, \eta_i \in \{1, \dots, n\}$. Let $\sigma = \{\sigma_1, \dots, \sigma_u\}$, $\eta = \{\eta_1, \dots, \eta_v\}$. Note that

$$v = \begin{cases} u+2 & J \in \mathcal{J}^{n-2} \\ u+1 & J \in \mathcal{J}^{n-1} \end{cases}.$$

Let $J \in \mathcal{J}^a$ be increasing. We introduce the following conditions for $\varepsilon, \varepsilon' \in E_n$:

Condition $D_{J,\varsigma}$:

- $w_1(\varepsilon) + w_1(\varepsilon') = w_2(J)$,
- if $w_1(\varepsilon) - w_1(\varepsilon') = \sum_{k=1}^s \varepsilon_{\gamma_k} - \sum_{k=1}^t \varepsilon_{\delta_k}$, then $\sigma \subseteq \gamma$, and $t \equiv r + \varsigma \pmod{2}$.

Here $\varsigma \in \{0, 1\}$.

Condition D_J :

- $w_1(\varepsilon) + w_1(\varepsilon') = w_2(J)$,
- if $w_1(\varepsilon) - w_1(\varepsilon') = \sum_{k=1}^s \varepsilon_{\gamma_k} - \sum_{k=1}^t \varepsilon_{\delta_k}$, then $\sigma \subseteq \gamma$, and $t \equiv r \pmod{2}$.

Here w_1, w_2 are the weight functions defined by (4.17), (4.18). We set $\gamma = \{\gamma_1, \dots, \gamma_s\}$, $\delta = \{\delta_1, \dots, \delta_t\}$. It is easy to see $\sigma \cup \eta = \gamma \cup \delta$.

In the following Conjectures 5.10–5.11, $\mathfrak{g} = D_n^{(1)}$, and σ and η are disjoint subsets of $\{1, \dots, n\}$.

Conjeture 5.10. Let J be an increasing element of \mathcal{J}^{n-2} , and let $\varsigma = 0, 1$. Then we have

$$\begin{aligned} & (-1)^{\frac{u(u-1)}{2}} \mathcal{Q}_{J,z}^{(n-2)} x_{\sigma}^{1-2u} x_{\eta}^{1-u} \prod_{\substack{i,j \neq i \in J, \\ i \prec j}} (-x_i^{-1}) \prod_{\substack{1 \leq k \leq u, 1 \leq l \leq u+2, \\ \sigma_k \succ \eta_l}} \left(-\frac{x_{\sigma_k}}{x_{\eta_l}} \right) \\ &= \sum (-1)^{\frac{s(s+1)}{2}} \mathcal{R}_{\varepsilon,qz}^{(n-\varsigma)} \mathcal{R}_{\varepsilon',q^{-1}z}^{(n-\varsigma)} x_{\gamma}^{-2u} x_{\delta}^{3-t} \prod_{\substack{1 \leq k \leq s, 1 \leq l \leq t, \\ \gamma_k \succ \delta_l}} \left(-\frac{x_{\gamma_k}}{x_{\delta_l}} \right), \end{aligned}$$

where the sum is taken over all $\varepsilon, \varepsilon' \in E_{n,\varsigma}$ satisfying Condition $D_{J,\varsigma}$, where γ and δ are determined there.

Conjeture 5.11. Let J be an increasing element of \mathcal{J}^{n-1} . Then we have

$$\begin{aligned} & (-1)^{\frac{u(u-1)}{2}} \mathcal{Q}_{J,z}^{(n-1)} x_{\sigma}^{-1-2u} x_{\eta}^{-u} \prod_{\substack{i,j \neq i \in J, \\ i \prec j}} (-x_i^{-1}) \prod_{\substack{1 \leq k \leq u, 1 \leq l \leq u+1, \\ \sigma_k \succ \eta_l}} \left(-\frac{x_{\sigma_k}}{x_{\eta_l}} \right) \\ &= \sum (-1)^{\frac{t(t-1)}{2}} \mathcal{R}_{\varepsilon,z}^{(n-1)} \mathcal{R}_{\varepsilon',z}^{(n)} x_{\gamma}^{-2u} x_{\delta}^{1-t} \prod_{\substack{1 \leq k \leq s, 1 \leq l \leq t, \\ \gamma_k \succ \delta_l}} \left(-\frac{x_{\gamma_k}}{x_{\delta_l}} \right), \end{aligned}$$

where the sum is taken over all $\varepsilon \in E_{n,1}$, $\varepsilon' \in E_{n,0}$ satisfying Condition D_J , where γ and δ are determined there.

For $n = 4$, we have checked by Mathematica 5.0 the conjectures hold up to certain degree. For general n , we prove in Appendix D the conjectures specialized to the ordinary characters $\bar{\mathcal{Q}}_J^{(a)}$ and $\bar{\mathcal{R}}_{\varepsilon}^{(n-1)}, \bar{\mathcal{R}}_{\varepsilon}^{(n)}$.

6. CONCLUSION

In this section, we give a summary of this paper. We have done the following things:

- (i) To each weight of the fundamental representation $V^{(a)}$, we associated a formal series $\mathcal{Q}_{J,z}^{(a)}$ or $\mathcal{R}_{\varepsilon,z}^{(a)}$. We expect that with some simple factors the formal series are q -characters of certain irreducible modules of $U_q(\mathfrak{b})$.
- (ii) Under suitable identifications, using relations for the connection coefficients implied by the ψ -system, we proposed the following conjecture relations for the series $\mathcal{Q}_{J,z}^{(a)}$ and $\mathcal{R}_{\varepsilon,z}^{(a)}$:
 - $\mathfrak{g} = A_n^{(1)}, B_n^{(1)}, C_n^{(1)}, D_n^{(1)}$, Proposition 5.1. This is the Plücker type relations which is in fact the defining relations of the series $\mathcal{Q}_{J,z}^{(a)}$;
 - $\mathfrak{g} = A_n^{(1)}$, Theorem 5.2. This is the Wronskian identity for the q -characters of $U_q(\mathfrak{b})$ to be proved in Appendix C.
 - $\mathfrak{g} = B_n^{(1)}$, Conjecture 5.3.
 - $\mathfrak{g} = C_n^{(1)}$, Conjectures 5.5-5.8.
 - $\mathfrak{g} = D_n^{(1)}$, Conjectures 5.10-5.11.

For the last three cases, we support our conjectures by checking the following. For the special cases $\mathfrak{g} = B_2^{(1)}$, $\mathfrak{g} = C_2^{(1)}$, we proved the conjectures by direct computations. For the cases $\mathfrak{g} = B_3^{(1)}$ and $\mathfrak{g} = D_4^{(1)}$, we checked the conjectures up to some degrees by Mathematica. When specialized to characters, the conjectured relations hold in all cases. This will be proved in Appendix D or E.

We hope we have presented reasonable grounds to suggest that the correspondence between the connection coefficients of certain differential equations and the q -characters of the Borel subalgebra $U_q(\mathfrak{b})$ supplies an effective way to find polynomial relations. Certainly this is only the first step, and more serious checks are desirable along with attempts toward proving these identities.

APPENDIX A. VECTOR REPRESENTATION OF ${}^L\mathfrak{g}$

Following [27, Appendix 2], we give an explicit realization of the vector representation of ${}^L\mathfrak{g}$. The symbol $E_{i,j}$ stands for the matrix unit $(\delta_{i,a}\delta_{j,b})_{a,b=1,\dots,N}$ where $N = \dim V^{(1)}$.

${}^L\mathfrak{g} = A_n^{(1)}$:

$$\begin{aligned} e_0 &= E_{n+1,1}, \quad e_i = E_{i,i+1} \quad (1 \leq i \leq n), \\ f_0 &= E_{1,n+1}, \quad f_i = E_{i+1,i} \quad (1 \leq i \leq n), \\ h_0 &= -E_{1,1} + E_{n+1,n+1}, \quad h_i = E_{i,i} - E_{i+1,i+1} \quad (1 \leq i \leq n), \\ h_\rho &= \text{diag}\left(-\frac{n}{2}, \dots, \frac{n}{2}\right), \\ \ell &= \text{diag}(\mu_1, \dots, \mu_{n+1}) \quad \left(\sum_{i=1}^{n+1} \mu_i = 0\right). \end{aligned}$$

${}^L\mathfrak{g} = A_{2n-1}^{(2)}$:

$$\begin{aligned} e_0 &= \frac{1}{2}(E_{1,2n-1} + E_{2,2n}), \quad e_i = E_{i+1,i} + E_{2n+1-i,2n-i} \quad (1 \leq i \leq n-1), \\ e_n &= E_{n+1,n}, \\ f_0 &= 2(E_{2n-1,1} + E_{2n,2}), \quad f_i = E_{i,i+1} + E_{2n-i,2n+1-i} \quad (1 \leq i \leq n-1), \\ f_n &= E_{n,n+1}, \\ h_0 &= E_{1,1} + E_{2,2} - E_{2n-1,2n-1} - E_{2n,2n}, \\ h_i &= -E_{i,i} + E_{i+1,i+1} - E_{2n-i,2n-i} + E_{2n+1-i,2n+1-i} \quad (1 \leq i \leq n-1), \\ h_n &= -E_{n,n} + E_{n+1,n+1}, \\ h_\rho &= \text{diag}\left(-n + \frac{1}{2}, \dots, -\frac{1}{2}, \frac{1}{2}, \dots, n - \frac{1}{2}\right), \\ \ell &= \text{diag}(\mu_1, \dots, \mu_n, -\mu_n, \dots, -\mu_1). \end{aligned}$$

${}^L\mathfrak{g} = D_{n+1}^{(2)}$:

$$\begin{aligned} e_0 &= E_{n+2,1} + E_{2n+2,n+2}, \quad e_i = E_{i,i+1} + E_{2n+2-i,2n+3-i} \quad (1 \leq i \leq n-1), \\ e_n &= 2(E_{n,n+1} + E_{n+1,n+3}), \\ f_0 &= 2(E_{1,n+2} + E_{n+2,2n+2}), \quad f_i = E_{i+1,i} + E_{2n+3-i,2n+2-i} \quad (1 \leq i \leq n-1), \\ f_n &= E_{n+1,n} + E_{n+3,n+1}, \\ h_0 &= 2(-E_{1,1} + E_{2n+2,2n+2}), \\ h_i &= E_{i,i} - E_{i+1,i+1} + E_{2n+2-i,2n+2-i} - E_{2n+3-i,2n+3-i} \quad (1 \leq i \leq n-1), \\ h_n &= 2(E_{n,n} - E_{n+3,n+3}), \\ h_\rho &= \text{diag}(-n, \dots, -1, 0, 0, 1, \dots, n), \\ \ell &= \text{diag}(\mu_1, \dots, \mu_n, 0, 0, -\mu_n, \dots, -\mu_1). \end{aligned}$$

${}^L\mathfrak{g} = D_n^{(1)}$:

$$\begin{aligned} e_0 &= \frac{1}{2}(E_{2n-1,1} + E_{2n,2}), \quad e_i = E_{i,i+1} + E_{2n-i,2n+1-i} \quad (1 \leq i \leq n-1), \\ e_n &= 2(E_{n-1,n+1} + E_{n,n+2}), \\ f_0 &= 2(E_{1,2n-1} + E_{2,2n}), \quad f_i = E_{i+1,i} + E_{2n+1-i,2n-i} \quad (1 \leq i \leq n-1), \\ f_n &= \frac{1}{2}(E_{n+1,n-1} + E_{n+2,n}), \end{aligned}$$

$$\begin{aligned}
h_0 &= -E_{1,1} - E_{2,2} + E_{2n-1,2n-1} + E_{2n,2n}, \\
h_i &= E_{i,i} - E_{i+1,i+1} + E_{2n-i,2n-i} - E_{2n+1-i,2n+1-i} \quad (1 \leq i \leq n-1), \\
h_n &= E_{n-1,n-1} + E_{n,n} - E_{n+1,n+1} - E_{n+2,n+2}, \\
h_\rho &= \text{diag}(-n+1, \dots, 0, 0, \dots, n-1), \\
\ell &= \text{diag}(\mu_1, \dots, \mu_n, -\mu_n, \dots, -\mu_1).
\end{aligned}$$

APPENDIX B. EXAMPLE OF PROCEDURE TO DEDUCE ψ -SYSTEM

In this section let $\mathfrak{g} = A_2^{(1)}$. We give more details about the procedure to obtain the ψ -system.

Using the matrix realization of the Chevalley generators in Appendix A, we find that the $L\mathfrak{g}$ -valued connection \mathcal{L} equals

$$\begin{aligned}
\mathcal{L} &= \frac{d}{dx} - \frac{\ell}{x} + E_{1,2} + E_{2,3} + PE_{3,1}, \\
P(x, E) &= x^{3M} - E, \quad (M > 0, E \in \mathbb{C}).
\end{aligned}$$

Let $V^{(1)} = \mathbb{C}^3$ be the vector representation of $L\mathfrak{g}$ with the standard basis $\{v_1, v_2, v_3\}$. Since $V^{(2)} = V_{\frac{1}{2}}^{(1)} \wedge V_{\frac{1}{2}}^{(1)}$, one can chose $u_1 = v_1 \wedge v_2$, $u_2 = v_1 \wedge v_3$ and $u_3 = v_2 \wedge v_3$ as a basis of $V^{(2)}$.

Then the actions of \mathcal{L} on the space $V^{(1)}$ and $V^{(2)}$ are given as

$$\begin{aligned}
\mathcal{L}^{(1)} &= \frac{d}{dx} - \frac{1}{x} \begin{pmatrix} \mu_1 & 0 & 0 \\ 0 & \mu_2 & 0 \\ 0 & 0 & \mu_3 \end{pmatrix} + \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ P & 0 & 0 \end{pmatrix}, \\
\mathcal{L}^{(2)} &= \frac{d}{dx} - \frac{1}{x} \begin{pmatrix} \mu_1 + \mu_2 & 0 & 0 \\ 0 & \mu_1 + \mu_3 & 0 \\ 0 & 0 & \mu_2 + \mu_3 \end{pmatrix} + \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ P & 0 & 0 \end{pmatrix},
\end{aligned}$$

respectively. The element Λ acts on $V^{(1)}$ and $V^{(2)}$ both as

$$\begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix}.$$

Hence it is easy to get $\mu^{(1)} = \mu^{(2)} = 1$, and the corresponding eigenvectors $\mathbf{u}^{(1)} = v_1 + v_2 + v_3$, $\mathbf{u}^{(2)} = u_1 + u_2 + u_3$.

Let $\psi^{(a)}(x, E)$ be the $V^{(a)}$ -valued canonical solution characterized by the behavior at ∞ :

$$\psi^{(a)}(x, E) = e^{-\frac{\mu^{(a)}}{M+1}x^{M+1}} x^{-Mh_\rho} \left(\mathbf{u}^{(a)} + o(1) \right) \quad (x \rightarrow \infty).$$

For $a = 1, 2$, we have the embedding of $L\mathfrak{g}$ -modules

$$(B.1) \quad \iota : \bigwedge^2 V_{\frac{1}{2}}^{(a)} \hookrightarrow V^{(a-1)} \otimes V^{(a+1)}$$

where $V^{(0)} = V^{(3)} = \mathbb{C}$. On the space $V^{(a-1)} \otimes V^{(a+1)}$, the functions $\phi = \iota(\psi_{-\frac{1}{2}}^{(a)} \wedge \psi_{\frac{1}{2}}^{(a)})$ and $\phi = \psi^{(a-1)} \otimes \psi^{(a+1)}$ both satisfy the equation $\mathcal{L}\phi = 0$ which is obvious by $\mathcal{L}\psi^{(i)} = 0$ ($i = a, a \pm 1$). Considering the behaviors at ∞ , we have

$$\begin{aligned}
&\iota(\psi_{-\frac{1}{2}}^{(a)} \wedge \psi_{\frac{1}{2}}^{(a)}) \\
&= e^{-\frac{\mu^{(a)}}{M+1}x^{M+1}(\omega^{-\frac{M+1}{2}} + \omega^{\frac{M+1}{2}})} x^{-Mh_\rho} \left((\omega^{-\frac{M+1}{2}} h_\rho \mathbf{u}^{(a)}) \wedge (\omega^{\frac{M+1}{2}} h_\rho \mathbf{u}^{(a)}) \right) \quad (x \rightarrow \infty),
\end{aligned}$$

and

$$\psi^{(a-1)} \otimes \psi^{(a+1)} = e^{-\frac{\mu^{(a-1)} + \mu^{(a+1)}}{M+1}x^{M+1}} x^{-Mh_\rho} (\mathbf{u}^{(a-1)} \otimes \mathbf{u}^{(a+1)}) \quad (x \rightarrow \infty).$$

A direct computation gives $\mu^{(a)}(\omega^{-\frac{M+1}{2}} + \omega^{\frac{M+1}{2}}) = 2\mu^{(a)} \cos \frac{\pi}{3} = \mu^{(a-1)} + \mu^{(a+1)}$, and

$$\iota((\omega^{-\frac{M+1}{2}} h_\rho \mathbf{u}^{(a)}) \wedge (\omega^{\frac{M+1}{2}} h_\rho \mathbf{u}^{(a)})) = c \mathbf{u}^{(a-1)} \otimes \mathbf{u}^{(a+1)}$$

with c a constant multiple. Hence we find that the functions $\phi = \iota(\psi_{-\frac{1}{2}}^{(a)} \wedge \psi_{\frac{1}{2}}^{(a)})$ and $\phi = \psi^{(a-1)} \otimes \psi^{(a+1)}$ also have the same behavior

$$\phi(x, E) = O \left(\exp \left(-\frac{2\mu^{(a)} \cos \frac{\pi}{h\nu}}{M+1} x^{M+1} \right) \right) \quad (x \rightarrow \infty)$$

in a sector containing $x > 0$. Since such a solution is unique upto a constant multiple, we conclude that (after adjusting the constant multiple)

$$\iota(\psi_{-\frac{1}{2}}^{(a)} \wedge \psi_{\frac{1}{2}}^{(a)}) = \psi^{(a-1)} \otimes \psi^{(a+1)} \quad (a = 1, 2).$$

APPENDIX C. PROOF OF THEOREM 5.2

In this section we give a proof of the ‘quantum Wronskian identity’ in Theorem 5.2. The first algebraic proof of this identity for Q -operators goes back to the work [10] which uses the q -oscillator representations of the Borel subalgebra of $U_q(\widehat{\mathfrak{sl}}_2)$. Subsequently this construction has been extended to $U_q(\widehat{\mathfrak{sl}}_3)$ in [11]. A direct approach using the L operators has been developed recently in the case of Yangian $Y(\mathfrak{gl}_N)$. However we have not been able to find a proof for $U_q(\widehat{\mathfrak{sl}}_n)$ written in the literature. We give here a direct, computational proof as an identity for the q -characters. As mentioned in the introduction, the injectivity of the q -character morphism then ensures that the identity holds true for Q -operators.

Let $\mathfrak{g} = A_n^{(1)}$, and set

$$\mathcal{A}_{i,z} = y_{i,q^{-1}z}^{-1} y_{i,qz}^{-1} y_{i-1,z} y_{i+1,z}, \quad i = 1, \dots, n.$$

One can rewrite $\mathcal{Q}_{i,z}^{(1)}$ as

$$\mathcal{Q}_{i,z}^{(1)} = y_{n,q^{n+1}z}^{-1} \sum_{k_{i+1}, \dots, k_{n+1}} \prod_{j=i}^n \left(\left(\frac{x_{j+1}}{x_i} \right)^{k_{j+1}} \mathcal{A}_{j,q^{j-2} \sum_{\mu=i}^{n+1} k_\mu z}^{-1} \right).$$

For $i \in \mathcal{J}$, we define

$$\begin{aligned} \Delta_{i,z}^{(0)} &= y_{n,q^{n+1}z} \mathcal{Q}_{i,z}^{(1)}, \\ \Delta_{i,z}^{(a)} &= \mathcal{A}_{n+1-a,q^{n+1-a}z}^{-1} \left(\Delta_{i,z}^{(a-1)} - \Delta_{i,q^{-2}z}^{(a-1)} \frac{x_{n+2-a}}{x_i} \right), \quad a = 1, \dots, n. \end{aligned}$$

By the above definition and the formula of $\mathcal{Q}_{i,z}^{(1)}$, we have the following lemma.

Lemma C.1. *For $i \in \mathcal{J}$ and $0 \leq a \leq n$, we have*

$$\Delta_{i,z}^{(a)} = \begin{cases} \sum_{k_{i+1}, \dots, k_{n+1-a}} \prod_{j=i}^{n-a} \left(\left(\frac{x_{j+1}}{x_i} \right)^{k_{j+1}} \mathcal{A}_{j,q^{j-2} \sum_{\mu=i}^{n+1-a} k_\mu z}^{-1} \right) & a < n+1-i, \\ 1 & a = n+1-i, \\ 0 & a > n+1-i. \end{cases}$$

Proposition C.2. *For an element $J = (j_1, j_2, \dots, j_a) \in \mathcal{J}^a$, we have*

$$\mathcal{Q}_{J,z}^{(a)} = y_{n+1-a,q^{n+1}z}^{-1} \begin{vmatrix} \Delta_{j_1,q^{1-a}z}^{(0)} x_{j_1}^0 & \Delta_{j_2,q^{1-a}z}^{(0)} x_{j_2}^0 & \cdots & \Delta_{j_a,q^{1-a}z}^{(0)} x_{j_a}^0 \\ \Delta_{j_1,q^{3-a}z}^{(1)} x_{j_1} & \Delta_{j_2,q^{3-a}z}^{(1)} x_{j_2} & \cdots & \Delta_{j_a,q^{3-a}z}^{(1)} x_{j_a} \\ \vdots & \vdots & \ddots & \vdots \\ \Delta_{j_1,q^{a-1}z}^{(a-1)} x_{j_1}^{a-1} & \Delta_{j_2,q^{a-1}z}^{(a-1)} x_{j_2}^{a-1} & \cdots & \Delta_{j_a,q^{a-1}z}^{(a-1)} x_{j_a}^{a-1} \end{vmatrix}.$$

Proof: From the definition we have

$$\mathcal{Q}_{J,z}^{(a)} = \prod_{k=1}^a y_{n,q^{n-a+2k}z}^{-1} \begin{vmatrix} \Delta_{j_1,q^{1-a}z}^{(0)} x_{j_1}^0 & \Delta_{j_2,q^{1-a}z}^{(0)} x_{j_2}^0 & \cdots & \Delta_{j_a,q^{1-a}z}^{(0)} x_{j_a}^0 \\ \Delta_{j_1,q^{3-a}z}^{(0)} x_{j_1}^0 & \Delta_{j_2,q^{3-a}z}^{(0)} x_{j_2}^0 & \cdots & \Delta_{j_a,q^{3-a}z}^{(0)} x_{j_a}^0 \\ \vdots & \vdots & \ddots & \vdots \\ \Delta_{j_1,q^{a-1}z}^{(0)} x_{j_1}^{a-1} & \Delta_{j_2,q^{a-1}z}^{(0)} x_{j_2}^{a-1} & \cdots & \Delta_{j_a,q^{a-1}z}^{(0)} x_{j_a}^{a-1} \end{vmatrix}.$$

For $k = 2, \dots, a$, we subtract the $(k-1)$ -th row multiplied by x_{n+1} from the k -th row and then extract the factor $\mathcal{A}_{n,q^{n-a-1+2k}z}^{-1}$. Thus we get

$$\begin{aligned} \mathcal{Q}_{I,z}^{(a)} &= \prod_{k=1}^a y_{n,q^{n-a+2k}z}^{-1} \prod_{k=2}^a \mathcal{A}_{n,q^{n-a-1+2k}z}^{-1} \\ &\times \begin{vmatrix} \Delta_{j_1,q^{1-a}z}^{(0)} x_{j_1}^0 & \Delta_{j_2,q^{1-a}z}^{(0)} x_{j_2}^0 & \cdots & \Delta_{j_a,q^{1-a}z}^{(0)} x_{j_a}^0 \\ \Delta_{j_1,q^{3-a}z}^{(1)} x_{j_1}^0 & \Delta_{j_2,q^{3-a}z}^{(1)} x_{j_2}^0 & \cdots & \Delta_{j_a,q^{3-a}z}^{(1)} x_{j_a}^0 \\ \vdots & \vdots & \ddots & \vdots \\ \Delta_{j_1,q^{a-1}z}^{(1)} x_{j_1}^{a-1} & \Delta_{j_2,q^{a-1}z}^{(1)} x_{j_2}^{a-1} & \cdots & \Delta_{j_a,q^{a-1}z}^{(1)} x_{j_a}^{a-1} \end{vmatrix}. \end{aligned}$$

Repeating the above steps in the order $b = 2, \dots, a-1$, by replacing x_{n+1} and $\mathcal{A}_{n,q^{n-a-1+2k}z}^{-1}$ with x_{n+2-b} and $\mathcal{A}_{n+1-b,q^{n-a-b+2k}z}^{-1}$ ($k = b+1, \dots, a$), we get

$$\begin{aligned} \mathcal{Q}_{I,z}^{(a)} &= \prod_{k=1}^a y_{n,q^{n-a+2k}z}^{-1} \prod_{b=1}^{a-1} \prod_{k=b+1}^a \mathcal{A}_{n+1-b,q^{n-a-b+2k}z}^{-1} \\ &\times \begin{vmatrix} \Delta_{j_1,q^{1-a}z}^{(0)} x_{j_1}^0 & \Delta_{j_2,q^{1-a}z}^{(0)} x_{j_2}^0 & \cdots & \Delta_{j_a,q^{1-a}z}^{(0)} x_{j_a}^0 \\ \Delta_{j_1,q^{3-a}z}^{(1)} x_{j_1}^0 & \Delta_{j_2,q^{3-a}z}^{(1)} x_{j_2}^0 & \cdots & \Delta_{j_a,q^{3-a}z}^{(1)} x_{j_a}^0 \\ \vdots & \vdots & \ddots & \vdots \\ \Delta_{j_1,q^{a-1}z}^{(a-1)} x_{j_1}^{a-1} & \Delta_{j_2,q^{a-1}z}^{(a-1)} x_{j_2}^{a-1} & \cdots & \Delta_{j_a,q^{a-1}z}^{(a-1)} x_{j_a}^{a-1} \end{vmatrix}. \end{aligned}$$

The proof is over by noting that

$$\prod_{k=1}^a y_{n,q^{n-a+2k}z}^{-1} \prod_{b=1}^{a-1} \prod_{k=b+1}^a \mathcal{A}_{n+1-b,q^{n-a-b+2k}z}^{-1} = y_{n+1-a,q^{n+1}z}^{-1}.$$

□

Now we give a proof of Theorem 5.2.

Proof of Theorem 5.2: By Proposition C.2, for $J = (1, 2, \dots, n+1)$, we have

$$\mathcal{Q}_{J,z}^{(n+1)} = \begin{vmatrix} \Delta_{1,q^{-n}z}^{(0)} x_1^0 & \Delta_{2,q^{-n}z}^{(0)} x_2^0 & \cdots & \Delta_{n+1,q^{-n}z}^{(0)} x_{n+1}^0 \\ \Delta_{1,q^{2-n}z}^{(1)} x_1^0 & \Delta_{2,q^{2-n}z}^{(1)} x_2^0 & \cdots & \Delta_{n+1,q^{2-n}z}^{(1)} x_{n+1}^0 \\ \vdots & \vdots & \ddots & \vdots \\ \Delta_{1,q^n z}^{(n)} x_1^n & \Delta_{2,q^n z}^{(n)} x_2^n & \cdots & \Delta_{n+1,q^n z}^{(n)} x_{n+1}^n \end{vmatrix}.$$

By Lemma C.1, we get

$$\mathcal{Q}_{J,z}^{(n+1)} = (-1)^{(n+1)n/2} \prod_{b=1}^n x_b^{n+1-b},$$

which implies the result. □

APPENDIX D. PROOFS OF IDENTITIES FOR CHARACTERS IN THE CASE $B_n^{(1)}$

In this section, we set $\mathfrak{g} = B_n^{(1)}$.

First we introduce some notation. For $i, j \in \mathcal{J}$ and an element $J \in \mathcal{J}^a$, we define

$$h_J = \prod_{\substack{k, \bar{k} \in J, \\ k \prec \bar{k}}} (x_k^{\frac{1}{2}} + x_k^{-\frac{1}{2}}), \quad \tilde{x}_{i,J} = \prod_{\substack{k \in J, \\ k \prec i}} (-x_i^{\frac{1}{2}\delta_{i,k}-1}) \prod_{\substack{k \in J^*, \\ i \prec k}} (x_k^{1-\frac{1}{2}\delta_{k,i}}),$$

$$g_{i,j} = f_{i,j} x_j^{1-\frac{1}{2}\delta_{j,i}} = -g_{j,i}, \quad g_{i,J} = \prod_{\substack{k \in J, \\ k \neq i}} g_{i,k},$$

where $f_{i,j}$ is defined in Subsection 4.3.

By (4.15) and the above notation, one has

$$\bar{Q}_J^{(a)} = \prod_{\substack{i \in J, j \in J^*, \\ i \prec j}} f_{i,j} \prod_{\substack{i, j \in J, \\ i \prec j}} ((x_j - x_i) f_{i,j}) = X_J^{-1} h_J \prod_{\substack{i \in J, j \in J^*, \\ i \prec j}} f_{i,j},$$

where X_J is given in (5.5). A direct computation gives

$$\prod_{\substack{i \in J^*, j \in J, \\ i \prec j}} f_{i,j} = \prod_{\substack{i \in J, j \in J^*, \\ i \prec j}} f_{i,j}, \quad \prod_{j \in J, i \prec j} f_{i,j} = \prod_{j \in J, j \prec i} f_{j,i},$$

$$\frac{\prod_{\substack{i \in J-j_k, j \in (J-j_k)^*, \\ i \prec j}} f_{i,j}}{\prod_{\substack{i \in J, j \in J^*, \\ i \prec j}} f_{i,j}} = \frac{\prod_{\substack{i \in J, \\ i \prec j_k}} g_{j_k, j}}{\prod_{\substack{j \in J^*, \\ j_k \prec j}} g_{j_k, j}} \tilde{x}_{j_k, J}.$$

Furthermore, for two elements J_1 and J_2 of \mathcal{J}^n and $j_k \in J_2 \cap J_1^*$, we have

$$\frac{x_{j_k, J_1} \tilde{x}_{j_k, J_2}}{x_{j_k, J_2} \tilde{x}_{j_k, J_1}} = -x_{j_k}^2 x_{J_1} x_{J_2}^{-1}.$$

Using the above formulas, the specialization of Conjecture 5.3 to the characters $\bar{Q}_J^{(a)}$ reduces to the following proposition.

Proposition D.1. *For two increasing elements $J_1 = (i_1, \dots, i_n)$, $J_2 = (j_1, \dots, j_n) \in \mathcal{J}^n$, we have*

$$(D.1) \quad x_{J_1}^{-\frac{1}{2}} h_{J_1^*} h_{J_2} - x_{J_2}^{-\frac{1}{2}} h_{J_2^*} h_{J_1} = x_{J_1}^{-1} x_{J_2}^{-1} \sum_{j \in J_2 \cap J_1^*} x_{J_1}^{\frac{3}{2}} x_j^{\frac{3}{2}} \frac{g_{j, J_1^*}}{g_{j, J_2^*}} h_{J_2-j} h_{J_1^*-j}.$$

Proof: Generally let

$$T_1 = \left\{ j \in J_1 \cap J_2^* \mid \bar{j} \notin J_1, \bar{j} \in J_2^* \text{ or } \bar{j} \in J_1, \bar{j} \notin J_2^* \right\},$$

$$T_2 = \left\{ j \in J_1 \cap J_2^* \mid \bar{j} \notin J_1, \bar{j} \notin J_2^* \right\}, \quad T_3 = \left\{ j \in J_1 \cap J_2^* \mid \bar{j} \in J_1 \cap J_2^*, j \prec \bar{j} \right\},$$

$$S_1 = \left\{ j \in J_2 \cap J_1^* \mid \bar{j} \notin J_2, \bar{j} \in J_1^*, \text{ or } \bar{j} \in J_2, \bar{j} \notin J_1^* \right\},$$

$$S_2 = \left\{ j \in J_2 \cap J_1^* \mid \bar{j} \notin J_2, \bar{j} \notin J_1^* \right\}, \quad S_3 = \left\{ j \in J_2 \cap J_1^* \mid \bar{j} \in J_2 \cap J_1^*, j \prec \bar{j} \right\},$$

then we have $J_1 \cap J_2^* = \bigcup_{1 \leq i \leq 3} T_i \cup \bar{T}_3$, $J_2 \cap J_1^* = \bigcup_{1 \leq i \leq 3} S_i \cup \bar{S}_3$, and $\#(J_1 \cap J_2^*) = \#(J_2 \cap J_1^*)$. Here

for a subset $S \subseteq \mathcal{J}$, we set $\bar{S} = \{\bar{i} \mid i \in S\}$.

Let

$$F_{J_1, J_2}(x_j) = x_j^{\frac{3}{2}} x_{J_1} x_{J_2}^{-1} \frac{h_{J_2-j} h_{J_1^*-j} g_{j, J_1^*}}{h_{J_1^*} h_{J_2} g_{j, J_2^*}}.$$

We also introduce a function

$$f(z) = \frac{1}{(1+z)} \frac{\prod_{k \in T_1} (z - x_k)}{\prod_{k \in S_1} (z - x_k)} \prod_{k \in S_2} \frac{zx_k - 1}{z - x_k} \frac{\prod_{k \in T_3} (z - x_k)(zx_k - 1)}{\prod_{k \in S_3} (z - x_k)(zx_k - 1)},$$

which has poles only at $z = -1, \infty$ and x_j ($j \in J_2 \cap J_1^* = \bigcup_{1 \leq i \leq 3} S_i \cup \bar{S}_3$).

A direct computation gives, for $j \in S_3$

$$\begin{aligned} & \text{Res}(f, x_j^{-1}) \\ &= \frac{1}{x_j(1+x_j^2)(1-x_j)} \frac{\prod_{k \in T_1} (1 - x_k x_j)}{\prod_{k \in S_1} (1 - x_k x_j)} \prod_{k \in S_2} \frac{x_k - x_j}{1 - x_k x_j} \frac{\prod_{k \in T_3} (x_j - x_k)(x_j x_k - 1)}{\prod_{\substack{k \in S_3, \\ k \neq j}} (x_j - x_k)(x_j x_k - 1)} \\ &= c F_{J_1, J_2}(x_j^{-1}), \end{aligned}$$

where $c = \frac{x_{T_3} x_{S_2}}{x_{S_3}}$. Similarly we obtain

$$\text{Res}(f, x_j) = c F_{J_1, J_2}(x_j) \quad (j \in S_i, i = 1, 2, 3).$$

On the other hand, by setting $h_i = x_i^{\frac{1}{2}} + x_i^{-\frac{1}{2}}$, it is easy to see that

$$\text{Res}(f, -1) = \frac{x_{T_3}}{x_{S_3}} \left(\frac{x_{T_1}}{x_{S_1}} \right)^{\frac{1}{2}} \frac{\prod_{k \in T_1} h_k}{\prod_{k \in S_1} h_k} \frac{\prod_{k \in T_3} h_k^2}{\prod_{k \in S_3} h_k^2}, \quad \text{Res}(f, \infty) = -\frac{x_{T_3} x_{S_2}}{x_{S_3}} = -c.$$

Since the right-hand side of (D.1) equals $h_{J_1^*} h_{J_2} x_{J_1}^{-\frac{1}{2}} \sum_{j \in J_2 \cap J_1^*} F_{J_1, J_2}(x_j)$, it simplifies to

$$h_{J_1^*} h_{J_2} x_{J_1}^{-\frac{1}{2}} (-c^{-1}) \left(-c + \frac{x_{T_3}}{x_{S_3}} \left(\frac{x_{T_1}}{x_{S_1}} \right)^{\frac{1}{2}} \frac{\prod_{k \in T_1} h_k}{\prod_{k \in S_1} h_k} \frac{\prod_{k \in T_3} h_k^2}{\prod_{k \in S_3} h_k^2} \right).$$

Now it is sufficient to prove

$$\left(\frac{x_{T_1}}{x_{S_1}} \right)^{\frac{1}{2}} \frac{\prod_{k \in T_1} h_k}{\prod_{k \in S_1} h_k} \frac{\prod_{k \in T_3} h_k^2}{\prod_{k \in S_3} h_k^2} = x_{S_2} \frac{x_{J_1}^{\frac{1}{2}} h_{J_2^*} h_{J_1}}{x_{J_2}^{\frac{1}{2}} h_{J_1^*} h_{J_2}}$$

which is immediate. We finish the proof. \square

APPENDIX E. PROOFS FOR IDENTITIES OF CHARACTERS OF CASE $C_n^{(1)}$ AND $D_n^{(1)}$

In this section we consider $\mathfrak{g} = C_n^{(1)}, D_n^{(1)}$.

By (4.15) and (4.4), Conjectures 5.5-5.8 reduce to the following identities.

$$\begin{aligned} \text{(E.1)} \quad & (-1)^{\frac{u(u+1)}{2}} x_{\sigma}^2 \prod_{\substack{1 \leq k \leq u, \\ 1 \leq l \leq u+1}} \frac{1}{(1 - x_{\sigma_k} x_{\eta_l})(x_{\sigma_k} - x_{\eta_l})} \prod_{1 \leq k \leq u+1} (1 - x_{\eta_k}^2) \\ &= \sum (-1)^{\frac{s(s+1)}{2}} x_{\gamma}^2 \prod_{1 \leq k < l \leq s} \frac{1}{1 - x_{\gamma_k} x_{\gamma_l}} \prod_{1 \leq k < l \leq t} \frac{1}{1 - x_{\delta_k} x_{\delta_l}} \prod_{\substack{1 \leq k \leq s, \\ 1 \leq l \leq t}} \frac{1}{x_{\gamma_k} - x_{\delta_l}}, \end{aligned}$$

$$\text{(E.2)} \quad (-1)^{\frac{u(u-1)}{2}} x_{\sigma}^2 \prod_{\substack{1 \leq k \leq u, \\ 1 \leq l \leq u+2}} \frac{1}{(1 - x_{\sigma_k} x_{\eta_l})(x_{\sigma_k} - x_{\eta_l})} \prod_{1 \leq k \leq u} (1 - x_{\sigma_k}) \prod_{1 \leq k \leq u+2} (1 + x_{\eta_k})$$

$$\begin{aligned}
&= \sum (-1)^{\frac{s(s-1)}{2}} x_\gamma^2 \prod_{1 \leq k < l \leq s} \frac{1}{1 - x_{\gamma_k} x_{\gamma_l}} \prod_{1 \leq k < l \leq t} \frac{1}{1 - x_{\delta_k} x_{\delta_l}} \prod_{\substack{1 \leq k \leq s, \\ 1 \leq l \leq t}} \frac{1}{x_{\gamma_k} - x_{\delta_l}}, \\
(E.3) \quad &(-1)^{\frac{u(u+1)}{2}} x_\sigma \prod_{\substack{1 \leq k \leq u, \\ 1 \leq l \leq u}} \frac{1}{(1 - x_{\sigma_k} x_{\eta_l})(x_{\sigma_k} - x_{\eta_l})} \prod_{1 \leq k \leq u} (1 - x_{\eta_k}^2) \\
&= \sum (-1)^{\frac{s(s+1)}{2}} x_\gamma \prod_{1 \leq k < l \leq s} \frac{1}{1 - x_{\gamma_k} x_{\gamma_l}} \prod_{1 \leq k < l \leq t} \frac{1}{1 - x_{\delta_k} x_{\delta_l}} \prod_{\substack{1 \leq k \leq s, \\ 1 \leq l \leq t}} \frac{1}{x_{\gamma_k} - x_{\delta_l}}, \\
(E.4) \quad &(-1)^{\frac{u(u-1)}{2}} x_\sigma \prod_{\substack{1 \leq k \leq u, \\ 1 \leq l \leq u+1}} \frac{1}{(1 - x_{\sigma_k} x_{\eta_l})(x_{\sigma_k} - x_{\eta_l})} \prod_{1 \leq k \leq u} (1 - x_{\sigma_k}) \prod_{1 \leq k \leq u+1} (1 + x_{\eta_k}) \\
&= \sum (-1)^{\frac{s(s-1)}{2}} x_\gamma \prod_{1 \leq k < l \leq s} \frac{1}{1 - x_{\gamma_k} x_{\gamma_l}} \prod_{1 \leq k < l \leq t} \frac{1}{1 - x_{\delta_k} x_{\delta_l}} \prod_{\substack{1 \leq k \leq s, \\ 1 \leq l \leq t}} \frac{1}{x_{\gamma_k} - x_{\delta_l}}.
\end{aligned}$$

In these formulas, σ, η are disjoint subsets of $\{1, \dots, n\}$, such that $(\#\sigma, \#\eta) = (u, u+1), (u, u+2), (u, u), (u, u+1)$ respectively. The sum is taken over partitions of $\sigma \cup \eta$ into subsets γ, δ satisfying $\sigma \subseteq \gamma$, where we set $s = \#\gamma, t = \#\delta$.

Similarly, in the case of $g = D_n^{(1)}$, by (4.15) and (4.5) the relevant identities read as follows.

$$\begin{aligned}
(E.5) \quad &(-1)^{\frac{u(u+1)}{2}} x_\sigma \prod_{\substack{1 \leq k \leq u, \\ 1 \leq l \leq u+2}} \frac{1}{(1 - x_{\sigma_k} x_{\eta_l})(x_{\sigma_k} - x_{\eta_l})} \prod_{1 \leq k \leq u} (1 - x_{\sigma_k}^2) \\
&= \sum (-1)^{\frac{t(t+1)}{2}} x_\delta \prod_{1 \leq k < l \leq s} \frac{1}{1 - x_{\gamma_k} x_{\gamma_l}} \prod_{1 \leq k < l \leq t} \frac{1}{1 - x_{\delta_k} x_{\delta_l}} \prod_{\substack{1 \leq k \leq s, \\ 1 \leq l \leq t}} \frac{1}{x_{\gamma_k} - x_{\delta_l}}, \\
(E.6) \quad &(-1)^{\frac{u(u-1)}{2}} \prod_{\substack{1 \leq k \leq u, \\ 1 \leq l \leq u+2}} \frac{1}{(1 - x_{\sigma_k} x_{\eta_l})(x_{\sigma_k} - x_{\eta_l})} \prod_{1 \leq k \leq u} (1 - x_{\sigma_k}^2) \\
&= \sum (-1)^{\frac{s(s-1)}{2}} \prod_{1 \leq k < l \leq s} \frac{1}{1 - x_{\gamma_k} x_{\gamma_l}} \prod_{1 \leq k < l \leq t} \frac{1}{1 - x_{\delta_k} x_{\delta_l}} \prod_{\substack{1 \leq k \leq s, \\ 1 \leq l \leq t}} \frac{1}{x_{\gamma_k} - x_{\delta_l}}.
\end{aligned}$$

Here $(\#\sigma, \#\eta) = (u, u+2), (u, u+1)$ respectively, $s = \#\gamma, t = \#\delta$, and the sum is taken over the partitions satisfying $\sigma \subseteq \gamma$ keeping fixed the parity of s .

These formulas are related by specialization of variables: (E.1) is obtained from (E.3) by setting $x_{\sigma_u} = 0$, (E.2) is obtained from (E.4) by setting $x_{\sigma_u} = 0$, and (E.4) is obtained from (E.3) by setting $x_{\sigma_u} = 1$. Likewise (E.5) is obtained from (E.6) by setting $x_{\sigma_u} = \infty$. Hence it is sufficient to deal only with (E.3) and (E.6).

Define polynomials in $x = (x_1, \dots, x_m)$ and $y = (y_1, \dots, y_n)$,

$$\begin{aligned}
F^{(m,n)}(x, y) &= \prod_{1 \leq i < j \leq m} (1 - x_i x_j) \prod_{1 \leq i < j \leq n} (1 - y_i y_j) (y_j - y_i), \\
G_{J_1, J_2}^{(m,n)}(x, y) &= \prod_{\substack{1 \leq i \leq m, \\ j \in J_1}} (1 - x_i y_j) \prod_{\substack{1 \leq i \leq m, \\ j \in J_2}} (x_i - y_j) \\
&\quad \times \prod_{\substack{i, j \in J_1, \\ i < j}} (y_j - y_i) \prod_{i \in J_1, j \in J_2} (1 - y_i y_j) \prod_{\substack{i, j \in J_2, \\ i < j}} (y_i - y_j).
\end{aligned}$$

Set also $\delta(J_1, J_2) = \#\{(i, j) \in J_1 \times J_2 \mid i < j\}$. Then (E.3) and (E.6) can be rewritten respectively as follows:

$$(E.7) \quad \prod_{j=1}^n (1 - y_j^2) \cdot F^{(n,n)}(x, y) = \sum_{J_1, J_2} (-1)^{\delta(J_1, J_2)} \prod_{j \in J_2} y_j \cdot G_{J_1, J_2}^{(n,n)}(x, y),$$

$$(E.8) \quad (1 + \epsilon) \prod_{i=1}^{n-1} (1 - x_i^2) \cdot F^{(n-1,n)}(x, y) = \sum_{J_1, J_2} \epsilon^{\#J_2} (-1)^{\delta(J_1, J_2)} G_{J_1, J_2}^{(n-1,n)}(x, y).$$

Here the sum in the right-hand sides are taken over all partitions $J_1 \sqcup J_2$ of $\{1, \dots, n\}$, and $\epsilon = \pm 1$. For the second, we have added/subtracted the original sums which have a parity restriction.

In order to show (E.7), we prove it in a slightly more general form.

Proposition E.1. *Suppose $n \geq 2$, $0 \leq m \leq n$ and $m \equiv n \pmod{2}$. Then*

$$(E.9) \quad \delta_{m,n} \prod_{j=1}^n (1 - y_j^2) \cdot F^{(m,n)}(x, y) = \sum_{J_1, J_2} (-1)^{\delta(J_1, J_2)} \prod_{j \in J_2} y_j \cdot G_{J_1, J_2}^{(m,n)}(x, y).$$

Proof: Denote by $R^{(m,n)}$ the right-hand side of (E.9), and set

$$L^{(m,n)} = \prod_{j=1}^n (1 - y_j^2) \cdot F^{(m,n)}(x, y), \quad h^{(m)}(x, z) = \prod_{i=1}^m (1 - x_i z).$$

Let us rewrite $R^{(m,n)}$ as follows:

$$\begin{aligned} & \prod_{j=1}^n y_j^{-(m+n)/2} \cdot R^{(m,n)} \\ &= \sum_{J_1, J_2} (-1)^{\#J_2} \prod_{i \in J_1} (y_i^{-(m+n)/2} h^{(m)}(x, y_i)) \prod_{j \in J_2} (y_j^{(m+n)/2} h^{(m)}(x, y_j^{-1})) \\ & \times \prod_{\substack{i, j \in J_1, \\ i < j}} (y_j - y_i) \prod_{\substack{i \in J_1, j \in J_2, \\ i < j}} (y_j^{-1} - y_i) \prod_{\substack{i \in J_1, j \in J_2, \\ i > j}} (y_i - y_j^{-1}) \prod_{\substack{i, j \in J_2, \\ i < j}} (y_j^{-1} - y_i^{-1}) \\ &= \sum_{\epsilon_1, \dots, \epsilon_n = \pm 1} \prod_{j=1}^n (\epsilon_j y_j^{-\epsilon_j(m+n)/2} h^{(m)}(x, y_j^{\epsilon_j})) \prod_{1 \leq i < j \leq n} (y_j^{\epsilon_j} - y_i^{\epsilon_i}) \\ &= \det \left(y_i^{-(m+n)/2+j-1} h^{(m)}(x, y_i) - y_i^{(m+n)/2-j+1} h^{(m)}(x, y_i^{-1}) \right)_{1 \leq i, j \leq n}. \end{aligned}$$

It is easy to see that the last expression is 0 if $m = 0$ or $m = 1$. Assume $m \geq 2$. We prove (E.9) by induction on m . From the skew-symmetry of the determinant, $R^{(m,n)}$ is divisible by $\prod_{j=1}^n (1 - y_j^2) \prod_{1 \leq i < j \leq n} (1 - y_i y_j) (y_j - y_i)$. Let us show that it is divisible also by $\prod_{1 \leq i < j \leq m} (1 - x_i x_j)$.

If we set $x_m = x_{m-1}^{-1}$, then $h^{(m)}(x, z) = z(z + z^{-1} - x_{m-1} - x_{m-1}^{-1}) h^{(m-2)}(x'', z)$, where $x'' = (x_1, \dots, x_{m-2})$. Hence $R^{(m,n)}$ becomes proportional to

$$\det \left(y_i^{-(m-2+n)/2+j-1} h^{(m-2)}(x'', y_i) - y_i^{(m-2+n)/2-j+1} h^{(m-2)}(x'', y_i^{-1}) \right)_{1 \leq i, j \leq n},$$

which is 0 by the induction hypothesis. This shows that $R^{(m,n)}$ is divisible by $L^{(m,n)}$.

The total degree of $L^{(m,n)}$ is $l^{(m,n)} = m(m-1) + n(3n+1)/2$, while that of $R^{(m,n)}$ is at most $r^{(m,n)} = n^2/2 + mn + (m+n)^2/4 - m/2$. Notice that $l^{(m,n)} > r^{(m,n)}$ for $n > m$ and $l^{(n,n)} = r^{(n,n)}$. From this we conclude that if $n > m$ then $R^{(m,n)} = 0$, and if $n = m$ then $R^{(n,n)}$ is a constant multiple of $L^{(n,n)}$. The constant is shown to be 1 by setting $x = 0$ and using the Weyl denominator formula of type C_n ,

$$\prod_{j=1}^n y_j^n \cdot \det \left(y_i^{-n+j-1} - y_i^{n-j+1} \right)_{1 \leq i, j \leq n} = \prod_{j=1}^n (1 - y_j^2) \prod_{1 \leq i < j \leq n} (1 - y_i y_j) (y_j - y_i).$$

□

Proposition E.2. *Suppose $n \geq 3$ and $0 \leq m \leq n-1$. Then for $\epsilon = \pm 1$ we have*

$$(E.10) \quad (1 + \epsilon) \delta_{m, n-1} \prod_{i=1}^m (1 - x_i^2) \cdot F^{(m,n)}(x, y) = \sum_{J_1, J_2} \epsilon^{\#J_2} (-1)^{\delta(J_1, J_2)} G_{J_1, J_2}^{(m,n)}(x, y).$$

Proof: We denote the right hand-side of (E.10) by $R_\epsilon^{(m,n)}$, and set

$$K^{(m,n)} = \prod_{i=1}^m (1 - x_i^2) \cdot F^{(m,n)}(x, y).$$

By a calculation similar to the one in the proof of Proposition E.1 we find

$$\begin{aligned} & \prod_{j=1}^n y_j^{-(m+n-1)/2} \cdot R_\epsilon^{(m,n)} \\ &= \det \left(y_i^{-(m+n+1)/2+j} h^{(m)}(x, y_i) + \epsilon' y_i^{(m+n+1)/2-j} h^{(m)}(x, y_i^{-1}) \right)_{1 \leq i, j \leq n} \end{aligned}$$

where $\epsilon' = (-1)^{m+n-1} \epsilon$. From this, it is easy to see that $R_\epsilon^{(0,n)} = 0$. As in Proposition E.1, we prove the assertion by induction on m .

As before, one can verify that the left-hand side is divisible by $\prod_{1 \leq i < j \leq m} (1 - x_i x_j) \prod_{1 \leq i < j \leq n} (1 - y_i y_j)(y_j - y_i)$. Let us show that it is divisible also by $\prod_{i=1}^m (1 - x_i^2)$. If we set $x_m = \pm 1$, then we have

$$h^{(m)}(x, z) \Big|_{x_m = \pm 1} = (z^{-1/2} \mp z^{1/2}) \cdot z^{1/2} h^{(m-1)}(x', z), \quad x' = (x_1, \dots, x_{m-1}).$$

Hence the determinant becomes proportional to

$$\det \left(y_i^{-(m+n)/2+j} h^{(m-1)}(x', y_i) \mp \epsilon' y_i^{(m+n)/2-j} h^{(m-1)}(x, y_i^{-1}) \right)_{1 \leq i, j \leq n},$$

which is 0 by the induction hypothesis.

The total degree of $K^{(m,n)}$ is $k^{(m,n)} = m(m+1) + 3n(n-1)/2$, while that of $R_\epsilon^{(m,n)}$ is at most $r_\epsilon^{(m,n)} = mn + n(n-1)/2 + (m+n)^2/4 - \delta/4$, where $\delta = 0, 1$ is determined by $\delta \equiv m - n \pmod{2}$. Then $k^{(m,n)} > r_\epsilon^{(m,n)}$ if $m \leq n - 2$, and $k^{(n-1,n)} = r_\epsilon^{(n-1,n)}$. Hence we find that $R_\epsilon^{(m,n)} = 0$ if $m \leq n - 2$, and $R_\epsilon^{(n-1,n)}$ is a constant multiple of $K^{(n-1,n)}$. The constant can be found by setting $x = 0$ and using the Weyl denominator formula of type D_n ,

$$\prod_{j=1}^n y_j^n \cdot \det \left(y_i^{-n+j} + \epsilon y_i^{n-j} \right)_{1 \leq i, j \leq n} = (1 + \epsilon) \prod_{1 \leq i < j \leq n} (1 - y_i y_j)(y_j - y_i).$$

□

Acknowledgement. The author would like to express her deep gratitude to Professor Michio Jimbo for his valuable advices, great helps and kindness, and to Professor Hidetaka Sakai for his constant encouragement during the course of the present work. She wishes to thank also Professor Junji Suzuki for helpful discussions.

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